

Antivortices in the mixed paramagnetic-orbital pair breaking regime

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Abstract

We consider a thin superconducting layer with a magnetic field direction slightly tilted from the plane-parallel orientation. In such a situation, which generalizes the FFLO state, both paramagnetic and orbital pair-breaking effects must be taken into account. We minimize the quasiclassical free energy to obtain the stable order parameter structure near the upper critical field, extending previous work [U.Klein et al., J. Low Temp. Phys. 118 (2000) 91] with regard to magnetic field contributions. We find two-dimensional periodic states, characterized by Landau quantum numbers $n > 1$ showing antivortices, e.g. a unit-cell with three order parameter zeros (two vortices and an antivortex) carrying totally a single flux quantum. The underlying physical mechanism is different from a recent prediction of antivortices in mesoscopic samples of triangular shape [L. F. Chibotaru et al., Phys. Rev. Lett. 86 (2001) 1323].

Key words: Antivortices; FFLO state; orbital pair breaking; paramagnetic pair breaking

A superconducting antivortex changes its phase by -2π when encircling the centre and carries a corresponding flux quantum of negative sign. Such entities are sometimes introduced as purely mathematical constructs (called mirror vortices), in order to fulfill the Ginzburg-Landau(GL) boundary conditions at surfaces of superconductors [1]. They play also a role above the Kosterlitz-Thouless transition in thin superconducting films. However, individual antivortices have apparently not yet been identified experimentally.

Recently, a two-dimensional periodic state with a unit cell carrying a single flux quantum, two ordinary vortices, and a single antivortex has been predicted [2,3]. This state was found by minimizing the free energy for a thin-film configuration where both paramagnetic and orbital pair-breaking takes effect. Very recently, nucleation of antivortices in a mesoscopic superconducting triangle has been predicted [4]. Good agreement of the theoretical phase boundary with experimental data has been found [4]. The physical origin of the antivortices is quite different in these two cases.

In the first case, we have a bulk phenomenon related to competing pair breaking effects; in the second case one has a combination of size effects and symmetry requirements. Formally, antivortices appear in both cases as a consequence of a mismatch between the number of zeros of the relevant set of basis functions and the required number of flux quanta per unit cell.

In this note we report on a study of antivortices in the mixed paramagnetic-orbital pair-breaking regime for superconductors of *finite* GL-parameter κ . Thereby, previous work [2,3] for London limit superconductors is extended to include magnetic field contributions.

We consider a quasi-two-dimensional superconductor in an external magnetic field whose direction φ can be varied (φ is the angle between the field and the conducting x,y -plane) [5]. We study a clean superconductor with a cylindrical Fermi surface of axial symmetry. This model describes conventional layered superconductors with (nearly) perfect 2D-behavior, i.e. extreme anisotropy and nearly perfect decoupling of the conducting layers.

The complete phase boundary H_{c2} between normal-conducting and superconducting state has been cal-

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culated by Shimahara and Rainer [6]. It has a cusp-like shape [1,4] with different pieces of the H_{c2} curve belonging to different values of the Landau quantum number n ($n = 0, 1, \dots$). For large φ one finds $n = 0$, dominating orbital pair breaking and the usual vortex state. With decreasing φ states with higher n appear close to the plane-parallel orientation. For $\varphi \rightarrow 0$ one finds $n \rightarrow \infty$, only paramagnetic pair breaking remains, and the stable state [6] is the Fulde-Ferrell-Larkin-Ovchinnikov (FFLO) state [7–13].

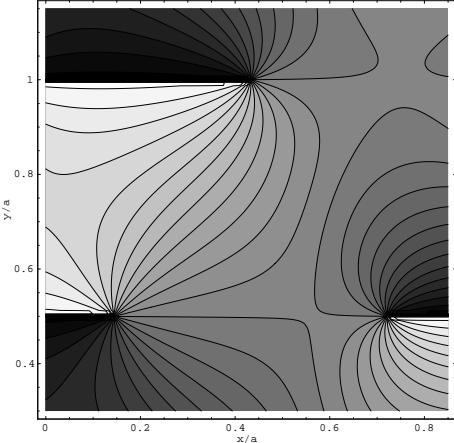


Fig. 1. Contour plot of phase of order parameter in triangular unit cell for $n = 2$ showing three order parameter zeros and associated branch cuts.

The structure of the superconducting states below H_{c2} has been calculated in Ref. [2] for a high- κ material. The quasi-periodic order parameter for a 2D-lattice with an angle α between the two basis vectors of length a and b may be constructed by standard methods:

$$\psi_n(x, y) = AC_n \sum_{m=-\infty}^{m=+\infty} \exp \left(-2\pi i \frac{b}{a} \cos \alpha \frac{m(m+1)}{2} \right) \times \exp \left(\frac{i}{\hbar} k_m x \right) h_n (y - mb \sin \alpha). \quad (1)$$

Here, $k_m = (2\pi\hbar/a)m$ and the function h_n is proportional to a Hermite polynomial of order n (see Ref. [2] for further details). To find the stable lattice structure (the parameters a, b, α) for given n, T , the quasiclassical free energy including nonlinear terms has to be minimized. For $n \geq 2$ several stable states showing antivortices have been found. The $n = 2$ state (see Fig. 3 of Ref. [2]) appearing at the largest φ , is the one most easily accessible experimentally.

This treatment has been generalized with regard to effects of finite κ [14]). Taking magnetic field contributions into account allows one to recover (from the $n = 0$ state) Abrikosovs solution in the GL-limit and to identify the transition between type II and type I behavior. An essential result with regard to antivortex states

is the fact that the $n = 2$ structure, found in Ref. [2] remains stable (in contrast to most other states) for a very wide range of κ -values (calculations have been performed for $0.1 \leq \kappa \leq 100$ at two T). The equilibrium structure is the same as in [2]. To show the presence of antivortices more explicitly, the *phase* of the order parameter is plotted in Fig. 1 (its absolute value is shown in Fig.3 of [2]). There are three zeros of the order parameter per unit cell, each one being the starting point of an associated branch cut. When traversing a branch cut from the black side to the white side, the phase increases by 2π . Encircling the vortex centers in anticlockwise direction, one sees that the upper vortex is an antivortex (phase change by -2π) while the two lower vortices are of the ordinary ($+2\pi$) type. The unit cell as a whole carries a single flux quantum, corresponding to a total phase change of $+2\pi$.

Summarizing, we considered a 2D superconducting film in a tilted magnetic field where both orbital and paramagnetic pair-breaking takes effect. Using a material with a circular Fermi surface and an arbitrary GL parameter κ , a quasiperiodic superconducting state ($n = 2$) with a vortex-antivortex pair and an ordinary vortex per unit cell should appear for applied field nearly parallel (within an angle of the order of 1 degree) to the conducting plane.

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References

- [1] M. Tinkham, *Introduction to Superconductivity*, McGraw-Hill, New York, 1975.
- [2] U. Klein, D. Rainer, H. Shimahara, J. Low Temp. Phys. **118** (2000) 91.
- [3] U. Klein, Physica B **284-288** (2000) 713.
- [4] L. F. Chibotaru, A. Ceulemans, V. Bruyndoncx, V. V. Moshchalkov, Phys. Rev. Lett. **86** (2001) 1323.
- [5] L. N. Bulaevskii, Sov. Phys. JETP **38** (1974) 634.
- [6] H. Shimahara, D. Rainer, J. Phys. Soc. Jpn **66** (1997) 3591.
- [7] P. Fulde, R. A. Ferrell, Phys. Rev. **135** (1964) A550.
- [8] A. I. Larkin, Y. N. Ovchinnikov, Sov. Phys. JETP **28** (1969) 1200.
- [9] H. Burkhardt, D. Rainer, Ann. Physik **3** (1994) 181.
- [10] M. S. Nam, J. A. Symington, J. Singleton, S. J. Blundell, A. Ardavan, J. A. A. J. Perenboom, M. Kurmoo, P. Day, J. Phys.: Condens. Matter **11** (1999) L477.
- [11] S. Manalo, U. Klein, J. Phys.: Condens. Matter **12** (2000) L471.
- [12] H. Shimahara, Phys. Rev. **B62** (2000) 14541.
- [13] S. Manalo and U. Klein, Phys. Rev. **B65** (2002) 144510.
- [14] U. Klein, to be published