

# Magnetic excitations in the spin-glass phase of $\text{La}_{2-x}\text{Sr}_x\text{CuO}_4$

Hideto Goka<sup>a,1</sup>, Shin-ichi Kuroshima<sup>a</sup> Masaki Fujita<sup>a</sup> Kazuyoshi Yamada<sup>a</sup>,  
Haruhiro Hiraka<sup>b</sup>, Yasuo Endoh<sup>b</sup>, Christopher D. Frost<sup>c</sup>

<sup>a</sup>*Institute for Chemical Research, Kyoto University, Gokasyo, Uji 611-0011, Japan*

<sup>b</sup>*Institute of Material Research, Tohoku University, Sendai 980-8577, Japan*

<sup>c</sup>*Rutherford Appleton Laboratory, Chilton, Didcot OX11 0QX, United Kingdom*

## Abstract

Magnetic excitations of  $\text{La}_{1.95}\text{Sr}_{0.05}\text{CuO}_4$  were investigated using pulsed neutron inelastic scattering over a wide range of energy up to 300 meV. In contrast to the optimally doped superconductor the dynamical magnetic susceptibility  $\chi''(\omega)$  around  $(\pi, \pi)$  in the spin-glass phase monotonically decreases with increasing  $\omega$  and nearly saturates beyond  $\sim 50$  meV. We found that the energy dependence of the  $q$ -position in the latter energy region is ascribed to a two-dimensional spin-wave dispersion relation with the nearest neighbor interaction  $J$  of  $108 \pm 6$  meV, which is smaller by  $\sim 20\%$  than that of  $\text{La}_2\text{CuO}_4$ .

*Key words:*  $\text{La}_{2-x}\text{Sr}_x\text{CuO}_4$ ; spin excitation; neutron scattering

## 1. Introduction

It is widely believed that magnetism is relevant for the superconductivity in the high- $T_c$  cuprate. In the hole-doped  $\text{La}_{2-x}\text{Sr}_x\text{CuO}_4$  (LSCO), systematic studies on the low energy spin excitation have been performed by neutron scattering in a wide hole concentration range, and the results of these studies clearly demonstrate a close relation of the low-energy spin excitations with the superconducting transition [1].

On the other hand, the high-energy neutron scattering studies have been reported only for the non-doped and optimally doped LSCO, and these studies revealed distinct feature in the energy spectrum between the nondoped and optimally doped LSCO [2]-[5]. Therefore, particularly, information on the spin excitations for the lightly doped or underdoped LSCO is very important.

Then we started a high-energy neutron scattering study on the spin excitations of  $\text{La}_{1.95}\text{Sr}_{0.05}\text{CuO}_4$

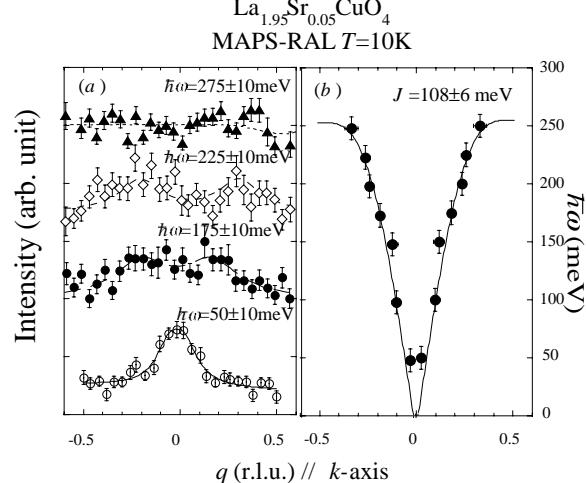


Fig. 1. (a):Peak profiles around  $(\pi, \pi)$  at various energies. (b):The dispersion relation of  $\text{La}_{1.95}\text{Sr}_{0.05}\text{CuO}_4$  in the energy region larger than 50 meV.

which is in the spin-glass phase near the superconducting phase boundary.

<sup>1</sup> Corresponding author. Present address: Institute for Chemical Research, Kyoto University, Gokasyo, Uji 611-0011, Japan  
E-mail: goka@msk.kuicr.kyoto-u.ac.jp

## 2. Experimental

Single crystals of  $\text{La}_{1.95}\text{Sr}_{0.05}\text{CuO}_4$  were grown by a Traveling-Solvent-Floating-Zone (TSFZ) method. Total volume of the assembled crystals is about 10 cc. Pulsed neutron scattering experiments were performed on the chopper spectrometer MAPS newly installed in ISIS Facility at Rutherford Appleton Laboratory. The crystals were aligned with the  $c$ -axis parallel to the incident beam. Measurements were performed using incident neutron energies  $E_i = 21.5, 30, 50, 100, 200$ , and 400 meV. Most of the measurements were performed at  $T=10$  K and for the measurement with  $E_i=50$  meV we studied temperature dependence up to  $T=200$  K.

## 3. Results

Fig. 1(a) shows peak profiles taken by a scan through the position  $(\pi, \pi)$  at various energies between 50 and 275 meV. The peak height is arbitrary. It is shown that the peak split enlarges as the energy increases. The relation between the width of the peak split and the energy-transfer is shown in Fig. 1(b). This relation is well ascribed to the 2D antiferromagnetic dispersion relation. From the least-squares fitting of the data to the relation we evaluated a nearest neighbor magnetic interaction  $J=108\pm6$  meV which is smaller by  $\sim 20\%$  than that reported for non-doped LSCO [2]. However, for the results in the lower energy region below around 50 meV, it is very difficult to apply the abovementioned dispersion relation.

In the middle panel of Fig. 2 (a) we show the energy spectrum of dynamical magnetic susceptibility  $\chi''(\omega)$  of  $\text{La}_{1.95}\text{Sr}_{0.05}\text{CuO}_4$  integrated in  $q$  around  $(\pi, \pi)$ . The  $\chi''(\omega)$  in the spin-glass phase decreases monotonically up to around 50 meV and gradually saturates at higher energy region. Such energy dependence of  $\chi''(\omega)$  is clearly distinct from both data of nondoped and optimally doped samples as shown in the top and bottom panels of Fig. 2 (a), respectively. The upper panel of Fig. 2 (b) shows the in-plane peak-width  $\kappa$  of  $\text{La}_{1.95}\text{Sr}_{0.05}\text{CuO}_4$  as a function of  $\sqrt{T^2 + (\hbar\omega/k_B)^2}$ . We refer the result of  $\kappa$  of  $\text{La}_{1.86}\text{Sr}_{0.14}\text{CuO}_4$  in the lower panel [6].

## 4. Discussion

The present high-energy pulsed neutron experiment newly revealed two features in the magnetic excitations of  $\text{La}_{1.95}\text{Sr}_{0.05}\text{CuO}_4$ . One is that the inelastic response at high energy region beyond around 50 meV is similar to that of the nondoped antiferromagnetic insula-

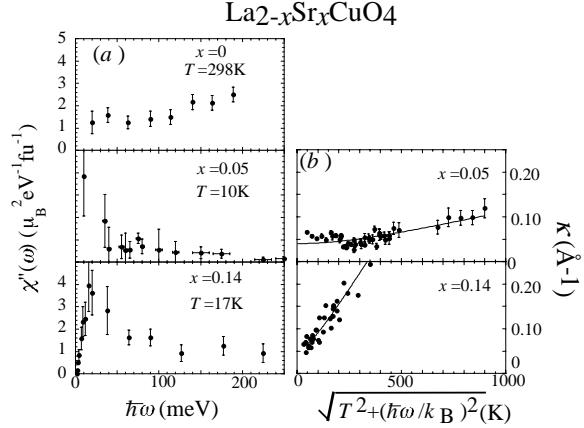


Fig. 2. (a): Energy spectra of  $\chi''(\omega)$  of  $\text{La}_{2-x}\text{Sr}_x\text{CuO}_4$  nondoped LSCO (top),  $\text{La}_{1.95}\text{Sr}_{0.05}\text{CuO}_4$  (middle), and  $\text{La}_{1.86}\text{Sr}_{0.14}\text{CuO}_4$  (bottom). Note that  $\chi''(\omega)$  of  $\text{La}_{1.95}\text{Sr}_{0.05}\text{CuO}_4$  is arbitrary, and we cannot compare the absolute value of  $\chi''(\omega)$ . (b): Relation between the Lorentzian peak-width  $\kappa$  and  $\sqrt{T^2 + (\hbar\omega/k_B)^2}$  for  $\text{La}_{1.86}\text{Sr}_{0.14}\text{CuO}_4$  (lower) [6], and  $\text{La}_{1.95}\text{Sr}_{0.05}\text{CuO}_4$  (upper).

tor and the magnetic interaction evaluated by using a 2D spin-wave dispersion relation is reduced by about 20% upon hole doping with  $x=0.05$ . Second is that the  $\chi''(\omega)$  below around 50 meV increases with decreasing the energy. Also the  $\chi''(\omega)$  in this energy region show substantial temperature dependence. At present we have no clear explanation on the enhancement of  $\chi''(\omega)$  in the low energy region. However, charge stripe model is one of possible scenarios where two distinct types of magnetic excitations naturally coexist, one originates from the nondoped antiferromagnetic region with the large energy-scale and the other from the hole doped region with the small energy-scale.

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## References

- [1] K. Yamada et al., Phys. Rev. B **57** (1998) 6165.
- [2] S. M. Hayden et al., Phys. Rev. Lett. **67** (1991) 3622.
- [3] K. Yamada et al., J. Phys. Soc. Jpn. **64** (1995) 2742.
- [4] S. M. Hayden et al., Phys. Rev. Lett. **76** (1996) 1344.
- [5] Y. Endoh et al., J. Phys. Soc. Jpn. **69** (2000) Suppl. B. 16.
- [6] G. Aeppli et al., Science **278** (1997) 1432.