

Charge and orbital ordered states in $\text{Nd}_{1-x}\text{Sr}_{1+x}\text{MnO}_4$ ($x = 0.67, 0.75$)

Hironori Nakao^{a,1}, Tadashi Satoh^a, Junko Satoh^a, Youichi Murakami^{a,b},
Masato Kubota^{c,2}, Yusuke Wakabayashi^c, Hiroshi Sawa^c, Tsuyoshi Kimura^d,
Yoshinori Tokura^d

^a*Department of Physics, Tohoku University, Sendai 980-8578, Japan*

^b*Synchrotron Radiation Research Center, JAERI, Hyogo 679-5143*

^c*Photon Factory, Institute of Materials Structure Science, KEK, Tsukuba, 305-0801, Japan*

^d*Department of Applied Physics, University of Tokyo, Tokyo 113-8656, Japan*

Abstract

The charge and orbital ordered states in the single-layered manganite crystals, $\text{Nd}_{1-x}\text{Sr}_{1+x}\text{MnO}_4$ ($x = 0.67, 0.75$), have been investigated. We have observed the superlattice spots with modulation wave vector, $q = (\delta, \delta, 0)$, $\delta = \frac{1}{8}, \frac{1}{4}$ [$\sim \frac{1}{6}, \frac{1}{3}$] in $x = 0.75[0.67]$ by resonant and non-resonant x-ray scattering. These results show that not only the atomic displacement pattern but also the charge-orbital structure has the unit cell of $4\sqrt{2}a \times \sqrt{2}a \times c$ [$3\sqrt{2}a \times \sqrt{2}a \times c$].

Key words: charge ordering; orbital ordering; resonant x-ray scattering;

1. Introduction

In the electric and magnetic properties of manganites, the charge, orbital and spin degrees of freedom each play important roles. The typical system, a layered-perovskite $\text{La}_{1-x}\text{Sr}_{1+x}\text{MnO}_4$, have intensively been studied so far [1]. For example, in $\text{La}_{0.5}\text{Sr}_{1.5}\text{MnO}_4$, the charge and orbital ordering affects the transport and the magnetic structure. These charge, orbital and spin ordered states were confirmed by the resonant x-ray scattering [2] and the neutron scattering [3]. However, the ordered states with $x > 0.5$ have seldom been investigated owing to the chemical phase separation [1]. By substituting Nd for La, Kimura *et al.* have recently succeeded in the growing the single crystals, $\text{Nd}_{1-x}\text{Sr}_{1+x}\text{MnO}_4$ ($0.67 \leq x \leq 1.0$) [4]. By an electron diffraction, they have found the superstructures for $x = 0.67, 0.75$ concerning with a

steep rise of the resistivity. The superlattices observed here indicate only the atomic displacements which may induced by the charge-orbital ordering. To confirm the charge-orbital ordered states, we have studied the $\text{Nd}_{1-x}\text{Sr}_{1+x}\text{MnO}_4$ ($x = 0.67, 0.75$) using resonant and non-resonant x-ray scattering. We discuss the charge-orbital ordered state by comparing with Wigner-crystal and bi-stripe models in this paper.

2. Experimental results and discussions

X-ray scattering experiments were carried out at the beam lines 4C of the Photon Factory in KEK. The incident X-ray energy is about 6.55 keV of Mn *K*-edge and the energy resolution is about 2 eV. A four-circle diffractometer equipped with a closed cycle He cryostat was used. At room temperature, the crystal structure is space group $I4/mmm$ ($a = 3.82$, $c = 12.5$ Å for $x = 0.75$). All the indices are denoted in the $I4/mmm$ setting in this paper.

¹ E-mail: nakao@iyo.phys.tohoku.ac.jp

² Present address: ERATO-JST, Tsukuba, Ibaraki 305-8562

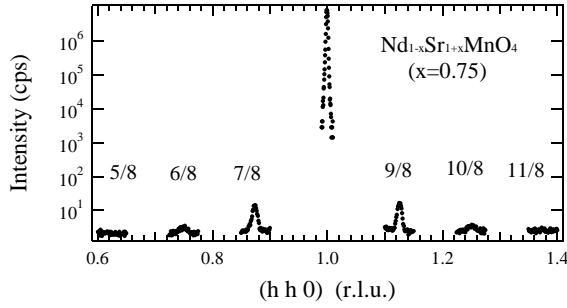


Fig. 1. $(h h 0)$ scan at resonant energy for $x = 0.75$ at $T = 8$ K.

We have utilized resonant and non-resonant x-ray scattering. The resonant x-ray scattering is sensitive to charge-orbital ordered state, while the non-resonant x-ray scattering is caused by the atomic displacement. To determine the charge-orbital ordered states, we have searched the superlattice spots of the low temperature phase at Mn K -edge, $E = 6.555$ keV (a resonant energy). The scan along $(h h 0)$ for $x = 0.75$ are presented in Fig. 1. Not only the modulation wave vector, $q = (\delta, \delta, 0)$, $\delta = \frac{1}{8}$, which was reported by the study of an electron diffraction [4], but also the weak signal at $\delta = \frac{1}{4}$ were observed. These signals observed at resonant energy may include non-resonant signals. The en-

ergy dependence at $(\frac{5}{4}, \frac{5}{4}, 0)$ shows the interference process, involving the resonant and non resonant scattering. This non-resonant signal may be attributed to the displacement of the oxygen atoms. In the superlattice spots $(2 \pm \delta, \pm \delta, 0)$ around the $(2 0 0)$ reflection, the strong intensities of non-resonant scattering have been observed. We have observed no resonant intensity owing to the strong non-resonant scattering.

The resonant and non-resonant scattering with $\delta = \frac{1}{8}, \frac{1}{4}$ indicates that the unit cell of the charge-orbital and atomic displacement pattern is $4\sqrt{2}a \times \sqrt{2}a \times c$ in the low temperature phase. The experimental results are examined by Wigner-crystal [6] and bi-stripe models [7]. However, these models can not explain the intensity ratio among the superlattice intensities. It would rather indicate the sinusoidal charge-orbital distribution along a -axis, which was discussed in ref. [8]

The same measurements for $x = 0.67$ have also been performed. The modulation wave vectors with $\delta \sim \frac{1}{6}, \frac{1}{3}$ have been observed by resonant and non-resonant scattering. We note that the superlattice intensity is much weaker than that for $x = 0.75$. The reason is still an open question.

In conclusion, we have studied the low temperature phase in $\text{Nd}_{1-x}\text{Sr}_{1+x}\text{MnO}_4$ ($x = 0.67, 0.75$) using resonant and non-resonant x-ray scattering technique. It has been elucidated that not only the atomic displacement but also the charge-orbital ordering have the superstructure of $\delta = \frac{1}{8}, \frac{1}{4}$ for $x = 0.75$ and $\delta = \frac{1}{6}, \frac{1}{3}$ for $x = 0.67$.

Acknowledgements

This work was supported by a Grant-In-Aid for Scientific Research from the Ministry of Education, Science and Culture, Japan, and Core Research for Evolutional Science and Technology. This study has been performed under the approval of the Photon Factory Program Advisory Committee (Proposal No. 2001S2-002).

References

- [1] Y. Moritomo *et al.*, Phys. Rev. B **51** (1995) 3297; W. Bao *et al.*, Solid State Commun. **98** (1996) 55.
- [2] Y. Murakami *et al.*, Phys. Rev. Lett. **80** (1998) 1932; Y. Wakabayashi *et al.*, J. Phys. Soc. Jpn. **70** (2001) 1194.
- [3] B. J. Sternlieb *et al.*, Phys. Rev. Lett. **76** (1996) 2169.
- [4] T. Kimura *et al.*, Phys. Rev. B **65** (2001) 020407.
- [5] Y. Murakami *et al.*, Phys. Rev. Lett. **81** (1998) 582.
- [6] P. G. Radaelli *et al.*, Phys. Rev. B **59** (1999) 14440.
- [7] Mori *et al.*, Nature **392** (1998) 473.
- [8] T. Nagai *et al.*, Phys. Rev. B **65** (2002) 060405.

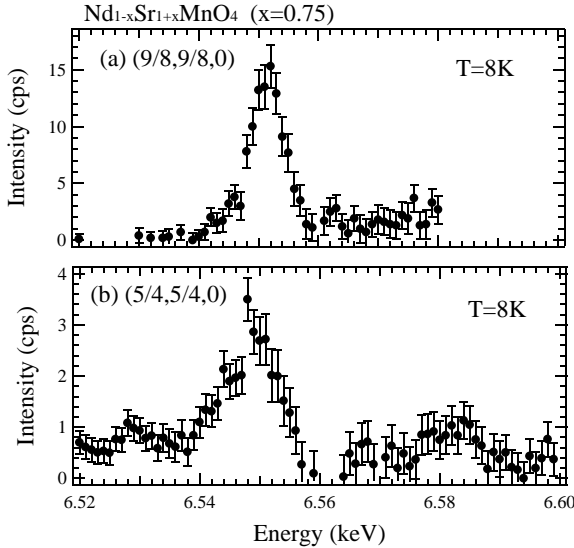


Fig. 2. Energy dependence of scattering intensities at (a) $(\frac{9}{8}, \frac{9}{8}, 0)$ and (b) $(\frac{5}{4}, \frac{5}{4}, 0)$.

ergy dependence of the scattering intensities has been measured at $(\frac{9}{8}, \frac{9}{8}, 0)$ and $(\frac{5}{4}, \frac{5}{4}, 0)$ so that the resonant signal can be discriminated from the non-resonant one. The energy spectrum at $(\frac{9}{8}, \frac{9}{8}, 0)$ shows a clear resonant feature, which is similar to the LaMnO_3 case [5], without the non-resonant signal. On the other hand,