

# NMR studies of the electron-doped hafnium nitride superconductor

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## Abstract

We report on Nuclear Magnetic Resonance measurements on an oriented polycrystalline sample of the layered superconductor  $\text{Li}_{0.48}(\text{THF})_{0.3}\text{HfNCl}$  having  $T_c \sim 26$  K.  $^{35}\text{Cl}$ -NMR signals were observed around zero Knight shift, suggesting that the partial Fermi-level density of states,  $N^{Cl}(E_F)$ , at Cl site is practically nothing, and the superconductivity is derived from the [HfN] double-honeycomb network. These results reconfirm that  $\text{Li}_{0.48}(\text{THF})_{0.3}\text{HfNCl}$  is a quasi-two-dimensional superconductor.

**Key words:** Layered superconductor ;HfNCl ;Nuclear Magnetic Resonance ; Low carrier system

In 1998, Yamanaka and his coworkers discovered a new type of superconductor, Li-doped  $\beta$ -HfNCl [1]. This material has attracted a great deal of attention because of the variety of its physical property. The relatively high transition temperature of  $T_c = 25.5$  K is realized by a small amount of Li-intercalation to the layered insulator HfNCl with a band gap of  $\sim 4$  eV. On intercalation, the interlayer distance  $d$  increases from  $9.23$  Å of  $\beta$ -HfNCl to  $18.7$  Å of  $\text{Li}_{0.48}(\text{THF})_{0.3}\text{HfNCl}$ , as schematically shown in Fig.1, and electrons are believed to be doped into the double HfN layer.

The bulk superconductivity suddenly appears at  $T_c \sim 25.5$  K for the doping contents of  $x \sim 0.13$ .  $T_c$  is almost constant ( $\sim 25.5$  K) up to  $x \sim 0.5$  but gradually decreases to  $\sim 24.4$  K toward  $x \sim 1$  [1]. A question why such a high  $T_c$  is realized in the vicinity of insulating phase naturally arises.

Uemura *et. al.* pointed out the exotic nature of superconducting properties near the Bose-Einstein condensation limit from muon spin relaxation ( $\mu\text{SR}$ ) measurements [6]. Recently, we reported the quasi-two-dimensional superconducting character in this

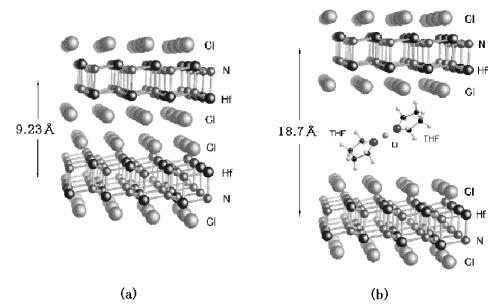


Fig. 1. Schematic structural model of (a) pristine HfNCl and (b)  $\text{Li}_{0.48}(\text{THF})_{0.3}\text{HfNCl}$ .

system from dc-magnetization and H-, Li-NMR measurements and clarified that the superconductivity is derived from the HfNCl layer [7,8]. Furthermore, we reported that the small Fermi level density of states ( $\sim 0.25$  states/eV) fail to explain the origin of the high  $T_c$  in terms of the conventional BCS model[9]. Anyway, whether or not the Cl block layers are related to the occurrence of superconductivity remain to be

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clarified. In this report, we present Cl-NMR results in the oriented  $\text{Li}_{0.48}(\text{THF})_{0.3}\text{HfNCl}$ .

Detailed sample preparation and experimental procedure were reported previously [1,8,9]. Magnetization measurements were performed up to  $H = 15\text{T}$  to check the superconducting transition temperature.

Figure 2 shows the  $T$  dependence of  $^{35}\text{Cl}$ -NMR spectra of  $\text{Li}_{0.48}(\text{THF})_{0.3}\text{HfNCl}$  measured in the magnetic field  $H \sim 9.4\text{T}$ . Note that the  $^{35}\text{Cl}$ -NMR spectra were observed around zero Knight shift. Here, the origin of the Knight shift is referred to the center-of-mass position of the spectrum of the pristine HfNCl, which is the isotropic chemical shift,  $\sim 120\text{ ppm}$ , with respect to spectrum of standard 1M-LiCl. The linewidth of  $\sim 600\text{ ppm}$  is comparable to the chemical shift range,  $0 \sim 1000\text{ ppm}$ , in insulating covalent compounds such as  $\text{CCl}_4$ . Thus, it is reasonable to consider that the  $^{35}\text{Cl}$ -NMR shift is due to the chemical shift

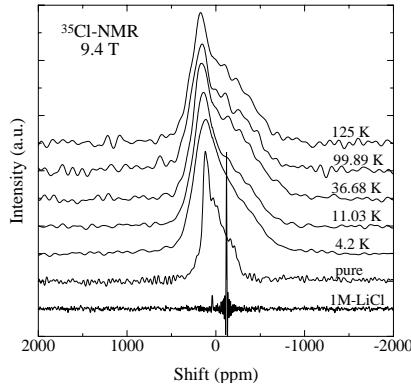


Fig. 2. Temperature dependence of  $^{37}\text{Cl}$ -NMR spectra.

Figure 3 shows the  $T$  dependence of  $^{37}\text{Cl}$  NMR shift. Below  $T_c = 20.5\text{ K}$ , the decrease of  $^{35}\text{Cl}$  NMR shift below  $T_c$ ,  $\sim -8\text{ ppm}$  ( $\sim 0.6\text{ Oe}$ ), is explained by the superconducting diamagnetic contribution in the vortex lattice, as discussed previously [7,8]; superconducting diamagnetic shift  $H_{dia}$  is estimated to be  $0.4\text{ Oe}$  at  $H = 94\text{ kOe}$  by using the relation  $H_{dia} = H_{c1}\ln(0.381e^{-0.5}d/\xi)/\ln\kappa$  [10] for  $\kappa = \sqrt{\kappa_{ab}\kappa_{||c}} \sim 151$ ,  $\xi = \sqrt{\xi_{ab}\xi_c} \sim 307$ ,  $H_{c1}^{ab} = 9\text{ Oe}$ ,  $d = 160\text{ \AA}$  which is the nearest neighbor vortex lattice spacing at  $94\text{ kOe}$ . These results indicate that the partial Fermi level density of states,  $N^{Cl}(E_F)$ , at Cl site is considerably small. Together with the previous reports [8],  $[\text{HfN}]_2$ -honeycomb network plays a major role in occurrence of the superconductivity in this system. These results are consistent with the results from band calculations; the conduction band has a two-dimensional (2D) character originating in planer hafnium  $d_{xy}$  and  $d_{x^2-y^2}$  hybridized with nitrogen  $p_x, p_y$  [3-5].

In summary, we measured  $^{35}\text{Cl}$ -NMR in the  $T$  range of  $4 \sim 150\text{K}$  across  $T_c$  for  $\text{Li}_{0.48}(\text{THF})_{0.3}\text{HfNCl}$ .

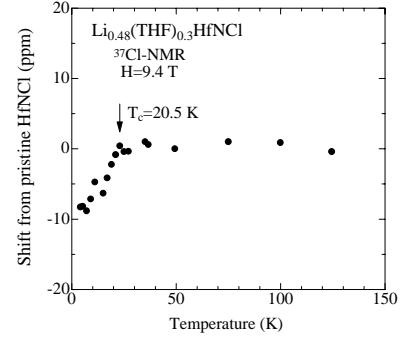


Fig. 3. Temperature dependence of  $^{37}\text{Cl}$ -NMR shift.

Present studies demonstrate that the superconductivity in this system is derived from the two-dimensional  $[\text{HfN}]_2$  honeycomb-network and consistent with the previous results [7,8]. In order to clarify the question why such a high  $T_c$  is realized in the vicinity of insulating phase,  $^{15}\text{N}$ -NMR measurements are now in progress and details will be reported elsewhere.

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