

# Effect of Pressure on the Superconducting Properties of $\text{Ho}_{0.5}\text{Y}_{0.5}\text{Ni}_2\text{B}_2\text{C}$

Gendo Oomi <sup>a,1</sup>, Daisuke Masaoka <sup>b</sup>, Tomoko Kagayama <sup>b</sup>, Noritaka Kuroda <sup>b</sup>,  
B.K. Cho <sup>c,2</sup>, P.C. Canfield <sup>c</sup>

<sup>a</sup> Department of Physics, Kyushu University, Ropponmatsu, Fukuoka 810-8560, Japan

<sup>b</sup> Department of Mechanical Engineering and Materials Science, Kumamoto University, Kumamoto 860-8555, Japan

<sup>c</sup> Ames Laboratory and Department Physics and Astronomy, Iowa State University, Ames, Iowa 50011, USA

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## Abstract

Electrical resistance and lattice constants of  $\text{Ho}_{0.5}\text{Y}_{0.5}\text{Ni}_2\text{B}_2\text{C}$  have been measured at high pressure in order to clarify the competition between antiferromagnetism and superconductivity. It is found that the superconducting transition temperature and the upper critical fields decrease with increasing pressure.

*Key words:* high pressure; superconductivity; borocarbide

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In the rare earth borocarbides  $\text{RNi}_2\text{B}_2\text{C}$  ( $\text{R}$ : rare earth element), the antiferromagnetism is well known to coexist with superconductivity for the materials with  $\text{R}=\text{Ho}, \text{Tm}$  and  $\text{Er}$ [1,2]. The structure of the borocarbides is a variant of the  $\text{ThCr}_2\text{Si}_2$  type, which consists of alternating layers of  $\text{RC}$  planes and  $\text{Ni}_2\text{B}_2$  slabs. Due to this fact, the physical properties are highly anisotropic. Among these compounds, the Ho-borocarbides have been studied extensively because the interplay of superconductivity and antiferromagnetism is the most significant, which is demonstrated by the large values of  $T_N/T_c$  of about 0.63. In the system of  $\text{Ho}_x\text{Y}_{1-x}\text{Ni}_2\text{B}_2\text{C}$ , this ratio can be easily changed by controlling  $x$ . In the present work, we attempted to measure the electrical resistance at high pressure and high magnetic field to make clear the interplay of antiferromagnetism and superconductivity under high pressure. The lattice constants are also measured under high pressure.

In this work we selected the sample of  $\text{Ho}_{0.5}\text{Y}_{0.5}\text{Ni}_2\text{B}_2\text{C}$  ( $x=0.5$ ), in which  $T_c$  is 12 K and  $T_N$  is 1.7 K, i.e.,

this material is near the border of the disappearance of antiferromagnetism. The values of  $T_N/T_c$  is about 0.14, which is the same as that of  $\text{TmNi}_2\text{B}_2\text{C}$ . Single crystals are grown by a  $\text{Ni}_2\text{B}$  flux method. The details of the preparation and the characterization were reported elsewhere[3]. High pressure was generated by using piston-cylinder device. The details of the present apparatus were described previously[4]. Electrical resistance was measured using standard four-probe method. The lattice constants are measured by means of X-ray diffraction under high pressure[5].

Fig.1 shows the lattice constants,  $a/a_0$ ,  $c/c_0$  and  $V/V_0$ , as a function of pressure up to 15 GPa at room temperature, where  $a_0, c_0$  and  $V_0$  are the values at ambient pressure. The structure is stable up to 15 GPa. The compression curve is found to be approximated by straight lines within experimental errors. The linear compressibilities,  $\kappa_i = -i^{-1}(\partial i/\partial P)$  ( $i=a$  or  $c$ ), are  $1.3$  and  $2.1 \times 10^{-3}$  GPa $^{-1}$  for  $a$  and  $c$  axis, respectively. The volume compressibility is  $4.7 \times 10^{-3}$  GPa $^{-1}$ , i.e., the bulk modulus is 214 GPa, which is a little bit larger than that of  $x=1$  ( $\text{HoNi}_2\text{B}_2\text{C}$ ), 192 GPa[6]. Fig.2 shows the electrical resistivity  $\rho(T)$  under high pressure.  $\rho(T)$  decreases smoothly with decreasing temperature followed by a sudden decrease at  $T_c$ . It is found

<sup>1</sup> E-mail: oomi@rc.kyushu-u.ac.jp

<sup>2</sup> present address: Department of Materials Science & Engineering, K-JIST 500-712, South Korea

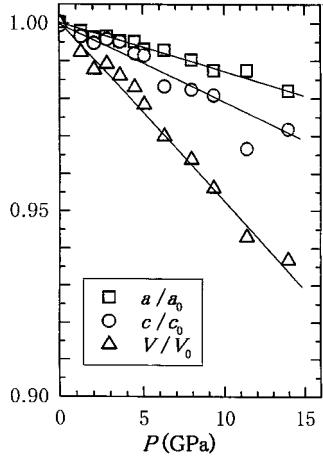


Fig. 1.  $a/a_0$ ,  $c/c_0$  and  $V/V_0$  of  $\text{Ho}_{0.5}\text{Y}_{0.5}\text{Ni}_2\text{B}_2\text{C}$  at 4.2 K as a function of pressure

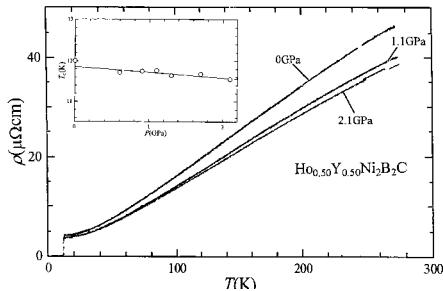


Fig. 2. Electrical resistivity  $\rho(T)$  of  $\text{Ho}_{0.5}\text{Y}_{0.5}\text{Ni}_2\text{B}_2\text{C}$  under high pressure.  $T_c$  is shown in the inset as a function of pressure.

that  $T_c$  decreases with increasing pressure: the pressure dependence is shown in the inset of Fig. 2. The pressure coefficient is  $-0.14 \text{ K/GPa}$ , which is smaller than those of  $x=1$  and  $\text{TmNi}_2\text{B}_2\text{C}$  [7,8].

Fig. 3 shows the upper critical fields  $H_{c2}$  of  $x=0.5$  as a function of pressure.  $H_{c2}$  is found to decrease almost linearly with increasing pressure having a pressure coefficient of  $(1/H_{c2})(\partial H_{c2}/\partial P) = -0.14 \text{ GPa}^{-1}$ , which is larger than that of  $\text{YNi}_2\text{B}_2\text{C}$ ,  $-0.038 \text{ GPa}^{-1}$ , but smaller than that of  $\text{HoNi}_2\text{B}_2\text{C}$ ,  $-0.70 \text{ GPa}^{-1}$  [9,10].

Here we compare the present results briefly with those of  $\text{TmNi}_2\text{B}_2\text{C}$  using Grüneisen parameters  $\Gamma$  because  $T_N/T_c$  of  $x=0.5$  is nearly the same as that of  $\text{TmNi}_2\text{B}_2\text{C}$ , 0.14.  $\Gamma$  is defined as  $\Gamma = -(\partial \ln E / \partial \ln V)$ , where  $E$  is the characteristic energy. The values of  $\Gamma$  for  $T_c$  are estimated by using the bulk modulus obtained here to be, -3 and -19 for  $x=0.5$  and  $\text{TmNi}_2\text{B}_2\text{C}$ , respectively. The detailed reason for this difference is not clear at present, but it should be noted that the coefficient of  $T^2$  term in  $\rho(T)$  of  $\text{TmNi}_2\text{B}_2\text{C}$  above  $T_c$  is 3 times larger than that of  $x=0.5$ . This fact implies that the effect of spin fluctuation is more significant

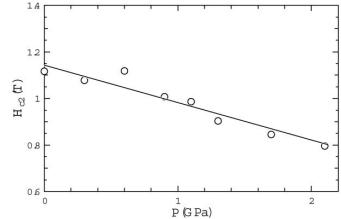


Fig. 3. The upper critical fields  $H_{c2}$  of  $\text{Ho}_{0.5}\text{Y}_{0.5}\text{Ni}_2\text{B}_2\text{C}$  at 4.2 K as a function of pressure.

in  $\text{TmNi}_2\text{B}_2\text{C}$  than in  $x=0.5$  and this may give rise to a large value of  $\Gamma$ . Furthermore we also estimate  $\Gamma$  for  $H_{c2}$ . The values are 31 and 38 for  $x=0.5$  and  $\text{TmNi}_2\text{B}_2\text{C}$ , respectively. These are comparable each other, which is largely different from the case of  $T_c$ . This is due to the fact that the antiferromagnetic ordering temperature of both sample is around 2 K, i.e., we observe  $H_{c2}$  near  $T_N$ . Near 4.2 K, the antiferromagnetism competes with the superconductivity becomes to give large value of  $\Gamma$ .

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