

# Third harmonic *ac* susceptibility measurements on $MgB_2$ bulk: irreversibility line and frequency dynamic behaviour

Daniele Di Gioacchino <sup>a,1</sup>, Paolo Tripodi <sup>b</sup>, Umberto Gambardella <sup>a</sup>

<sup>a</sup> *I.N.F.N., Frascati National Laboratory, via Enrico Fermi 40, 00044 Frascati, Italy*

<sup>b</sup> *E.N.E.A.-Centro Ricerche Frascati, via Enrico Fermi 45, 00044 Frascati, Italy*

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## Abstract

The third harmonic of the *ac* susceptibility of  $MgB_2$  high pressure (HP) bulk samples have been measured at the frequency of 1070Hz as a function of the temperature for different *dc* magnetic fields. The irreversibility line (IL) has been extracted from the temperature onsets. The IL has a  $(1 - t)^n$  dependence, with an "effective power law" exponent  $n=1.27$ . This exponent takes into account the interaction between the thermal fluctuations and flux pinning feature. Another IL data of HP samples, drawn out of VSM measurements at few Hertz has similar behaviour. This shows that the superconducting response is independent from frequency in the ranges of the temperature and the magnetic field studied. This analysis supports the fact that when the magnetic field is on and the critical state is formed the frequency creep decay is outside the frequency range analyzed.

*Key words:*  
 $MgB_2$ ; irreversibility line; dynamic transport

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## 1. Introduction

The recent discovery of the superconductivity in  $MgB_2$  [1], has stimulated large interest as this material shows various peculiar behaviours [2]: a) an electron-phonon superconducting behaviour, like conventional BCS superconductors; b) grain alignment does not play a crucial role in superconducting current transport [3].

Thus, from the viewpoint of practical applications, it is important to understand the magnetic IL as a function of the temperature [4] [5]. IL is generally defined in a resistive way by the appearance of a resistive state at each field [2] [3], or by VSM magnetic measurements [3]. A different way to detect the IL and the non-linear dynamics of the vortex is to measure the high harmonic susceptibility components,  $\chi_3 + i\chi_3''$  [6], in respect to the temperature and the DC magnetic field. This procedure permits to have a more sensitive

method and the contribution of the linear losses that limit the superconducting properties is ignored [7]. In this paper we present the onset of  $\chi_3$ , measured at 1070Hz as function of temperature and DC magnetic field of the  $MgB_2$  bulk samples. Comparisons of the measured IL of  $MgB_2$  samples with other IL data by VSM measurements of similar samples have been done.

## 2. Results and discussion

$MgB_2$  bulk samples of  $(1 \times 2 \times 5.6 \text{ mm}^3)$  produced by sintered reaction [8] have been measured. The critical temperature is above 38K. The high harmonics susceptibility has been measured in a double coil susceptometer (TARI facilities). The sample is placed on a sapphire holder, this set-up is cooled in a thermally controlled He gas flow cryostat provided with 8 Tesla superconducting magnet. The ac driving field has 1070Hz frequency and 6 Gauss amplitude. The external mag-

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<sup>1</sup> E-mail: daniele.digioacchino@lnf.infn.it

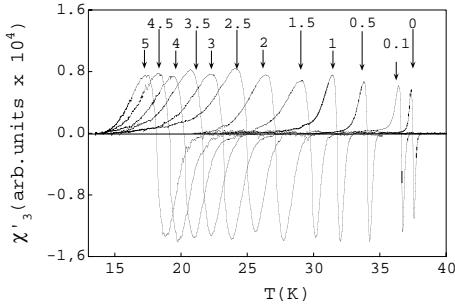


Fig. 1.  $\chi'_3$  of  $MgB_2$  bulk sample at  $B_{ac}=6G$ ,  $f=1070Hz$ ,  $B_{dc}=0-5T$ .

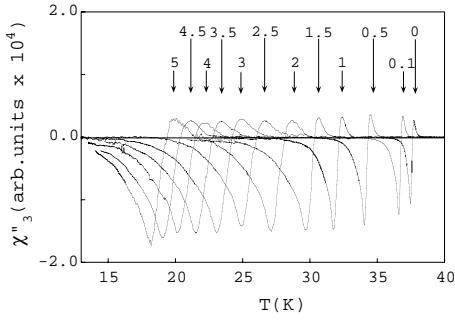


Fig. 2.  $\chi''_3$  of  $MgB_2$  bulk sample at  $B_{ac}=6G$ ,  $f=1070Hz$ ,  $B_{dc}=0-5T$ .

netic field is parallel to the length of the sample. In fig.1 the real and imaginary parts of  $\chi_3$  are shown for dc magnetic field  $B_{dc}$  in the range (0-5) Tesla.

The  $\chi_3$  signal rises when the current-voltage characteristic becomes non-linear and the multiharmonic components are generated. The criterion used to define the IL is determined by the temperature at which the third harmonic amplitudes  $\chi_3$  start to deviate from the zero value. IL determined by the onset of  $\chi_3$  measurements and the fit function  $(1 - T/T_c)^n$ , described in the vortex glass/liquid theory, are shown in fig. 2. The  $n$  adjustable parameter describes how the thermal fluctuations can give rise to the vortex lattice melting [9]. In presence of quantum fluctuations  $n$  is smaller than 2 [10]. This description of the melting line with an "effective power law" exponents is a compact way description because it depends on: disorder, dimensionality, lattice flux elastic response; normally described as scaling law of the pinning versus temperature and magnetic field [11]. However, it is difficult to extract clear physical evidence from the above power law [10].

Moreover, in fig.2 another IL data (x), determined by VSM measurements at few Hertz [3] are shown. The data comparison shows that in  $MgB_2$  bulk, under these magnetic fields and temperature conditions, the flux dynamics remains similar in the frequency range of induced current. Infact [12] in the three-dimensional phase diagram H-T-J of superconductor, the current

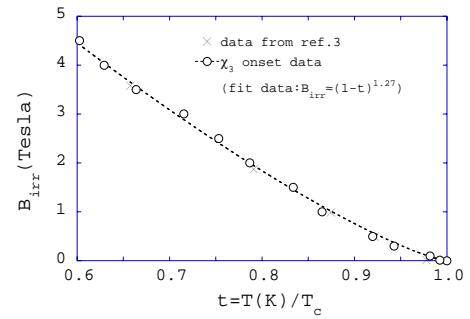


Fig. 3. IL line of the  $MgB_2$  bulk sample based of  $\chi_3$  onsets (—o—). (x) are IL data based on VSM measurements .

axes can be interpreted as time axes (or 1/frequency) via the relation:

$$j(t) \propto [\ln(\frac{t}{t_0})]^{-\frac{1}{\mu}}. \quad (1)$$

### 3. Conclusion

An accurate measurement of the irreversibility line (IL) of the  $MgB_2$  bulk sample has been done from the third harmonic of the *ac* susceptibility temperature onsets. The temperature behaviour of the IL like  $(1 - T/T_c)^{1.27}$  has been found which shows a correction of the simple quadratic power-law. This value takes into account the complex interaction between the thermal fluctuations and flux pinning feature. Moreover there is an indication that the critical state formed in the presence of magnetic field (0-5 Tesla) is independent from the frequency range (1Hz-1KHz) of the induced current.

### References

- [1] J. Nagamatsu, N. Nakagawa, T. Muranaka, Y. Zenitani, J. Akimitsu *Nature* **410** (2001) 63.
- [2] C. Buzea, T. Yamashita, *Superc. Science and Tech.* **14** (2001) R115.
- [3] H. Kumakura, Y. Takano, H. Fujii, K. Togano, H. Kito, H. Ihara *Physica C* **363** (2001) 179.
- [4] D.C. Larbalestier, *et al*, *Nature* **410** (2001) 186.
- [5] H. H. Wen, *et al*, *Physica C* **363** (2001) 170.
- [6] A. Shaulov, D. Dorman, *Appl. Phys. Lett.* **53** (1988) 2680.
- [7] D. Di Gioacchino *et al*, *Phys. Rev. B* **59** (1999) 11539.
- [8] Edison S.P.A. patent pending.
- [9] A. Houghton, R.A. Pelcovits, A. Subdø, *Phys. Rev. B* **40** (1989) 6763
- [10] G. Blatter, *et al*, *Rev. Mod. Phys.* **66** (1994) 1234-1235
- [11] W. A. Fietz, W.W. Webb, *Phys. Rev. B* **178** (1969) 657
- [12] G. Blatter, *et al*, *Rev. Mod. Phys.* **66** (1994) fig.21 1228