

Specific heat measurements of pyrochlore-type $R_2Mo_2O_7$ ($R=Nd-Ho$)

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Abstract

Specific heat measurements have been performed on pyrochlore-type $R_2Mo_2O_7$ ($R=Nd-Ho$) single crystals. For the ferromagnetic compounds with $R=Nd-Gd$ ($T_C \approx 50-90$ K), Schottky-like specific heat anomalies corresponding to the level splitting for the 4f electrons of R^{3+} are observed at low temperatures, in contrast to the specimens with $R=Tb-Ho$ which show spin-glass-like order below $T_g \approx 20-25$ K. Low-temperature specific heat data of $Sm_2Mo_2O_7$ under a magnetic field of 10 T are presented.

Key words: pyrochlore; specific heat; spin-glass;

Pyrochlore oxides with a general formula $A_2B_2O_7$ have a face centered cubic structure and each of R and Mo atoms form a three-dimensional network of corner-sharing tetrahedra. In particular, there has been great deal of interest in pyrochlore molybdates $R_2Mo_2O_7$, because the systems with $R=Y$ and $Tb-Er$ are insulating, geometrically frustrated and undergo a phase transition into a spin-glass-like state despite the absence of apparent structural disorder[1,2], while the systems with $R=Nd-Gd$ exhibit ferromagnetic order with metallic conductivities. Very recently, it has been pointed that the magnetic phase boundary is correlated with the metal-insulator crossover[3] and the structural variation at T_C in $Nd_2Mo_2O_7$ suggests some double-exchange interaction mechanism in metallic $R_2Mo_2O_7$ [4].

In the present work, we have performed the precise measurements of specific heat for $R_2Mo_2O_7$ ($R=Nd-Ho$) using single crystal specimens, to throw further light on the magnetic ordering. To our knowledge, the specific heat measurements have been so far limited on the polycrystalline specimens of $R_2Mo_2O_7$ with $R=Sm, Gd, Y$ [5]. Single crystals of $R_2Mo_2O_7$ were

grown in an Ar atmosphere by a floating-zone method and confirmed to be a single crystal by Laue reflection. Specific heat was measured by a thermal relaxation method.

In metallic $R_2Mo_2O_7$, the ferromagnetic ordering of Mo^{4+} ions occurs at T_C , whereas the ordering of R^{3+} ions are interpreted to gradually develop below T_C and become significant below T_F ($\ll T_C$), as evidenced by recent neutron scattering experiments for $Nd_2Mo_2O_7$ ($T_C \sim 93$ K, $T_F \sim 20$ K)[6]. We show the temperature T dependence of magnetic specific heat C_m for $Sm_2Mo_2O_7$ in Fig.1. In deriving C_m , pyrochlore-type $Y_2Ti_2O_7$ which is insulating and nonmagnetic was used to eliminate the lattice contribution C_{latt} . In Fig. 1, a sharp anomaly corresponding to the ferromagnetic ordering of Mo^{4+} ions is observed at T_C (≈ 80 K), in contrast to the behavior of the polycrystalline specimen in the previous work[5]. The released magnetic entropy $S = \int C_m/T dT$ between 70 and 90 K we estimate is only ~ 10 percent of that expected for the ordering of Mo^{4+} ($4d^2$, $S=1$), similar to perovskite manganites[7]. The anomaly at T_C was also observed for the specimen with $R=Nd$ but not for $R=Gd$, which is in the vicinity of spin-glass-like state.

In Fig. 1, a Schottky-like anomaly, which is associ-

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ated with the ordering of Sm^{3+} ions, is observed at low- T and was also observed for the specimens with $\text{R}=\text{Nd}$ and Gd . The estimated magnetic entropies for these anomalies nearly correspond to that of a two-level system for $\text{R}=\text{Nd}$ and Sm and an eight-level system for $\text{R}=\text{Gd}$, although the peaks of the anomalies are small and broad, compared with standard multilevel Schottky anomalies[5]. In the inset of Fig.1, low- T $C_m(T)$ data for $H=0$ and 10 T parallel to the (111) direction are displayed. In the inset, the peak temperature of the anomaly T_p becomes slightly lower by applying a magnetic field of 10 T, indicating that the splitting of the doublet for 10 T ($\Delta_{10\text{T}}$) is smaller than that for $H=0$ (Δ_0), though the behaviors with $\Delta_{H \neq 0} > \Delta_0$ have been found in various materials. To derive $\Delta_{10} < \Delta_0$, we need to assume that the ordering of the Sm^{3+} moments μ_{Sm} is arising from antiferromagnetic Sm-Mo interactions inferred by the spin arrangement in $\text{Nd}_2\text{Mo}_2\text{O}_7$, where the net magnetization of the Nd and Mo sublattices are opposite in direction to each other[6]. In this case, the effective molecular field at the Sm sites H_m would be antiparallel in average to the applied magnetic field H . Accordingly, we may give the splitting of the doublet as, $\Delta_0 = 2\mu_{\text{Sm}}H_m(0)$ for $H=0$, and $\Delta_{10\text{T}} = 2\mu_{\text{Sm}}[H_m(10) - 10]$ for $H=10$ T. These equations satisfy $\Delta_{10} < \Delta_0$ for $H_m(0) \leq H_m(10)$ and yield $\mu_{\text{Sm}} = 0.065\mu_B$ when $H_m(0) = H_m(10)$. This value of μ_{Sm} is fairly smaller than the free-ion value of $0.7\mu_B$, similar to that of Nd^{3+} in $\text{Nd}_{0.67}\text{Sr}_{0.33}\text{MnO}_3$ [7].

Next, we show the T -dependence of specific heat C_p for the specimens with $\text{R}=\text{Tb-Ho}$ in the temperature range $2 \leq T \leq 20$ K in Figs. 2(a)-2(c). As seen in Figs. 2(a) and 2(b), the specimens with $\text{R}=\text{Tb}$ and Dy do not show any Schottky-like anomaly at low- T , in contrast to those with $\text{R}=\text{Nd-Gd}$, although $C_p(T)$ for $\text{R}=\text{Dy}$ appears to show a very slight hump below 10 K. On the other hand, a low- T anomaly, which was a clear peak but not Schottky-like in shape in the $C_m (=C_p - C_{\text{latt}})$ versus T plot, is observed at ~ 5 K for $\text{R}=\text{Ho}$ in Fig. 2(c). No anomaly was observed at $\sim T_g$ for $\text{R}=\text{Tb-Ho}$. The behaviors of $C_p(T)$ for $\text{R}=\text{Tb}$ and Dy are also forming a contrast with those for $\text{Tb}_2\text{Ti}_2\text{O}_7$ and $\text{Dy}_2\text{Ti}_2\text{O}_7$, where $C_p(T)$ exhibits a peak at low- T [8,9]. It should be noted that the specific heat anomaly associated with R^{3+} ions remarkably depends on the magnetic order at the Mo site, indicating that the ordering of the R^{3+} moments are much affected by the Mo^{4+} moments through interactions. The absence of anomaly at low- T for $\text{R}=\text{Tb}$ and Dy is attributable to the R^{3+} moments involved in the spin freezing of the Mo^{4+} moments at T_g , so that the magnetic entropies are released in the wide T -range above and below T_g . In $\text{Ho}_2\text{Mo}_2\text{O}_7$, it has been found that the magnetic susceptibility obeys a Currie-Weiss law and is fairly slightly history dependent below T_g , contrasting the behaviors in $\text{Tb}_2\text{Mo}_2\text{O}_7$, suggesting that the Ho^{3+}

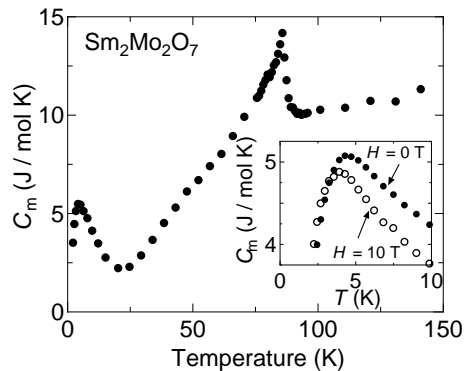


Fig. 1. Temperature dependence of magnetic specific heat C_m for $\text{Sm}_2\text{Mo}_2\text{O}_7$. Low- T $C_m(T)$ data for $H=0$ and 10 T are shown in the inset.

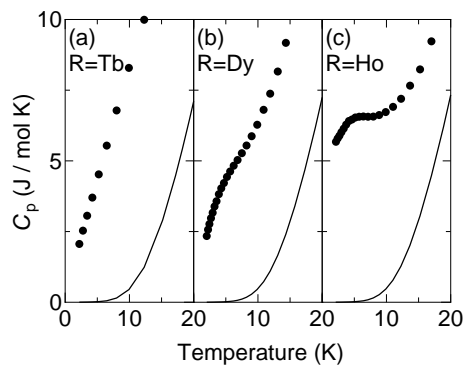


Fig. 2. Temperature dependence of specific heat for $\text{R}_2\text{Mo}_2\text{O}_7$ with $\text{R}=\text{Tb}$ (a), Dy (b), Ho (c). The solid lines indicate the lattice contributions.

moments are remain almost paramagnetic even below T_g [10]. The appearance of the low- T anomaly for $\text{R}=\text{Ho}$ is thought of due to relatively weak R-Mo couplings.

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