

Transport properties of the dense Kondo system $\text{Ce}_{0.5}\text{La}_{0.5}\text{B}_6$

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Abstract

We have measured the electrical resistivity of the dense Kondo system $\text{Ce}_{0.5}\text{La}_{0.5}\text{B}_6$ at very low temperatures. A large term which is in proportion to temperature squared is found in the resistivity of $\text{Ce}_{0.5}\text{La}_{0.5}\text{B}_6$ below ~ 0.3 K. The proportion coefficient of this term of $\text{Ce}_{0.5}\text{La}_{0.5}\text{B}_6$ is 37 times larger than that of CeB_6 .

Key words: dense Kondo effect; $\text{Ce}_{0.5}\text{La}_{0.5}\text{B}_6$; electrical resistivity ;magnetoresistivity

The magnetic phases of $\text{Ce}_x\text{La}_{1-x}\text{B}_6$ have attracted much attention because of their unusual profiles originating from the dense Kondo effect and intersite multipole interactions. Phase I is a paramagnetic phase, phase II an antiferro quadrupole (AFQ) and phase III an antiferromagnetic with AFQ ordering. Recently, additional phase IV was found in the substitutional alloys $\text{Ce}_x\text{La}_{1-x}\text{B}_6$ ($x = 0.75, 0.70, 0.65$). No Bragg peak is found in phase IV in the neutron diffraction measurements [1], while the clear λ -shaped peak is found in specific heat at the phase I-phase IV transition temperature $T_c \sim 1$ K [2]. In the case of $\text{Ce}_{0.5}\text{La}_{0.5}\text{B}_6$, only a broad maximum around 1 K is shown in the specific heat in zero field, and the specific heat of this system is in proportion to $T^{1.5}$ at lower temperatures. The order parameter of phase IV and the ground state of $\text{Ce}_{0.5}\text{La}_{0.5}\text{B}_6$ are both current controversy [3]. Some previous studies have been made on the resistivity of $\text{Ce}_{0.5}\text{La}_{0.5}\text{B}_6$ [4,5]. In the present study, we have measured the resistivity of this system extending the temperature range down to ~ 50 mK.

The sample was grown by the floating zone method and cut to a rectangular shape $0.3 \times 0.7 \times 8$ mm³. We measured the electrical resistivity by the four-terminal DC method using a top loading type dilution refriger-

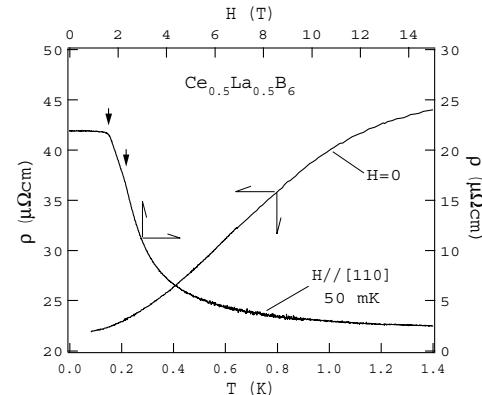


Fig. 1. Temperature dependence of the resistivity of $\text{Ce}_{0.5}\text{La}_{0.5}\text{B}_6$ for $H=0$ and magnetoresistivity for $T=50$ mK. Fields and the current are directed along the [110] and $[\bar{1}\bar{1}0]$ axes, respectively.

ator. The current was $0.15 \sim 0.6$ mA. For a good thermal contact to $^3\text{He}-^4\text{He}$ mixture, we glued the sample onto a copper plate $10 \times 15 \times 1$ mm³ using GE7031 grease and immersed them into the mixture directly.

The electrical resistivity at 0 T and the magnetoresistivity at 50 mK are shown in Fig. 1. Magnetic fields and the current are directed along the [110] and $[\bar{1}\bar{1}0]$ axes, respectively. In the case of $\text{Ce}_{0.75}\text{La}_{0.25}\text{B}_6$, rapid

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decrease in the resistivity was reported in association with the transition from phase I into phase IV [5]. However, we found no indication of the phase I-phase IV transition in the resistivity for $H = 0$ in $\text{Ce}_{0.5}\text{La}_{0.5}\text{B}_6$. Probably, $\text{Ce}_{0.5}\text{La}_{0.5}\text{B}_6$ is in phase I in the low T , low H region. Two anomalies indicated by bold arrows are shown in the magnetoresistivity of $\text{Ce}_{0.5}\text{La}_{0.5}\text{B}_6$. One is the bending point at 1.6 T and the other is a weak bending point at around 2.3 T. Under the magnetic fields $H//[\text{110}]$, $\text{Ce}_{0.5}\text{La}_{0.5}\text{B}_6$ may be in phase III in the range $1.6 \text{ T} < H < 2.3 \text{ T}$. The anomalies at 1.6 T and at 2.3 T may correspond to the phase I-phase III transition and the phase III-phase II transition, respectively. We illustrate the magnetic phase diagrams of $\text{Ce}_{0.5}\text{La}_{0.5}\text{B}_6$ in Fig. 2. As shown in Fig. 1, magnetoresistivity at 15 T is only 11% of that at 0 T. This implies that the large residual resistivity of $\text{Ce}_{0.5}\text{La}_{0.5}\text{B}_6$ under low fields is originated from the dense Kondo effect, *i. e.*, from the large cross section of randomly distributed Ce ions.

The temperature dependence of the resistivity at very low temperatures is shown in Fig. 3. The resistivity follows the formula $\rho(T) = \rho_0 + AT^2$ below $\sim 0.3 \text{ K}$. Here, ρ_0 is the residual resistivity and A is the proportional coefficient. The magnitude of A of $\text{Ce}_{0.5}\text{La}_{0.5}\text{B}_6$ for $H = 0$ is $30.5 \mu\Omega\text{cmK}^{-2}$ which is 37 times larger than that of CeB_6 ($0.832 \mu\Omega\text{cmK}^{-2}$) [6]. Although Ce ions are distributed randomly in $\text{Ce}_{0.5}\text{La}_{0.5}\text{B}_6$, a strongly correlated Fermi liquid state is likely realized in this system. When the long range ordering is induced by the applied field, A becomes small. The magnitude of A at 3 T is 53 % of that at 0 T. The large A under lower fields may be attributed to the non-ordered ground state of $\text{Ce}_{0.5}\text{La}_{0.5}\text{B}_6$.

In conclusion, the large term which is in proportion to temperature squared is found in the resistivity of $\text{Ce}_{0.5}\text{La}_{0.5}\text{B}_6$ in the temperature range $0.1 \lesssim T \lesssim 0.3 \text{ K}$. This strongly suggests the formation of strongly correlated Fermi liquid state in this system.

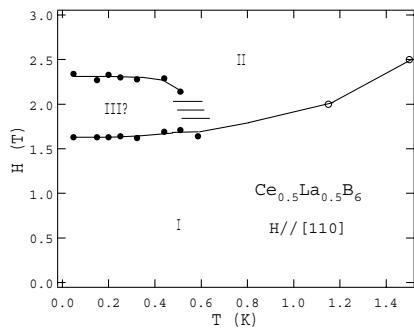


Fig. 2. Magnetic phase diagrams of $\text{Ce}_{0.5}\text{La}_{0.5}\text{B}_6$ for $H//[\text{110}]$. Closed circles are present results and open circles are taken from Ref. [7].

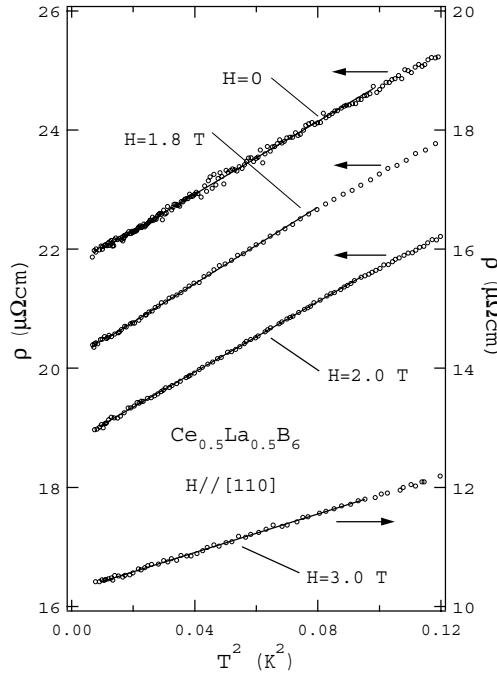


Fig. 3. Resistivity of $\text{Ce}_{0.5}\text{La}_{0.5}\text{B}_6$ at very low temperature. Fields and the current are directed along the $[\text{110}]$ and $[\bar{1}\bar{1}0]$ axes, respectively. Solid lines are fits according to the formula $\rho(T) = \rho_0 + AT^2$.

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