

Hall effect anomalies in Kondo-lattice CeAl₂

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Abstract

Anomalous Hall effect in paramagnetic and modulated antiferromagnetic (AFM) phases of a Kondo-lattice CeAl₂ has been studied in a wide range of temperatures (1.8-300K) and magnetic fields (up to 80 kOe). It was shown that the large anomalous skew scattering component R_H^a has a broad maximum around $T \approx 4$ K depressing by a factor of 3 in $H \approx 80$ kOe. A non-monotonous behavior of anomalous magnetic component R_H^m in magnetic fields below 40 kOe which also demonstrates a narrow peak of $R_H^m(T)$ at $T = T_N \approx 3.85$ K is likely to be attributed to the AFM-domains reorientation process at liquid helium temperatures. The problem of two magnetic phase transitions and complicated activation type behavior of R_H^a in this intermetallic “coexistence compound” are discussed.

Key words: spin fluctuations, Kondo-lattice, antiferromagnetism

An anomalous transport properties and especially unusual Hall effect have been observed in a number of Kondo-lattice compounds and qualitatively explained in the framework of skew scattering models [1,2]. Here the results of Hall coefficient and magnetoresistance measurements are presented for the “coexistence compound”- magnetic Kondo-lattice CeAl₂. The issue of this study was twofold – (i) to verify in details the validity of the skew scattering models [1,2] and (ii) to investigate the mysterious magnetic phase transitions [3,4] in CeAl₂ with the help of high precision transport measurements.

The experiments have been carried out on high quality polycrystalline samples of this cubic C15 Laves-phase material in wide temperature range 1.8-300K in magnetic fields up to 80 kOe. To measure the angular dependencies of Hall resistivity and magnetoresistance in CeAl₂ the sample was rotated in magnetic field \mathbf{H} around the current \mathbf{I} axis in transversal ($\mathbf{I} \perp \mathbf{H}$) geom-

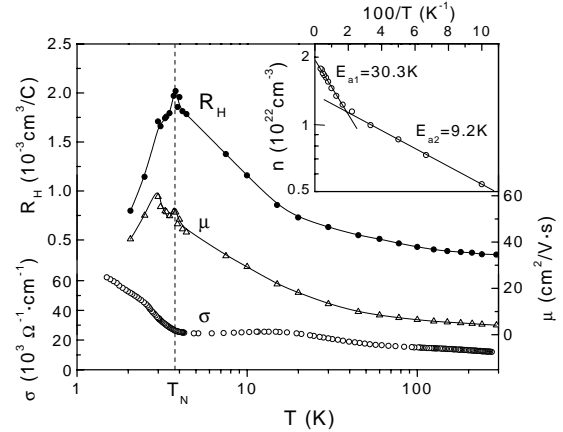


Fig. 1. Temperature dependencies of the Hall coefficient, conductivity and mobility in CeAl₂. Inset shows the activation behavior of the reciprocal Hall coefficient $n = (R_H^a e)^{-1}$.

etry. It was found [5] that a single harmonic behavior of the Hall resistivity is destroyed in magnetic fields $H \geq 3$ kOe and even harmonics appear to contribute

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in total Hall signal. To eliminate the effects of magnetoresistance contribution in the Hall resistivity the magnetoresistance measurements have been also carried out. The data analysis [5] allowed to deduce the anomalous components R_H^a and R_H^{am} of the Hall effect in CeAl₂. Temperature and magnetic field dependencies of R_H^a and R_H^{am} parameters are shown in Figs.1-2. Among these two contributions the large anomalous positive component R_H^a (skew scattering contribution [1,2]) demonstrates a broad maximum around $T \approx 4K$ (Fig.1) which is depressed drastically (by a factor of 3) in magnetic field $H \approx 80$ kOe (Fig.2). The decrease of magnitude of R_H^a in magnetic field depends only slightly from the temperature in the interval $3.4K < T_N < 4.2K$ in the vicinity of the broad maximum of R_H^a . Such a strong $R_H^a(H)$ dependence can be attributed to the depression of the Kondo-compensation mechanism in magnetic field. Indeed, in the case of CeAl₂, where the Kondo temperature is found to be $T_K \approx 5K$ [3], one can expect an essential decrease of amplitude of the Abrikosov-Suhl resonance in moderate magnetic field $H \leq 80$ kOe resulting to the reduction of the R_H^a component [1].

The most striking feature of the skew scattering component temperature dependence $R_H^a(T)$ in this intermetallic compound is a complicated **activation type behavior**. The plot $\log(R_H^a e)^{-1} = f(1/T)$ (inset of Fig.1) allows to establish two activation processes in transport with the energies $E_{a1} = 30.3 \pm 0.8K$ (in the interval 70-300K) and $E_{a2} = 9.2 \pm 0.1K$ (10-40K). Additionally, the reciprocal mobility of charge carriers $\mu^{-1}(T) = (\sigma(T) R_H(T))^{-1}$ (inset of Fig.2) is characterized by Curie-Weiss type behavior $\mu^{-1}(T) \sim \chi^{-1}(T) \sim (T + \Theta_i)$ in these intervals with $\Theta_1 = -350 \pm 20K$ and $\Theta_2 = -7.5 \pm 0.5K$ correspondingly. Among these two findings the first one is out of the conclusions of the skew scattering models [1,2], while the analytical dependence $\mu(T) \sim \chi(T)(1 - \chi(T))$ predicted in [1,2] for temperature range $T \gg T_K \approx 5K$ is very similar to that one observed in this study. However, to analyze in detail the data of Figs.1-2 one needs also to estimate the effects of real crystal electric fields in CeAl₂ ($\Delta_{CF1} \approx 100K$ [3]) additionally to approaches [1,2].

The anomalous magnetic contribution in the Hall coefficient R_H^{am} is characterized by (i) a narrow maximum at $T = T_N \approx 3.85K$ (Fig.1) and (ii) a non-monotonous behavior of R_H^{am} in the magnetic field below 40 kOe (Fig.2). A maximum of the anomalous magnetic Hall coefficient $R_H^{am}(H)$ is observed at $H_m \approx 15$ kOe, moreover, the amplitude of this peak increases dramatically when the temperature decreases below T_N (Fig.2). Following to the arguments suggested in [6] the $R_H^{am}(H)$ anomaly can likely be attributed to the AFM-domains reorientation process at liquid helium temperatures.

Another interesting feature of the $R_H^{am}(H)$ dependencies is the change of sign which occurs in the mag-

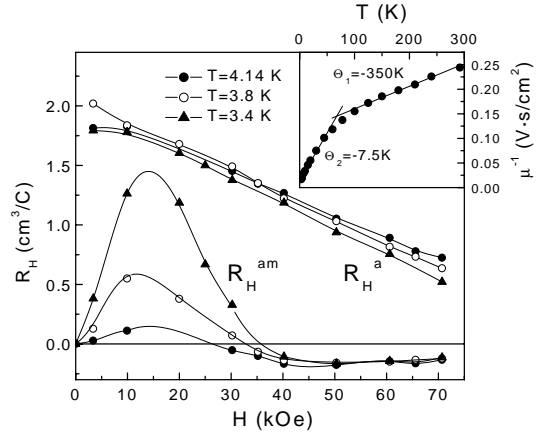


Fig. 2. Anomalous components of the Hall coefficient R_H^a and R_H^{am} versus applied magnetic field in CeAl₂. Inset shows the reciprocal mobility $\mu_H^{-1}(T)$ temperature dependence.

netic fields interval 30-40 kOe (Fig.2). It is important to stress here that the $R_H^{am}(H)$ sign inversion points are placed just above the AFM phase boundary and especially for the cases of $T = 4.14K$ and $T = 3.8K$ can be attributed to “the unknown magnetic phases” [3–6] on the $H-T$ phase diagram of CeAl₂. Thus, the last finding contributes in favor of conclusion [3,6] of a rather complicated magnetic behavior with short range ferromagnetic correlations which have been established in CeAl₂ just above the AFM phase boundary. Moreover, to support the arguments of [3] in favor of two magnetic phase transitions in CeAl₂ the mobility $\mu_H(T)$ demonstrates two narrow peaks at 3.85K and 3.0K (Fig.1). To clarify in more details the H-T magnetic phase diagram in this “coexistence” compound the anomalous transport measurements are in progress now.

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