

# Bonding nature and wave function around the Fermi level of MgB<sub>2</sub>-related compounds

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## Abstract

We present first-principles molecular-orbital calculations for the 39-K superconductor MgB<sub>2</sub> and related diborides using the DV- $X\alpha$  method. The electronic structure of a slab cluster embedded in the Madelung potential is determined self-consistently in terms of charge transfer between Mg and B. In contrast with the simplest ionic picture Mg<sup>2+</sup>B<sub>2</sub><sup>-</sup>, the negative charge of B is estimated to be 0.39. Our results, which include total and partial density of states, overlap population in the neighboring B-B and Mg-B bonds and wave functions around the Fermi energy  $E_F$  consisting of B 2p<sub>x,y</sub> orbitals in the B planes, indicate the incomplete filling of the  $\sigma$  orbitals in MgB<sub>2</sub>. We discuss possible ways to achieve higher transition temperatures in non-periodic systems such as surfaces, interfaces and artificially layered hetero-structures based on this approach.

*Key words:* diborides; electronic structure; molecular orbitals; Madelung potential

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## 1. Introduction

There has been rapid progress in our understanding of the electronic structure of MgB<sub>2</sub> since the first announcement of superconductivity at 39 K [1]. In spite of the crystal-structural similarity to intercalated graphite with carbon replaced by boron, MgB<sub>2</sub> represents a new class of superconductors. The simplest ionic picture would suggest that the Mg donates two outer electrons to B planes. This picture is supported by recent band structure calculations [2]. Although MgB<sub>2</sub> is formally isoelectronic to graphite, its unique feature is the incomplete filling of  $\sigma$  bands corresponding to  $sp^2$ -hybridized bonding within the graphite-like B planes.

In this paper, we attempt to examine the relationship between the  $\sigma$  band structure in  $k$  space and the in-plane  $sp^2$  bonding nature in real space for MgB<sub>2</sub> and some related diborides using first-principles molecular-orbital calculations.

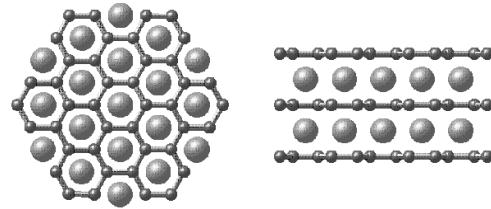


Fig. 1. Top view and cross-sectional view of our cluster model employed in the DV- $X\alpha$  calculation.

## 2. Calculational Methods

Our approach is based on the DV- $X\alpha$  method [3] that utilizes discrete variational (DV) integral and local density approximation ( $X\alpha$ ). MgB<sub>2</sub> occurs in the hexagonal omega (C32) structure. To reproduce its bulk electronic structure, a slab cluster, illustrated in Fig. 1, is embedded in the Madelung potential (MP), which is to be evaluated self-consistently in terms of charge transfer from Mg to B. Numerical atomic orbitals obtained in the crystal field are used as basis

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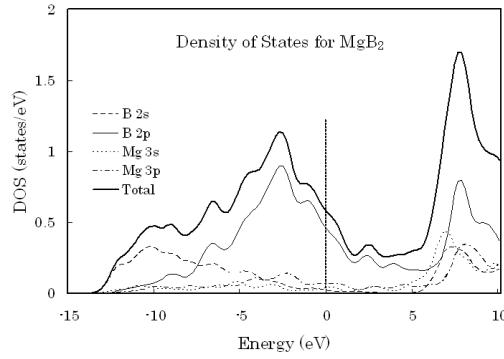


Fig. 2. Total and partial density of states for  $\text{MgB}_2$  ( $a = 3.084$ ,  $c = 3.522$ ).

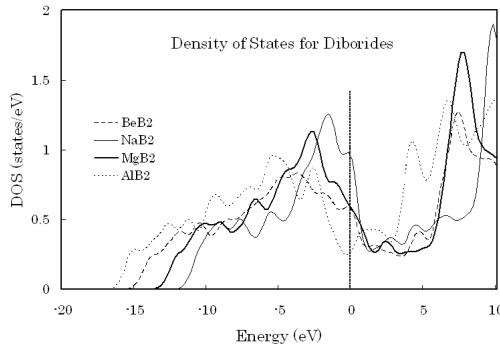


Fig. 3. Density of states for  $\text{BeB}_2$  ( $a = 2.87$ ,  $c = 2.85$ ),  $\text{NaB}_2$  ( $a = 3.02$ ,  $c = 4.19$ ),  $\text{MgB}_2$  and  $\text{AlB}_2$  ( $a = 3.009$ ,  $c = 3.262$ ).

functions. Unlike the band calculations, the periodicity of crystal lattices is not essential for this approach.

### 3. Results and Discussion

Figure 2 shows the total and partial density of states (DOS) for  $\text{MgB}_2$ . The states around the Fermi level  $E_F$  are dominantly derived from B 2p orbitals; thus, the B  $2p_{x,y}$  electrons contribute both to the in-plane B-B bonding and to two-dimensional superconductivity. In Fig. 3, the DOS curves of  $\text{BeB}_2$ ,  $\text{NaB}_2$ ,  $\text{MgB}_2$  and  $\text{AlB}_2$  are compared, where the converged negative charge of B are 0.32, 0.16, 0.39 and 0.52, respectively. Except for  $\text{AlB}_2$  with completely filled  $\sigma$  bands, the overall shapes are quite similar. The relatively small values of the charge transfers compared to the simple ionic picture indicate that the transferred electrons are not well localized at the B sites.

In the framework of Mulliken's population analysis, bond overlap population is a good measure of the covalency of the bonds. The calculated values of neighboring B-B (Mg-B) bonds are 0.76 (0.12), 0.74 (0.13),

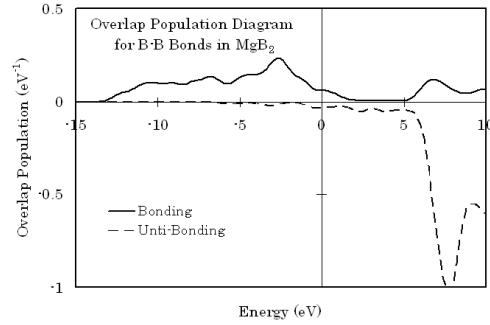


Fig. 4. Overlap population diagram for the neighboring B-B bonds in the B planes of  $\text{MgB}_2$ .

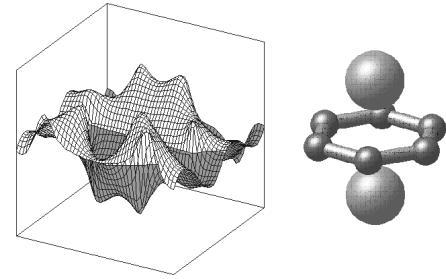


Fig. 5. Wave function at  $E_F$  in the B planes of  $\text{MgB}_2$ .

0.70 (0.20) and 0.66 (0.28) for  $\text{BeB}_2$ ,  $\text{NaB}_2$ ,  $\text{MgB}_2$  and  $\text{AlB}_2$ , respectively. The existence of  $\sigma$  holes can be understood from the overlap population diagram for the neighboring B-B bonds of  $\text{MgB}_2$  shown in Fig. 4. The bonding component of the overlap population extends over  $E_F$ . The results for  $\text{BeB}_2$  and  $\text{NaB}_2$  are also similar, apart from their  $\sigma$  band broadening.

The wave function at  $E_F$  (see Fig. 5) consists of dominantly B  $2p_{x,y}$  orbitals in the B planes of  $\text{MgB}_2$ . Its slightly asymmetric form suggests that the corresponding  $k$  lies in the vicinity of the  $\Gamma$  point in the Brillouin zone. As for  $\text{AlB}_2$ , these  $2p_{x,y}$  orbitals are completely filled and do not contribute to conduction.

In the process of iteration, we have found that depending on the MP, the relative location of the  $\sigma$  and  $\pi$  bands shifts on the energy scale; accordingly, the hole doping in the B planes is sensitive to the MP along  $c$  axis. We expect that higher  $T_c$  could be achieved in artificially layered systems with modified MP.

### References

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