

# Photoemission study of $\text{CeMIn}_5$ ( $M=\text{Rh, Ir}$ ) : nearly localized nature of $f$ electrons

S.-i. Fujimori <sup>a,1</sup> T. Okane <sup>a</sup> J. Okamoto <sup>a</sup> K. Mamiya <sup>a</sup> Y. Muramatsu <sup>a</sup> A. Fujimori <sup>a,b</sup>,  
T. Narimura <sup>c</sup> K. Kobayashi <sup>c</sup> K. Shimada <sup>c</sup> H. Namatame <sup>c</sup> M. Taniguchi <sup>c</sup>, H. Harima <sup>d</sup>,  
D. Aoki <sup>e</sup>, S. Ikeda <sup>e</sup> H. Shishido <sup>e</sup> Y. Tokiwa <sup>e</sup>, Y. Haga <sup>f</sup>, Y. Ōnuki <sup>e,f</sup>,

<sup>a</sup>Japan Atomic Energy Research Institute, SPring-8, Hyogo 679-5148, Japan

<sup>b</sup>Graduate School of Science, University of Tokyo, Bunkyo-ku, Tokyo 113-0033, Japan

<sup>c</sup>Hiroshima Synchrotron Radiation Center, Hiroshima University, Higashi-Hiroshima 739-8526, Japan

<sup>d</sup>Institute of Scientific and Industrial Research, Osaka University, Osaka 567-0043, Japan

<sup>e</sup>Graduate School of Science, Osaka University, Toyonaka, Osaka 560-0043, Japan

<sup>f</sup>Advanced Science Research Center, Japan Atomic Energy Research Institute, Ibaraki 319-1195, Japan

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## Abstract

We have performed  $4d$ - $4f$  ( $h\nu = 122$  eV) and  $3d$ - $4f$  ( $h\nu = 881$  eV) resonant photoemission spectroscopy studies on the heavy-fermion cerium compounds  $\text{CeMIn}_5$  ( $M=\text{Rh and Ir}$ ), which show competition between superconductivity and antiferromagnetism. The results suggest that the Ce  $4f$  electrons in both compounds are nearly localized. We have also found that although the Ce  $4f$  electrons in the superconducting  $\text{CeIrIn}_5$  are more delocalized than those in the antiferromagnetic  $\text{CeRhIn}_5$ , their electronic structures are very similar to each other.

*Key words:* photoemission spectroscopy;  $\text{CeRhIn}_5$ ;  $\text{CeIrIn}_5$

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In recent years, relationship between magnetism and superconductivity has attracted much attention. Especially, some magnetic  $f$ -electron compounds show superconductivity under high pressure, suggesting that magnetic interactions may play essential roles in the superconductivity[1,2].  $\text{CeMIn}_5$  ( $M=\text{Rh and Ir}$ ) are recently synthesized this class of Ce-based compounds. They are thought to be located near the quantum critical point (QCP) in Doniach's phase diagram, and are good target materials to study how the magnetic interactions are involved in the pairing mechanism. The  $\text{CeMIn}_5$  compounds crystallize in the tetragonal  $\text{HoCoGa}_5$ -type structure, which can be viewed as an alternating stack of the  $\text{CeIn}_3$  and  $M\text{In}_2$  layers. Therefore, they can be regarded as a quasi-two-

dimensional version of  $\text{CeIn}_3$ .  $\text{CeRhIn}_5$  is a heavy fermion (HF) antiferromagnet with  $T_N=3.8$  K at ambient pressure. According to their lattice constants, the  $\text{CeIn}_3$  layers experience a chemical pressure of  $\sim 14$  kbar relative to bulk  $\text{CeIn}_3$ . Hegger *et al.*[3] have discovered that  $\text{CeRhIn}_5$  undergoes a superconducting transition at  $P_C \sim 16.3$  kbar and  $T_C=2$  K. This  $P_C$  is actually lower than that of  $\text{CeIn}_3$ , supporting that  $\text{CeRhIn}_5$  corresponds to  $\text{CeIn}_3$  under high pressure. On the other hand,  $\text{CeIrIn}_5$  has a small lattice constant along the  $c$  axis compared with that of  $\text{CeRhIn}_5$ , and corresponds to  $\text{CeRhIn}_5$  under pressure.  $\text{CeIrIn}_5$  is a HF superconductor with  $T_C=0.4$  K at ambient pressure, which is again consistent with the picture that the  $\text{CeIn}_3$  layers in  $\text{CeIrIn}_5$  experience a high chemical pressure compared with those in  $\text{CeRhIn}_5$ .

In the present study, we have performed  $4d$ - $4f$  ( $h\nu = 122$  eV) and  $3d$ - $4f$  ( $h\nu = 881$  eV) resonant photoemission spectroscopy (RPES) measurements to study

<sup>1</sup> Corresponding author. Present address: Synchrotron Radiation Center, Japan Atomic Energy Research Institute, SPring-8, Mikazuki, Hyogo 679-5148 Japan. E-mail: fujimori@spring8.or.jp

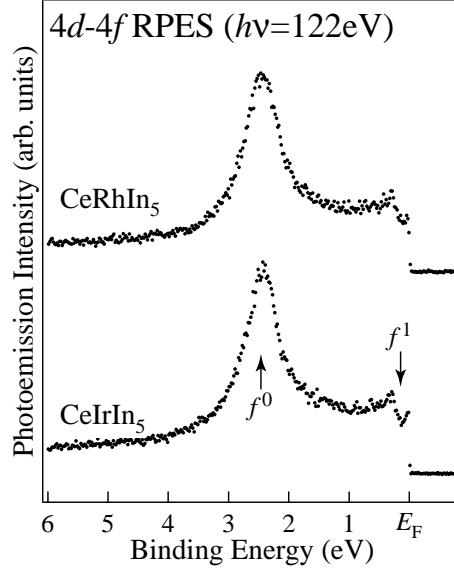


Fig. 1. 4d-4f resonant photoemission spectra of CeMIn<sub>5</sub>.

their electronic structures.

Single crystals of CeRhIn<sub>5</sub> and CeIrIn<sub>5</sub> were grown by the self-flux method described in Ref [3]. The 3d-4f resonant photoemission experiments ( $h\nu = 881.2$  eV) were performed at BL23SU of SPring-8. The sample temperature was kept at 15 K during the course of measurements, and overall energy resolution was set to about 200 meV. The 4d-4f resonant photoemission experiments ( $h\nu = 122$  eV) were performed at BL1 of Hiroshima Synchrotron Radiation Center (HiSOR). The sample temperature was kept at 10 K during the course of measurements, and overall energy resolution was set to about 20 meV

Figure 1 shows 4d-4f RPES spectra ( $h\nu = 122$  eV) of CeMIn<sub>5</sub>. The spectra consist of the  $f^1$  final-state peak located just below the Fermi level and the  $f^0$  final-state peak at about 2.4 eV. These spectral features have been observed in other Ce-based compounds, and understood within the framework of the single impurity Anderson model (SIAM)[4]. According to SIAM, the stronger the  $f$  electrons hybridize with conduction electrons, the stronger the  $f^1$  peak becomes. Here, the intensity of  $f^0$  final state peak, located at about 2.4 eV, is very strong, suggesting that the Ce 4f electrons are strongly localized.

Figure 2 shows the 3d-4f RPES spectra ( $h\nu = 881$  eV) of CeMIn<sub>5</sub>. The spectra consist of the  $f^1$  and  $f^0$  final-state peaks, same as 4d-4f RPES spectra, but their relative intensities are different from those in 4d-4f RPES spectra. The intensity of  $f^1$  final state peak is much stronger than those in 4d-4f RPES spectra. This is considered to be due to the surface insensitivities in the 3d-4f RPES experiments, compared with those in the 4d-4f RPES experiments[5].

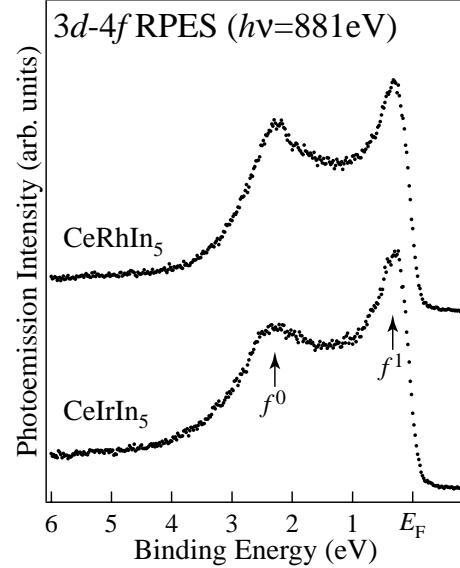


Fig. 2. 3d-4f resonant photoemission spectra of CeMIn<sub>5</sub>.

In these spectra, the intensity of  $f^0$  final state peak is still large even it is compared with those of the low  $T_K$  HF Ce-based compounds, such as CeRu<sub>2</sub>Si<sub>2</sub> ( $T_K = 22$  K)[5]. This again argues that Ce 4f electrons in these compounds are nearly localized. We also note that although both spectra are very similar, the  $f^1$  to  $f^0$  intensity ratio is somewhat large in CeIrIn<sub>5</sub> compared with that of CeRhIn<sub>5</sub>. This indicates that the hybridization between the conduction electrons and the Ce 4f electrons in CeIrIn<sub>5</sub> is slightly larger than that in CeRhIn<sub>5</sub>. This is consistent with the fact that CeIrIn<sub>5</sub> corresponds to CeRhIn<sub>5</sub> under pressure.

To summarize, we have found that Ce 4f electrons in CeMIn<sub>5</sub> ( $M=\text{Rh, Ir}$ ) are nearly localized. Although the Ce 4f electrons in CeIrIn<sub>5</sub> are more delocalized than those in CeRhIn<sub>5</sub>, they are more localized than low  $T_K$  compound like CeRu<sub>2</sub>Si<sub>2</sub>.

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