

# Effect of Pressure on the Electrical Resistivity of Heavy Fermion Antiferromagnet $\text{Ce}_2\text{RhIn}_8$

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## Abstract

Heavy-fermion antiferromagnet  $\text{Ce}_2\text{RhIn}_8$  is structurally related to the pressure-induced superconductor  $\text{CeRhIn}_5$ . We have measured pressure dependence of the electrical resistivity up to 1.6 GPa in the temperature range 1.5-300 K for  $\text{Ce}_2\text{RhIn}_8$ . We have found that the Néel temperature  $T_N$  decreases by pressure at the rate of  $-0.8 \text{ K GPa}^{-1}$ . This large decrease of  $T_N$  is a contrast to the case of  $\text{CeRhIn}_5$  in which  $T_N$  has very weak pressure dependence.

*Key words:* heavy fermion, pressure effect, electrical resistivity,  $\text{Ce}_2\text{RhIn}_8$

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## 1. Introduction

In recent years, heavy-fermion antiferromagnet  $\text{Ce}_n\text{RhIn}_{3n+2}$  ( $n=1, 2$  or  $\infty$ ) have attracted much attention, since pressure-induced superconducting transition were discovered for cubic  $\text{CeIn}_3$  ( $n=\infty$ ) and tetragonal  $\text{CeRhIn}_5$  ( $n=1$ ). At ambient pressure,  $\text{CeIn}_3$  and  $\text{CeRhIn}_5$  are antiferromagnet below  $T_N=10$  and 3.8 K, respectively.[1,2] Applying pressure, both exhibit superconducting transition with  $T_C=0.25$  K at  $P=2.5$  GPa for  $\text{CeIn}_3$  and  $T_C=2.0$  K at  $P=1.6$  GPa for  $\text{CeRhIn}_5$ . [2,3]  $\text{Ce}_n\text{RhIn}_{3n+2}$  can be viewed  $n$  layers of  $\text{CeIn}_3$  are stacked separated by a layer of  $\text{RhIn}_2$ . [4] Tetragonal  $\text{Ce}_2\text{RhIn}_8$ , which is a  $n=2$  member, has an intermediate crystal structure between  $\text{CeIn}_3$  ( $n=\infty$ ) and  $\text{CeRhIn}_5$  ( $n=1$ ).  $\text{Ce}_2\text{RhIn}_8$  orders at  $T_N=2.8$  K and the magnetic structures are the same as those in  $\text{CeIn}_3$ . [5] For further understanding magnetic and electronic properties of  $\text{Ce}_2\text{RhIn}_8$ , we have grown the single crystals of  $\text{Ce}_2\text{RhIn}_8$  and carried out the resistivity measurements under hydrostatic pressure.

## 2. Experimental

Single crystals of  $\text{Ce}_2\text{RhIn}_8$  were grown using a flux technique described elsewhere.[6] Powder X-ray diffraction spectra, which obtained using the crushed single crystals, were well indexed by the tetragonal  $\text{Ho}_2\text{CoGa}_8$ -type structure with lattice parameters  $a=4.664 \text{ \AA}$  and  $c=12.25 \text{ \AA}$ . [4] The crystal orientation was determined by a usual Laue method. The typical sizes of grown crystals were 3 mm along the  $ab$ -plane and 0.5 mm along the  $c$ -axis. Measurements of the electrical resistivity along the  $a$ -axis were carried out using a four-probe DC method under hydrostatic pressure in a clamp-type pressure cell.[6] The temperature dependence of resistivity were measured under fixed pressures up to 1.6 GPa in the range 1.5-300 K.

## 3. Results and Discussion

Figure 1 illustrates the resistivity along the  $a$ -axis for  $\text{Ce}_2\text{RhIn}_8$  as a function of logarithmic temperature  $\log T$  under the hydrostatic pressure  $P=0, 1.0$  and 1.6 GPa. In any pressure, the resistivity shows  $-\log T$  dependence below about 100 K, and exhibit a peak at  $T_{max}$ . As shown in the inset of Figure 1,  $T_{max}$  increases linearly

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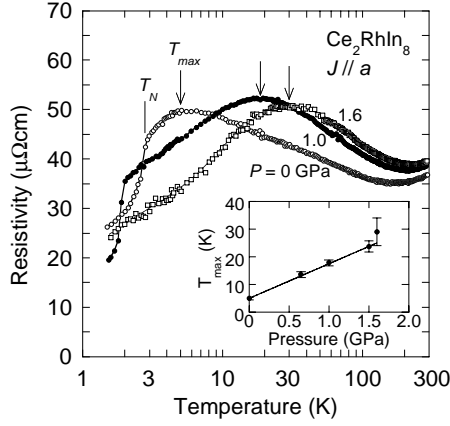


Fig. 1. The resistivity as a function of  $\log T$  for  $\text{Ce}_2\text{RhIn}_8$  along the  $a$ -axis at  $P = 0$  (open circle) and 1.6 GPa (open square). The arrows indicate the resistivity maximum temperature  $T_{max}$ . The inset shows pressure dependence of  $T_{max}$ .

by pressure with the rate  $+12.5 \text{ K GPa}^{-1}$ . Since  $T_{max}$  is roughly proportional to Kondo temperature  $T_K$ , the pressure dependence of  $T_K$  may be inferred in the inset of figure 1. With increasing the pressure, the resistivity increased linearly at high temperature region and the negative logarithmic slope becomes steeper. These results are typical for Ce-based heavy-fermion compounds.

At ambient pressure ( $P=0$ ), the resistivity exhibits a large drop at the Néel temperature  $T_N = 2.8 \text{ K}$  and gradually reaches residual resistivity about  $26 \mu\Omega\text{cm}$  with decreasing temperature. The details of the low temperature variations of resistivity under several fixed pressures are shown in Figure 2. Applying the pressure,  $T_N$  decreases as indicated by arrows in Figure 2. At 1.6 GPa, magnetic order is not observed above 1.5 K. As shown in Figure 3,  $T_N$  decreases linearly up to 1.5 GPa at the rate of  $-0.8 \text{ K GPa}^{-1}$ . This large decrease is contrast to the case of  $\text{CeRhIn}_5$ , in which  $T_N$  slightly increases with pressure.

As clearly seen in Figure 2, the low temperature resistivity is largely reduced by pressure. We could not determine the residual resistivity above 1.5 GPa, thus we plot the resistivity at 3 K,  $\rho(3 \text{ K})$ , versus pressure in Figure 3. The  $\rho(3 \text{ K})$  decreases linearly with pressure up to 1.5 GPa, and show a drop at 1.6 GPa.

Very recently pressure-induced superconducting transition and  $P$ - $T$  phase diagram for  $\text{Ce}_2\text{RhIn}_8$  were reported in ref. 7. Above 1.6 GPa, magnetic order is disappeared and only a superconductivity phase exists. [7] The decreasing rate of  $T_N$  with pressure is good agreement with our data. The sharp bend in the  $\rho(3 \text{ K})$  -  $P$  line at 1.6 GPa in Figure 3 may be related to the disappearance of magnetic order at this pressure.

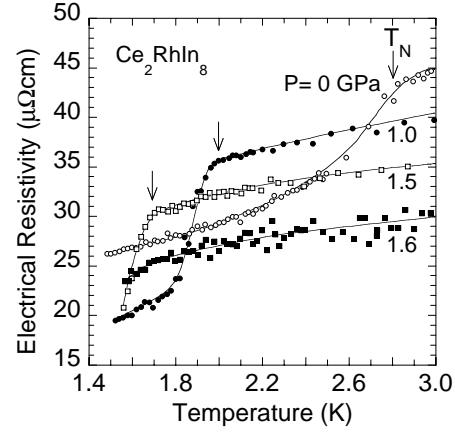


Fig. 2. The low temperature resistivity for  $\text{Ce}_2\text{RhIn}_8$  along the  $a$ -axis at several constant pressures  $P=0, 1.0, 1.5$  and 1.6 GPa. The arrows indicate the  $T_N$ .

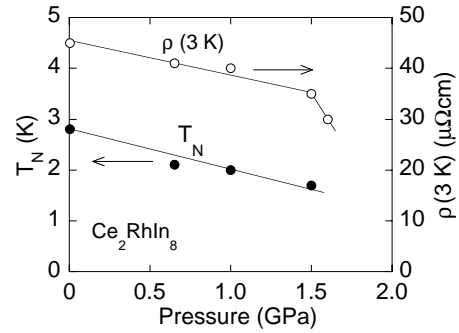


Fig. 3. The pressure dependences of the  $T_N$  and the resistivity at 3 K,  $\rho(3 \text{ K})$ , for  $\text{Ce}_2\text{RhIn}_8$ .

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