

# Comparative Study on the Anisotropic Properties of MgB<sub>2</sub>

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## Abstract

The anisotropy of upper critical field of MgB<sub>2</sub> has been studied on single crystals and poly-crystals by both transport and magnetic measurements. The angular dependence of  $H_{c2}$  shows deviation from the anisotropic Ginzburg-Landau model at lower temperatures. The value of anisotropy parameter is temperature and sample dependent, and is about 3 and 4.5 for single and poly crystals, respectively, at temperatures near  $T_c$ . However,  $H_{c2}^c(T)$  are almost the same for all samples. These features could be an indication of anisotropic s-wave superconductivity with pancake-like energy gap or resulted from the different impurity levels in these samples. The anisotropy of  $H_{c2}$  in *ab*-plane has also been measured and we set an upper bound of 1% for the in-plane anisotropy.

*Key words:* MgB<sub>2</sub> ;anisotropy ;upper critical field ;

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## 1. Introduction

MgB<sub>2</sub> has stimulated intense researches all over the world for its highest  $T_c$  among intermetallic compounds and simple structure. Its superconducting mechanism, two gap structure, anisotropic properties, surface superconductivity are issues of recent interests. Among these issues, the anisotropy of upper critical field is very important for both superconducting mechanism and applications of MgB<sub>2</sub>. However, the anisotropy of MgB<sub>2</sub> is still under debate due to the span of anisotropy parameter  $\gamma = H_{c2}^{ab}/H_{c2}^c$  from 1.1 to 13 [1]. It is not very clear why  $\gamma$  depends on temperature and samples [2]. It is also not clear whether the anisotropic behavior of MgB<sub>2</sub> obeys the anisotropic Ginzburg-Landau relation or not. In this paper, we report our comparative studies on  $\gamma$  for single crystals and poly-crystals by both transport and magnetic measurements. We also

report in-plane anisotropy of  $H_{c2}$  by carefully aligning the magnetic field in the plane.

## 2. Experimental

The single crystals of MgB<sub>2</sub> are synthesized under high pressure and its crystallinity has been checked by Laue x-ray photograph. The Laue x-ray diffraction shows a clear six-fold pattern. The superconducting transition temperature ( $T_c$ ) of single crystalline samples are 36-38 K and the transition width ( $\Delta T_c$ ) is smaller than 0.3 K. The single crystalline samples are thin platlets with thickness of about 25  $\mu m$ . The dense poly-crystalline samples are prepared by directly reacting Mg and B in Ta tube without pressure.  $T_c$  is about 38.5 K and the transition is very sharp (0.2 K). The upper critical fields are determined by transport measurements using standard four-probe method and by SQUID magnetometer. For transport measurements of the anisotropy of  $H_{c2}$ , we use a two-axis sample rotator and a vector magnet system with maximum fields

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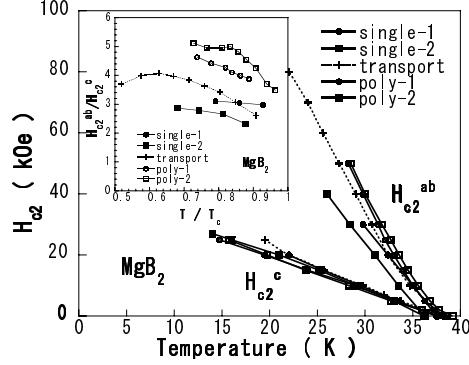


Fig. 1. Temperature dependence of the upper critical field  $H_{c2}^{ab}$  and  $H_{c2}^c$ . Inset shows the temperature dependence of anisotropy parameter  $\gamma = H_{c2}^{ab}/H_{c2}^c$  for different samples.

of 50 kOe and 30 kOe for transverse and longitudinal directions, respectively.

### 3. Results and Discussion

Figure 1 shows the temperature dependence of upper critical fields obtained from  $R - T$  and  $M - T$  curves. For single crystalline sample,  $H_{c2}$  are determined by the onset of diamagnetic transition or the kink points of  $R - T$  curves [3]. For poly-crystalline samples, the  $H_{c2}^c$  and  $H_{c2}^{ab}$  are estimated by the kink of the slope of  $M - T$  curves proposed by Bud'ko *et al.* [4]. The inset of Fig. 1 shows the temperature dependence of the anisotropy parameter  $\gamma = H_{c2}^{ab}/H_{c2}^c$ .  $H_{c2}^c$  and  $H_{c2}^{ab}$  have different temperature dependence and hence  $\gamma$  is also temperature-dependent, which implies a breakdown of the anisotropy of the band effective mass or may be related to the anisotropy of the energy gap structure of  $\text{MgB}_2$ .  $H_{c2}^{ab}(T)$  is obviously sample-dependent while  $H_{c2}^c(T)$  is almost independent of samples. These features may be resulted from two origins: (1) The pancake-like energy gap anisotropy proposed by Posazhennikova *et al.* [5], which may result in larger change of  $H_{c2}^{ab}$  while smaller change of  $H_{c2}^c$  for samples with different  $\gamma$  values and can result in temperature-dependent  $\gamma$ . (2) Difference in impurity levels, which can also result in this phenomenon because  $H_{c2}^c$  only depends on  $\xi_{ab}$  while  $H_{c2}^{ab}$  depends on both  $\xi_{ab}$  and  $\xi_c$ . So if only  $\xi_c$  is different for different sample,  $H_{c2}^{ab}$  and  $\gamma$  will change and  $H_{c2}^c$  will not change for samples with different impurity levels.

To check the anisotropic behavior of  $H_{c2}$ , we have measured the angular dependence,  $H_{c2}(\theta)$ , at different temperatures as shown in Fig. 2. Here  $\theta$  is the angle of the field from the c-axis. At temperatures near  $T_c$ , the  $H_{c2}(\theta)$  can be fitted by anisotropic GL relation, while at lower temperatures, it deviates from the GL theory.

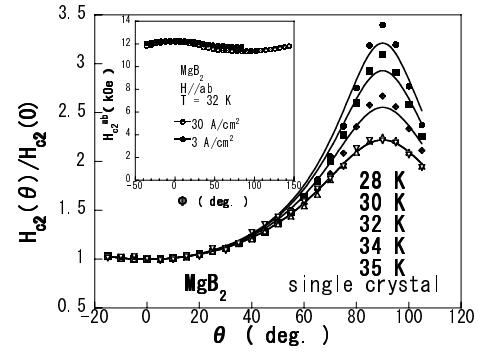


Fig. 2. Angle dependence of normalized upper critical field  $H_{c2}(\theta)$  at different temperatures determined by the peak of the derivative of  $R - H$  curves. Solid lines are the best fitting by GL model with  $\gamma = 3.21, 2.93, 2.55, 2.22, 2.22$ . Inset shows the angle dependence of  $H_{c2}^{ab}$  in ab-plane at  $32$  K determined by  $R - H$  curves with different transporting currents.

The peak at  $90^\circ$  is sharper than GL theory, which may be resulted from the effect of special energy gap structure of  $\text{MgB}_2$  or just the effect of surface superconductivity on the determination of  $H_{c2}$  [6].

We also measured the in-plane anisotropy of the upper critical field,  $H_{c2}(\phi)$ , by carefully aligning the magnetic field exactly in the  $ab$ -plane for each  $\phi$ . The current is passed along the  $a$ -axis, where we define the direction of the in-plane field  $\phi = 0^\circ$ . The inset of Fig. 2 shows  $H_{c2}(\phi)$  at  $T = 32\text{K}$  for two different current densities. As is evident from the figure, the dominant component of  $H_{c2}(\phi)$  has a two-fold symmetry rather than the six-fold symmetry expected from the hexagonal crystal structure. The two-fold symmetry is due to the Lorentz force as evidenced by the current density dependence and the minimum of  $H_{c2}(\phi)$  for a field perpendicular to the current. After subtracting the two-fold symmetry component, the six fold-symmetry component is less than 1%. It should be noted that a clear anisotropy of  $H_{c2}(\phi)$  more than 30% in hexagonal material  $\text{Cs}_x\text{WO}_3$  is reported in high quality single crystals [7].

This work is supported by Grant-Aid for Scientific Research from Ministry of Education, Culture, Sports, Science and Technology.

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