

Impossibility of superconducting state in multiwall carbon nanotubes and single crystal graphite

Anatoly I. Romanenko^{a,1}, Alexandr V. Okotrub^a, Lubov G. Bulusheva^a,
Olga B. Anikeeva^a, Nikolai F. Yudanov^a, Cheng Dong^b, Yongming Ni^b

^a*Institute of Inorganic Chemistry, Siberian Branch of Russian Academy of Science, Novosibirsk, Russia*

^b*National Laboratory for Superconductivity, Institute of Physics CAS, Beijing, China*

Abstract

The magnetic susceptibility χ as of pristine and brominated arc-produced sample of multiwall carbon nanotubes (MWNTs) so single crystal graphite (SCG) was measured from 4.5 to 400 K. An additional contribution $\Delta\chi(T)$ to diamagnetic susceptibility $\chi(T)$ of carbon materials was found at $T < 50$ K for all samples. It is shown that $\Delta\chi(T)$ are dominated by quantum correction to χ for interaction electrons (interaction effects-IE). The IE shows a crossover from two-dimensional to three-dimensional at $B = 5.5$ T for MWNTs and always three-dimensional for SCG. The effective interaction between electrons for interior layers as MWNTs so SCG are repulsion and the electron-electron interaction λ_c was estimated to be $\lambda_c \sim 0.26$ for MWNTs and $\lambda_c \sim 0.06$ for SCG.

Key words: Electron-electron interaction; Carbon nanotubes, Single crystal graphite; Brominated carbon nanotubes

From the time of discovery nanotubes one of the most important problems is the possibility of a superconducting state in them. There are the superconductivity in single-walled carbon nanotubes up to 20 K [1] and the possibility of superconductivity above 400 K in MWNTs [2]. The nature of a superconducting state is the electron-electron interaction. On the other hand electron-electron interaction is exhibited in electronic transport properties of conductors in normal state - so-called quantum interference effects - interaction effects (IE) [4,3]. IE are connected with the correction to density of states of conduction electrons in a results of quantum interferences of electrons at their diffuse motion in random conductors. But in such systems the one-particle processes, so-called weak localization (WL) [4] and weak antilocalization (WAL) [5], always accompany with IE. The observation of IE corrections

(IEC) to χ is very important for partition of including of WL, WAL, and IE to different physical properties. From all these corrections only IE contributes to χ .

For observation of IEC to χ it is necessary to divide the contributions connecting with IEC and much more on quantity a magnetic susceptibility of sample, and to exclude the contribution of paramagnetic impurities. Earlier, with the using of relaxation processes in Mo_2S_3 [6], we changed the contribution connected with IE correction to χ by quenching of high-temperature metastable state of a sample. As a result, with the using of difference contribution to χ in equilibrium and metastable states, we received the IEC to χ in the pure state [7]. In this work, with the using of chemical modification of a MWNTs sample (brominated), we singled out IEC to χ in laminated structures based on MWNTs and SCG.

The material which contained MWNTs was synthesized by arc discharge graphite evaporation. A nanotube content in the inner part of carbon deposit grown on the cathode was estimated by transmission electron microscopy to reach about 80% [8]. Tubes have from 2

¹ Corresponding author. Present address: Institute of Inorganic Chemistry, Siberian Branch of Russian Academy of Science, Novosibirsk, Russia. E-mail: romanenk@casper.che.nsk.su

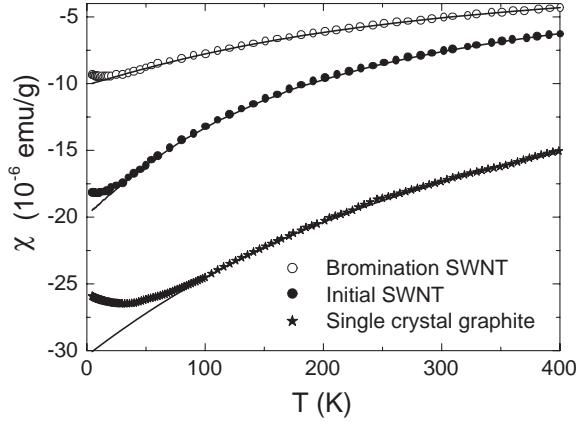


Fig. 1. The temperature dependence of magnetic susceptibility χ of SWNT, bromination SWNT, and SCG measurement at $H=0.01$ T.

to 30 shells with an outer diameter of 60-150 Å. Scanning electron microscopy revealed a predominant orientation of MWNTs was perpendicular to the deposit growth axis.

According to experimental and theoretical data, the basic contribution in χ of quasi-two-dimensional graphite gives orbital magnetic Susceptibility χ_{or} connected with extrinsic carriers n_0 [9]. Fig. 1 presents the magnetic susceptibility χ of investigation samples as a function of temperature. The data in Fig. 1 show that at $T < 50$ K there is an additional contribution $\Delta\chi_{or}(T) = \chi(T) - \chi_{or}(T)$ to $\chi(T)$. According to theoretical calculations [3] only electron-electron interactions can contribute to magnetic susceptibility. Absolute value of $\Delta\chi_{or}(T)/\chi_{or}(T)$ at all magnetic fields applied to the pristine sample increases with decreasing temperature that has been predicted for IEC in the systems characterized by electron-electron repulsion [3]. Hence, at $B = 5.5$ T a transfer from two-dimensional IEC to three-dimensional one takes place for MWNTs. At lower magnetic field the interaction length $L_I(T) = (\hbar D/k_B T)^{1/2}$ is much less than the magnetic length $l_B = (\hbar c/2eB)^{1/2}$, which in turn becomes dominant at high field. An estimation of the characteristic lengths gives respectively the value of $L_I(4.2K) = 130$ Å and the value of $l_B = 100$ Å at $B = 5.5$ T. IE correction for SCG was three-dimensional at all fields.

In summary, we have investigated the additional contribution to temperature dependence of $\Delta\chi_{or}(T)/\chi_{or}(T)$ of lamination structure of MWNTs and SCG at $T < 50$ K. It is shown that $\Delta\chi_{or}(T)/\chi_{or}(T)$ is connected with quantum correction to magnetic susceptibility for interaction electron. At low field this correction is two-dimensional for MWNTs. At $B = 5.5$ T was observed three-dimensional correction to magnetic susceptibility. IEC for SCG

remain three-dimensional at low fields up to 0.001 T. The crossover of IEC from 2d to 3d for MWNTs is connected with decreasing of magnetic length up to value less then typical mean diameter of nanotubes. It is shown that bromination of samples leads to increasing of n_0 from $n_{0ini} \sim 3 \times 10^{10}$ cm⁻² for pristine sample up to $n_{0Br} \sim 10^{11}$ cm⁻² for brominated samples. But $\Delta\chi_{or}(T)/\chi_{or}(T)$ did not changed when n_0 increases, what is in full agreement with theoretical predictions. From $\Delta\chi_{or}(4.5K)/\chi_{or}(4.5K)$, we estimated the constant of electron-electron interaction $\lambda_c \sim 0.26$ for MWNTs and $\lambda_c \sim 0.06$ for SCG. This interaction is repulsion for interior layers, which give the domination contribution to $\Delta\chi_{or}(T)/\chi_{or}(T)$ as a integration value. The repulsion character of electron-electron interaction exclude the possibility of superconducting states in interior layers of MWNTs and SCG.

Acknowledgements

The work was supported by Russian Foundation of Basic Research (Grants No: 00-02-17987; 02-02-06068), INTAS (Grant Nos 00-237, 01-254), and Siberian Branch of Russian Academy of Science (Grant No 61).

References

- [1] Z. K. Tang, Lingyun Zhang, N. Wang et al., *Science* **292** (2001) 2462.
- [2] Guo-meng Zhao, Y. S. Wang, preprint cond-mat/0111268, (2001).
- [3] B.L. Al'tshuler, A.G. Aronov, A.Yu. Zyuzin, *Sov. Phys. JETP* **53** (1983) 889.
- [4] P.A. Lee, T.V. Ramakrishnan, *Rev. Modern Physics* **57** (1985) 287.
- [5] S. Hikami, A.I. Larkin, Y. Nagaoka, *Prog. Theor. Phys.* **63(2)** (1980) 707.
- [6] A.I.Romanenko, A.K.Dzhunusov, I.N. Kuropynatntsk, and E.V. Kholopov, *Sov. Phys. JETP Lett.* **41** (1985) 237.
- [7] A.I.Romanenko, F.S. Rakhenkulov, V.N. Ikorski, P.S. Nikitin, *Sov. Phys. JETP Lett.* **42** 377 (1985) 377.
- [8] A.V. Okotrub et al., *Appl. Phys. A* **71** 1 (2000) 481.
- [9] A.S. Kotozonov, and S.V. Kuvshinnikov, *Phys. Lett. A* **229** (1997) 377.
- [10] J. Heremans, C.H. Olk, and D.T. Morelli, *Phys. Rev. B* **49** 15122 (1994) 15122.
- [11] F. Tsui, L. Jin and O. Zhou, *Appl. Phys. Lett. B* **76**, (2000) 1452.
- [12] H. Ajiki and Ando, *J. Phys. Soc. Jpn.* **62**, (1993) 2470.
- [13] J.P. Lu, *Phys. Rev Lett.* **74**, (1995) 1123.
- [14] A.M. Rao, P.C. Eklund, S. Bandow, A. Thess, and R.E. Smalley, **388** (1997) 257.
- [15] L.X. Benedict, S.G. Louie, and M.L. Cohen, *Solid State Commun.* **100** (1996) 177.