

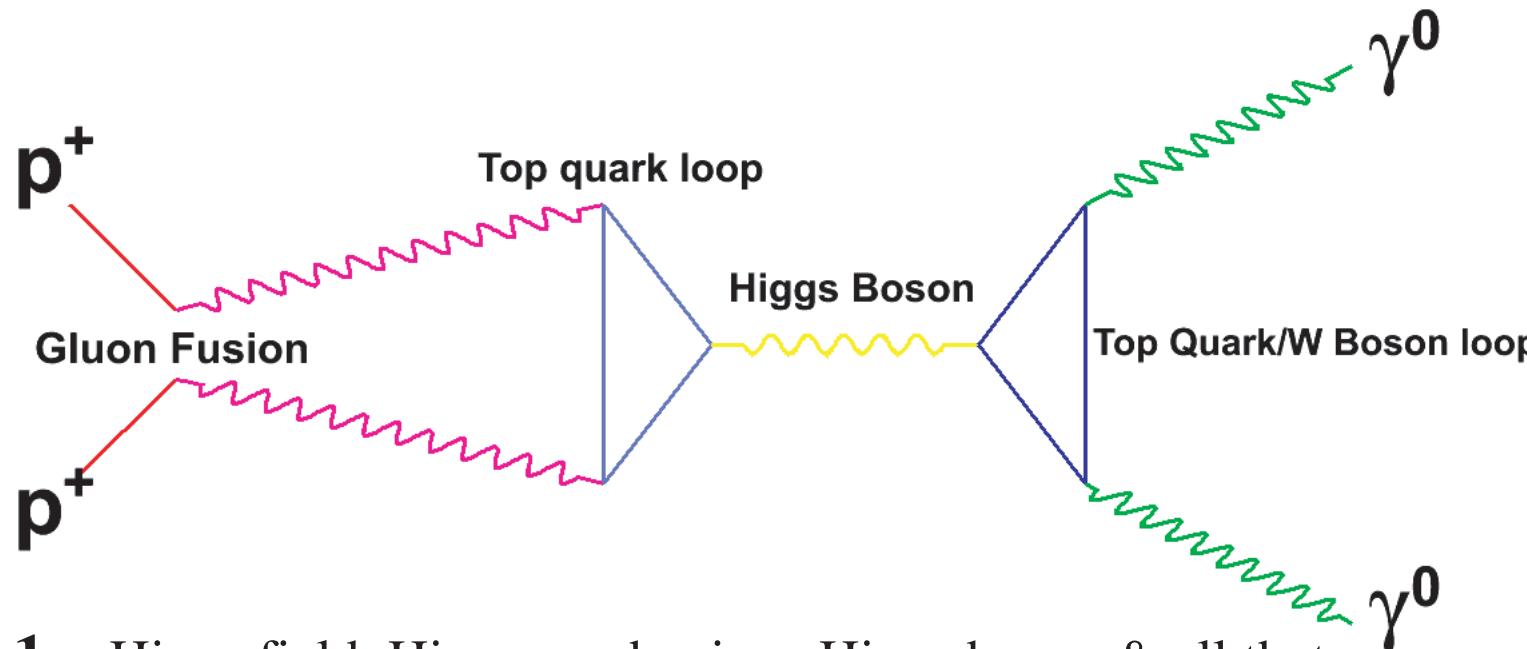
Higgs bosons in particle physics & condensed matter

$\hat{a}^+ \hat{a}$ Aalto University

G. Volovik

Landau Institute

RUSSIAN ACADEMY OF SCIENCES
L. D Landau
INSTITUTE FOR
THEORETICAL
PHYSICS



1. Higgs field, Higgs mechanism, Higgs boson & all that
2. ${}^3\text{He-B}$ & Nambu conjecture on masses of fermions & bosons
3. from ${}^3\text{He-B}$ to Standard Model: application of Nambu sum rule
4. do we see extra Higgs bosons ?
5. Higgs bosons from quench (cosmology & superconductors)
6. Majorana & Majorino



in collaboration with Mikhail Zubkov, ITEP, Moscow

Annual Gods Meeting, Matsue

Torsion & spinning strings, torsion instanton

Fermion zero modes on strings & walls

Antigravitating (negative-mass) string

Gravitational Aharonov-Bohm effect

Domain wall terminating on string

String terminating on domain wall

Monopoles on string & Boojums

Witten superconducting string

Soft core string, Q-balls

Z & W strings

skyrmions

Alice string

Pion string

Kibble mechanism

Dark matter detector

Primordial magnetic field

Baryogenesis by textures

Cosmological & Newton constants

dark energy

dark matter

Effective gravity

Bi-metric & conformal gravity

Graviton, dilaton

Spin connection

Rotating vacuum

Vacuum dynamics

conformal anomaly

ergoregion, event horizon

Hawking & Unruh effects

black hole entropy

Vacuum instability

Superfluidity of neutron star

vortices, glitches

Nambu-Jona-Lasinio

shear flow instability

Vaks-Larkin

Color superfluidity

Savvidi vacuum

Quark confinement, QCD cosmology

Intrinsic orbital momentum of quark matter

Emergence & effective theories

Weyl, Majorana & Dirac fermions

Vacuum polarization, screening - antiscreening, running coupling

Symmetry breaking (anisotropy of vacuum)

Parity violation -- chiral fermions

Vacuum instability in strong fields, pair production

Casimir force, quantum friction

Fermionic charge of vacuum

Higgs fields, gauge bosons, Higgs bosons title 2013

Momentum-space topology

Hierarchy problem, Supersymmetry

Neutrino oscillations

Chiral anomaly & axions

Spin & isospin, String theory

CPT-violation, GUT

Gap nodes

Low -T scaling

mixed state

Broken time reversal

1/2-vortex, vortex dynamics

Films: FQHE,
Statistics & charge of
skyrmions & vortices

Edge states;
flat band, Fermi arc
spintronics

1D fermions in vortex core
Critical fluctuations

Mixture of condensates

Vector & spinor condensates

BEC of quasiparticles,

magnon BEC & laser

meron, skyrmion, 1/2 vortex

General; relativistic;

spin superfluidity

multi-fluid

rotating superfluid

Shear flow instability

Magnetohydrodynamic

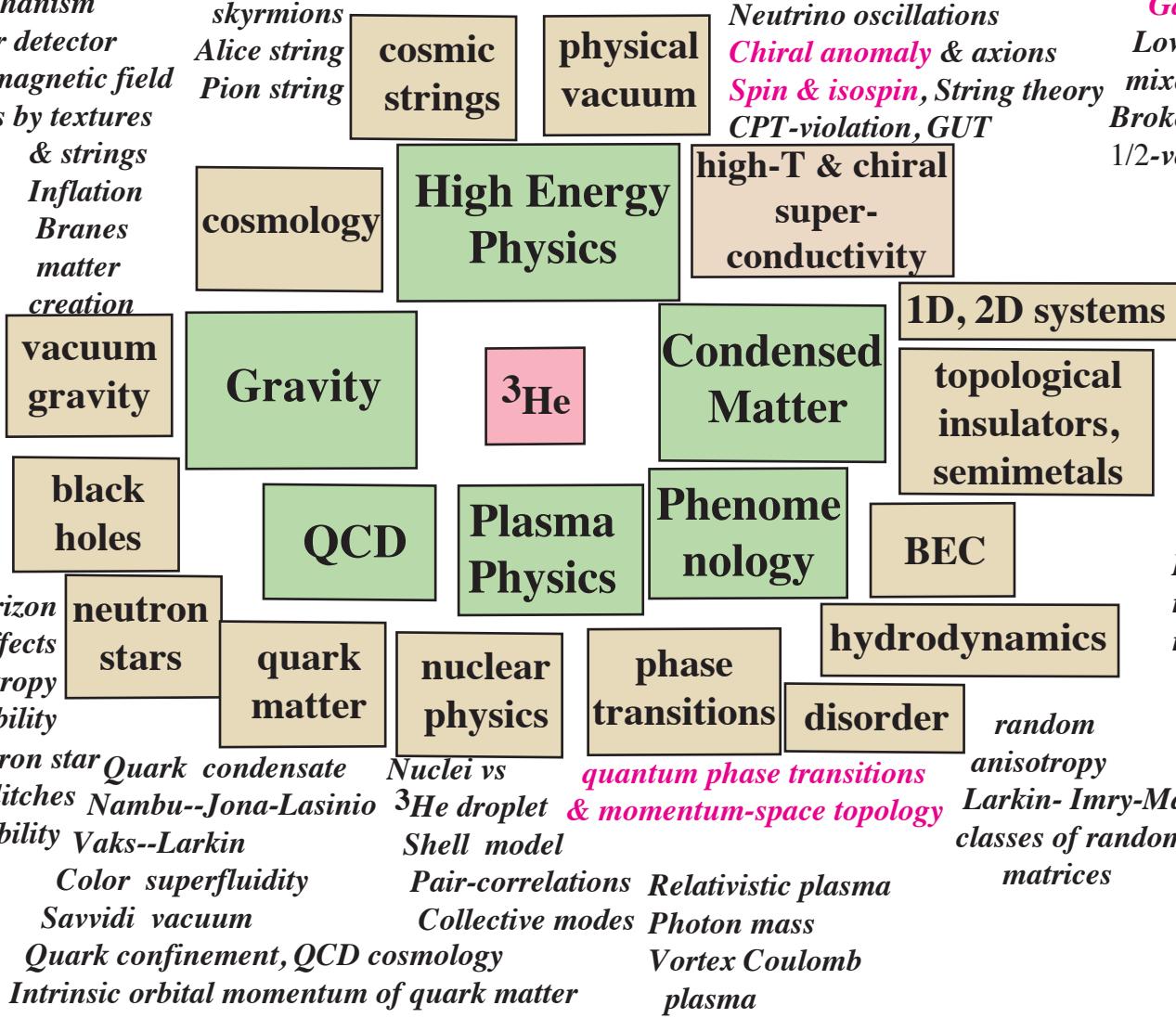
Turbulence of vortex lines

propagating vortex front

velocity independent

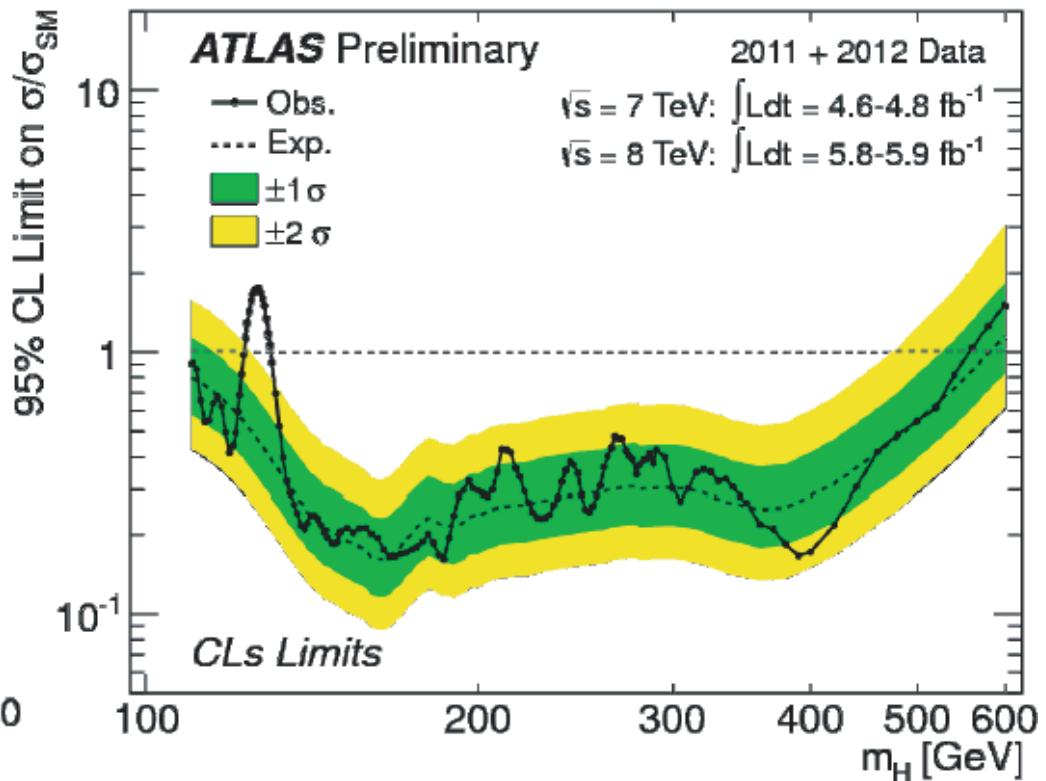
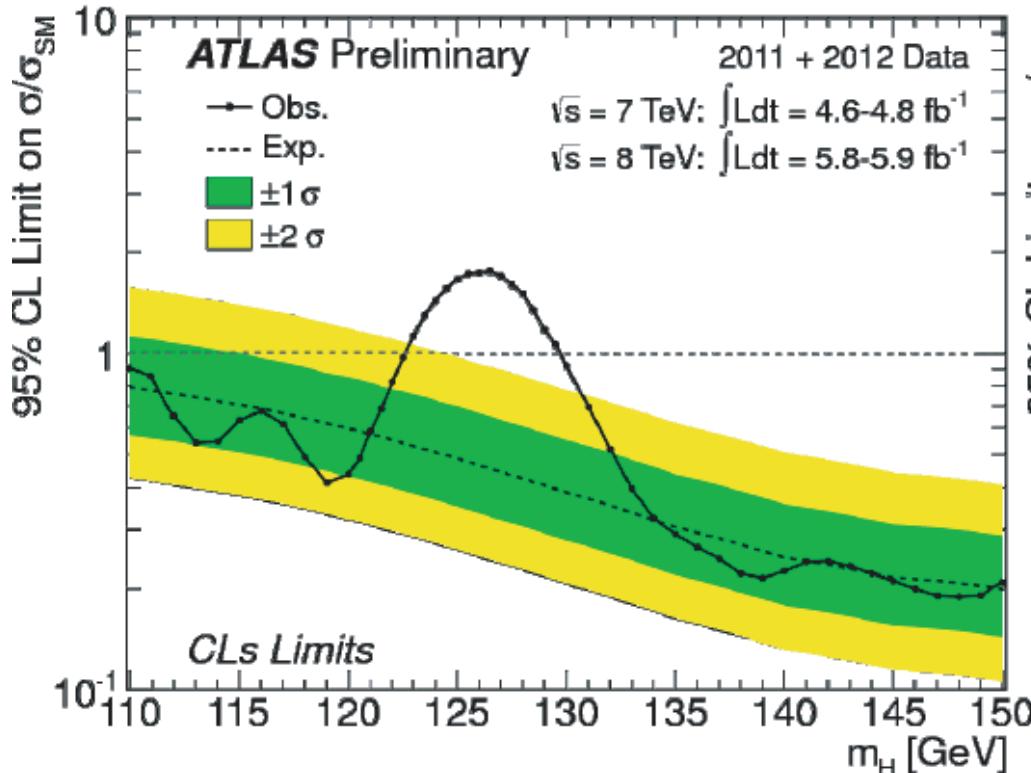
Reynolds number

3He Universe

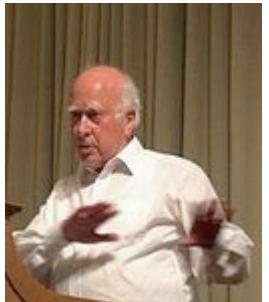


Updated combined SM Higgs exclusion

- Combining updated 2011+2012 $\gamma\gamma$ and 4-lepton analysis with others as before



- Expected exclusion with this dataset from 110-582 GeV
 - Observed exclusion from 110-122.6 GeV and 129.7-558 GeV
 - 111.7-121.8 GeV and 130.7-523 GeV excluded at 99% CL
 - Region around 126 GeV cannot be excluded



author of Higgs mechanism & massive Higgs bosons

co-authors of Higgs mechanism



Kibble

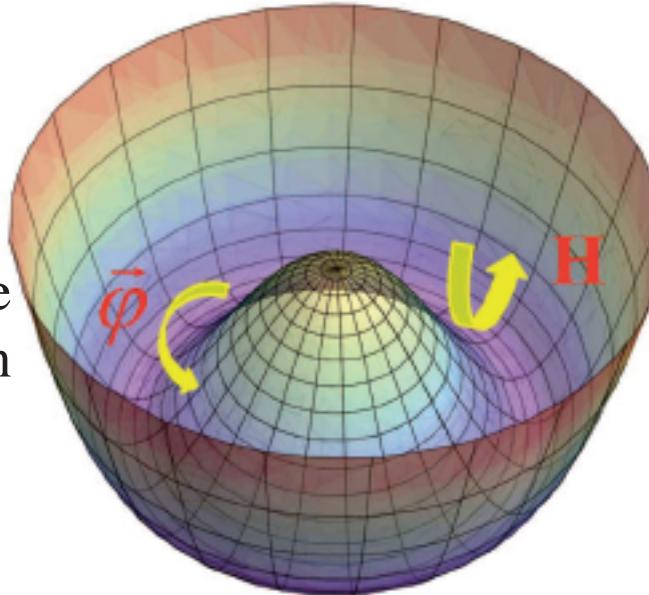
Guralnik

Hagen

Englert

Brout

Symmetry breaking, Higgs mechanism & massive particles



Standard Model

no NG bosons !

NG bosons

massive particles

Higgs boson(s)

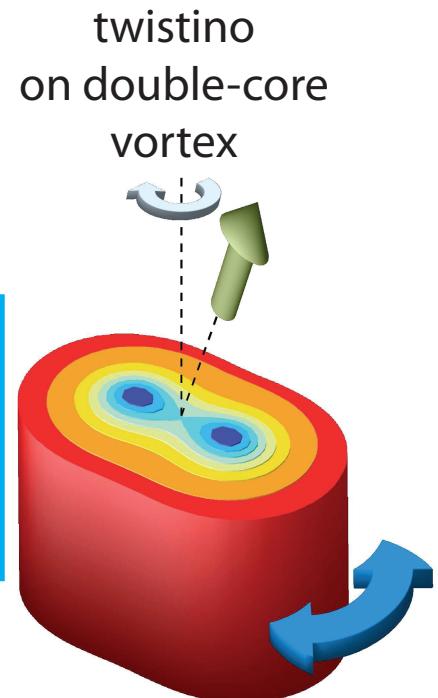
order parameter modes

W & Z gauge bosons

Meissner effect (expulsion of magnetic field)

quarks & leptons

gapped quasiparticles



Nambu-Goldstone (NG) bosons

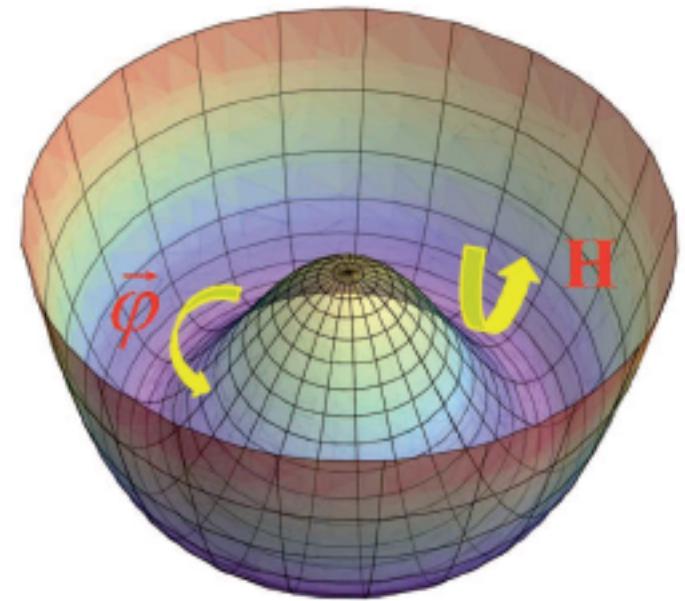
Goldstone's theorem:

spontaneous breakdown of a global continuous symmetry
leads to massless particle – Nambu-Goldstone boson

Y. Nambu, Phys. Rev. Lett. **4** (1960) 380

J. Goldstone, Nuovo Cim. **19** (1961) 154

J. Goldstone, A. Salam, S. Weinberg, Phys. Rev. **127** (1962) 965



cond-mat order parameter = Higgs field

$$\Psi = |\Psi| e^{i\varphi}$$

$$\Psi = \begin{pmatrix} \Psi_1 \\ \Psi_2 \end{pmatrix}$$

cond-mat NG bosons = gapless modes of order parameter:
phonons, spin waves, orbital waves

no NG bosons had been observed in particle physics

Why ?

Higgs mechanism

spontaneous breaking of gauge symmetry does not require Goldstone bosons,
their degrees of freedom deliver the longitudinal polarization modes of gauge bosons,
gauge bosons become massive (Z-boson, W-boson)

P. Anderson, Plasmons, gauge invariance and mass, Phys. Rev. **130** (1963) 439

F. Englert, R. Brout, Phys. Rev. Lett. **13** (1964) 321

P. W. Higgs, Phys. Lett. **12** (1964) 132 and Phys. Rev. Lett. **13** (1964) 508

G. S. Guralnik, C. R. Hagen, T. W. B. Kibble, Phys. Rev. Lett. **13** (1964) 585

cond-mat analog: Cooper pairing and Meissner effect in superconductors

Higgs, PRL **13** (1964) 508:

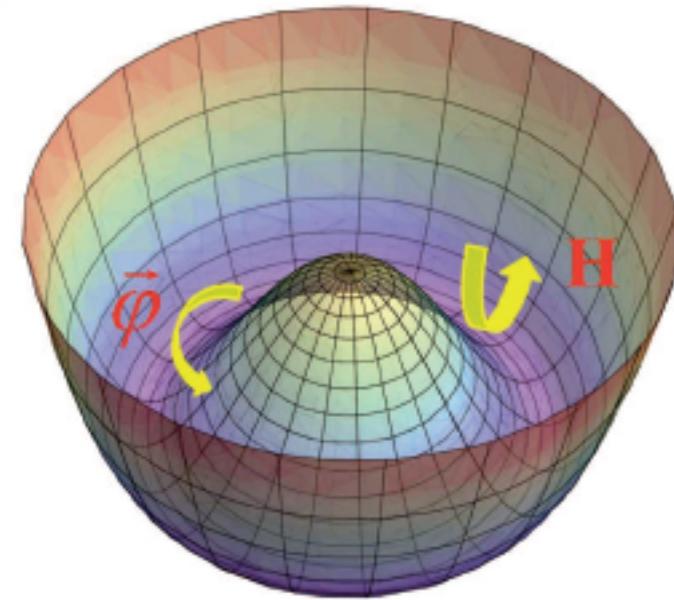
“It is to be expected that this feature will appear also in theories
in which the symmetry-breaking scalar fields are not elementary dynamic variables
but bilinear combinations of Fermi fields.”

“In the theory of superconductivity the scalar fields are associated with fermion pairs;
the doubly charged excitation responsible for the quantization of magnetic flux
is then the surviving member of a U(1) doublet.”

How many Higgs bosons ?

amplitude H-modes of Higgs field

P. W. Higgs, PRL **13** (1964) 508:



“The model of the most immediate interest is that in which the scalar fields form an **octet** under **SU(3)**... There are **2** massive scalar bosons ... (**2 Higgs bosons**) ... ; the remaining **6** components of the scalar octet combine with the corresponding components of the gauge-field octet to describe massive vector bosons (**6 massive gauge bosons**).”

Vdovin (1963)

Bosons in superfluid $^3\text{He-B}$: collective modes of order parameter

9-plet $A_{\alpha i}$ under $\text{SO}(3) \times \text{SO}(3) \times \text{U}(1)$

14 Higgs bosons + 4 NG modes (4 massive gauge bosons in gauged $^3\text{He-B}$)

Masses of elementary particles

elementary particles acquire mass due to interaction with Higgs field

cond-mat mass of elementary particles:

Bogoliubov- de Gennes particles with gap in superconductors and $^3\text{He-B}$

Ginzburg-Landau

vs

Standard Model

Order parameter:
complex scalar field

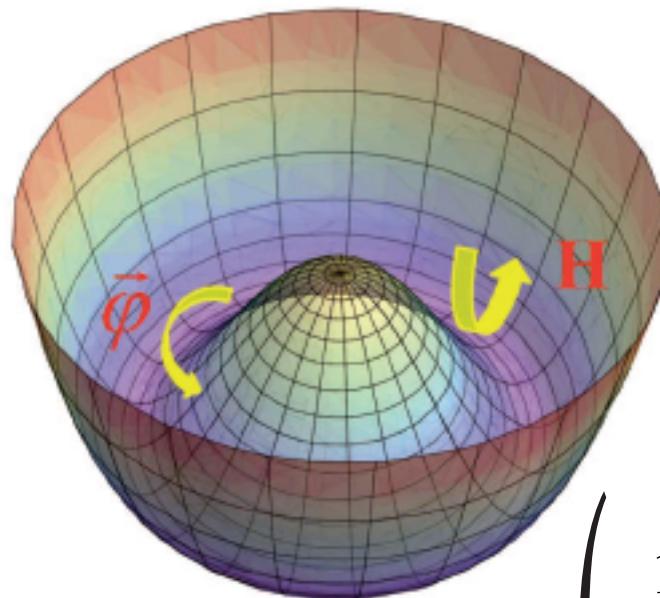
$$\Psi = |\Psi| e^{i\varphi}$$

$2 = 1+1$ collective modes:

$$\Psi = 1 + (H + i\varphi)$$

1 NG mode (phase φ mode)

1 amplitude H mode (Higgs boson)



Higgs field: complex spinor

$$\Psi = \begin{pmatrix} \Psi_1 \\ \Psi_2 \end{pmatrix}$$

$4 = 3+1$ collective modes

$$\Psi = \begin{pmatrix} 1 \\ 0 \end{pmatrix} + H \begin{pmatrix} 1 \\ 0 \end{pmatrix} + \delta\Psi$$

3 rotations

3 NG modes

1 amplitude H mode (Higgs boson H)

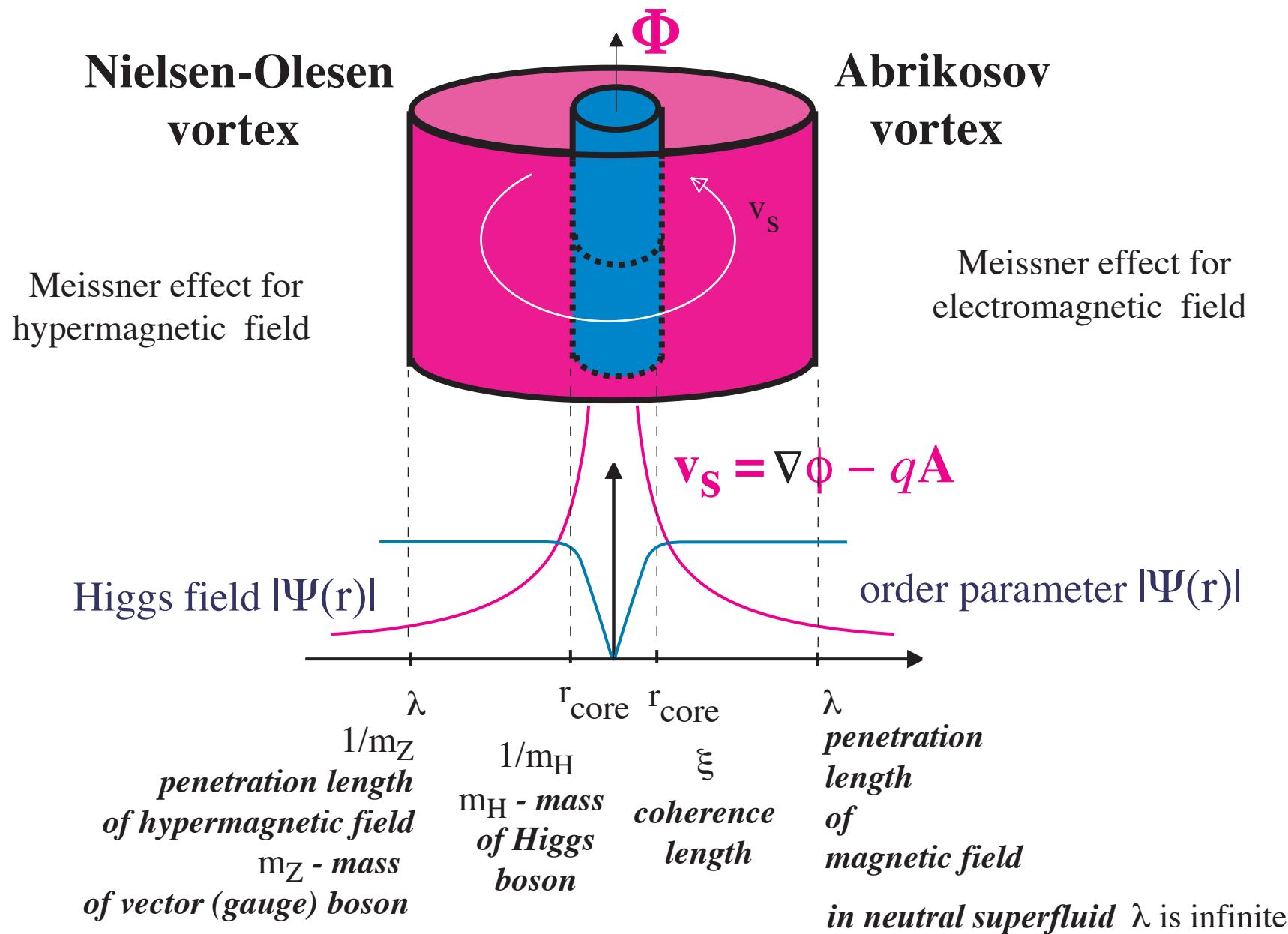
Standard Model requires only one Higgs boson

Meissner effect:

mass of gauge vector boson =
inverse penetration length

3 NG modes are absorbed into longitudinal
modes of 3 massive vector bosons
(charged W^+ , W^- and neutral Z)
this is called **Higgs mechanism**

Higgs mechanism in Standard Model & superconductor



$$E = |\nabla\Psi - iqA\Psi|^2 = |\Psi|^2(\nabla\phi - qA)^2 = \lambda^{-2}(\nabla\phi - qA)^2 = \rho_S v_S^2$$

from Standard Model Ginzburg-Landau to Standard Model BEC

nobody now believes that SM is complete theory,
all believe SM is effective low energy theory

composite Higgs is more natural than fundamental

Higgs, PRL **13** (1964) 508:

“... the symmetry-breaking scalar fields
are not elementary dynamic variables
but *bilinear combinations of Fermi fields.*”

Higgs field can be composite object as Cooper pair in superconductors

$$\Psi = \langle \bar{\chi}_L \chi_R \rangle$$

Standard Model as BCS theory in semimetal

BCS Cooper pairing

dynamical mixing of left & right particles

Order parameter as composite object

Higgs field as composite object

$$\Psi = \langle \chi \chi \rangle$$

$$\Psi = \langle \bar{\chi}_L \chi_R \rangle$$

phase & amplitude modes

Higgs bosons

$$\Psi = \langle \chi \chi \rangle_{\text{vac}} + \delta \Psi$$

$$\Psi = \langle \bar{\chi}_L \chi_R \rangle_{\text{vac}} + \delta \Psi$$

Ψ mixes electron and hole

Ψ mixes left and right Weyl particles

Hamiltonian for Bogoliubov-Nambu quasiparticles

Hamiltonian for Dirac particles

$$H = \begin{pmatrix} \epsilon(\mathbf{p}) & \Psi \\ \Psi^* & -\epsilon(\mathbf{p}) \end{pmatrix}$$

order parameter induces gap

$$E^2(\mathbf{p}) = \epsilon^2(\mathbf{p}) + |\Psi|^2$$

$$H = \begin{pmatrix} c\sigma \cdot \mathbf{p} & \Psi \\ \Psi^+ & -c\sigma \cdot \mathbf{p} \end{pmatrix}$$

Higgs field induces Dirac mass $M = |\Psi|$

$$E^2(\mathbf{p}) = c^2 p^2 + |\Psi|^2$$

generalization of Standard Model using hints from superfluid ^3He

superfluid ^3He

$$\text{order parameter } A_{\alpha j} = \langle \chi \sigma_\alpha \nabla_j \chi \rangle = \langle \chi \hat{O} \chi \rangle$$

in superfluid ^3He

$$\hat{O} = \sigma_\alpha \nabla_j$$

\hat{O} is in general operator, which may include generators of different symmetry groups

general composite Higgs field

$$\text{Higgs field } \Psi = \langle \bar{\chi} \hat{O} \chi \rangle$$

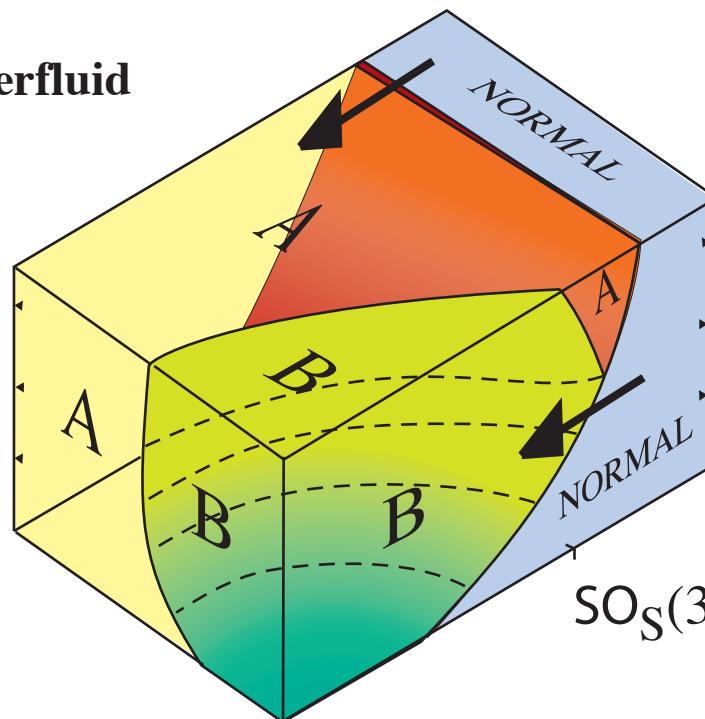
GUT in Standard Model
symmetry breaking phase transitions

$$SO(10) \rightarrow SU(3) \times SU_L(2) \times U(1) \rightarrow SU(3) \times U_Q(1)$$

symmetry breaking phase transitions
in superfluid ^3He

$$SO_S(3) \times SO_L(3) \times U(1) \rightarrow SO_S(2) \times U_Q(1) \quad Q \text{ is analog of electric charge}$$

$^3\text{He-A}$
topological chiral Weyl superfluid
two Dirac cones



$^3\text{He-B}$
time-reversal symmetric
topological superfluid

$$SO_S(3) \times SO_L(3) \times U(1) \rightarrow SO_J(3)$$

multiple Higgs bosons from superfluid ^3He

superfluid ^3He

$$\text{order parameter } A_{\alpha j} = \langle \chi \sigma_\alpha \nabla_j \chi \rangle$$

pair breaking & squashing modes

$$A_{\alpha j} = \delta_{\alpha j} + u_{\alpha j} + i v_{\alpha j}$$

\hat{O} is operator, which may include generators of different symmetry groups

Bogoliubov-Nambu quasiparticles

$$H = \begin{pmatrix} \epsilon(\mathbf{p}) & A_{\alpha j} \sigma_\alpha \cdot p_j \\ A_{\alpha j}^* \sigma_\alpha \cdot p_j & -\epsilon(\mathbf{p}) \end{pmatrix}$$

composite Higgs

$$\text{Higgs field } \Psi = \langle \bar{\chi} \hat{O} \chi \rangle$$

multiple Higgs bosons

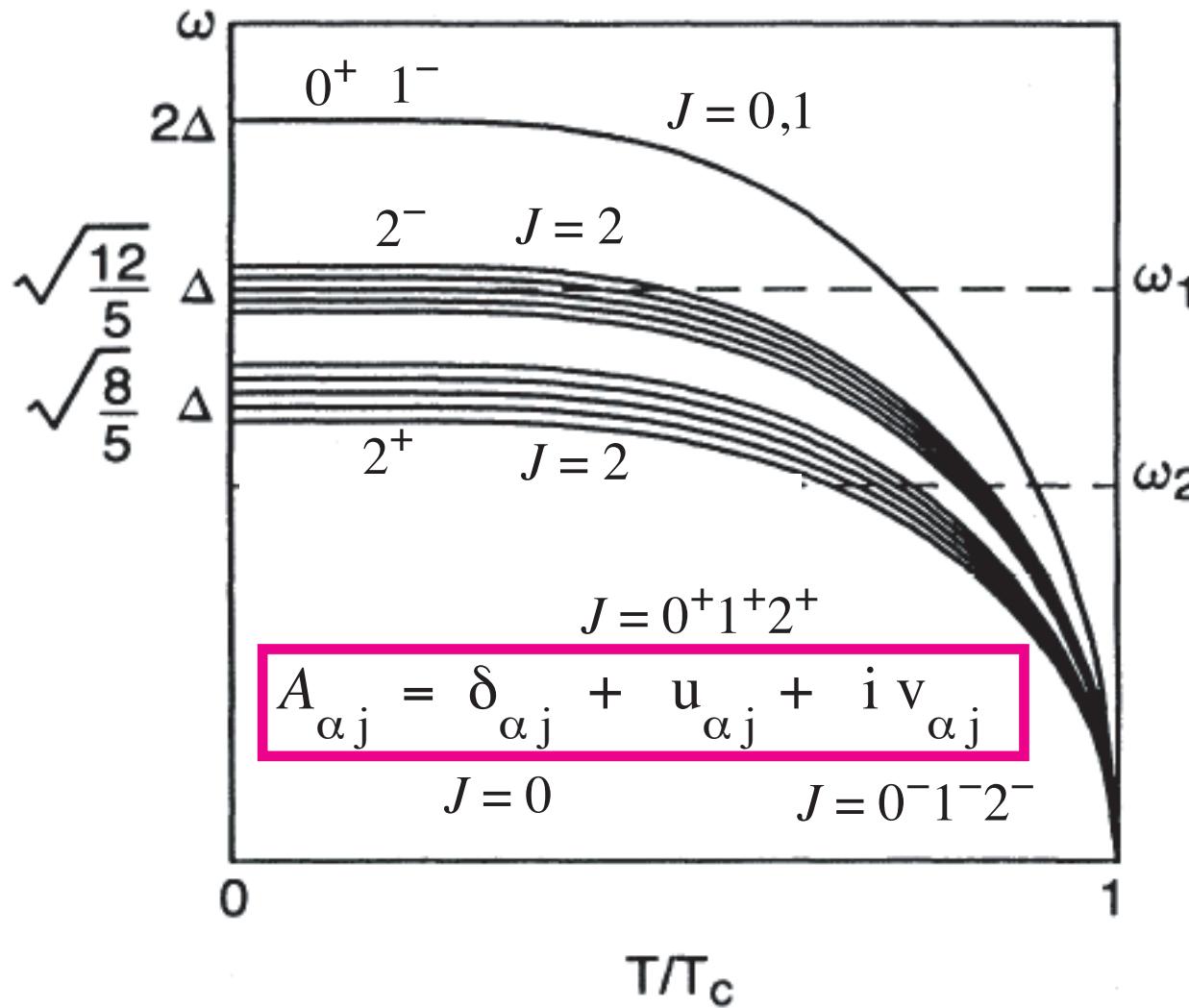
$$\Psi = \langle \bar{\chi} \hat{O} \chi \rangle_{\text{vac}} + \delta \Psi$$

massive quarks & leptons

$$H = \begin{pmatrix} c\sigma \cdot \mathbf{p} & & & \Psi \\ & c\sigma \cdot \mathbf{p} & & \\ & & -c\sigma \cdot \mathbf{p} & \\ \Psi^+ & & & -c\sigma \cdot \mathbf{p} \end{pmatrix}$$

4 Goldstone & 14 Higgs modes in $^3\text{He-B}$

$\text{SO}_S(3) \times \text{SO}_L(3) \times \text{U}(1) \longrightarrow \text{SO}_J(3)$



4 Goldstone modes

$$E_{0-} = E_{1+} = 0$$

14 Higgs bosons

$$E_{1-} = E_{0+} = 2\Delta$$

$$E_{2-} = (12/5)^{1/2}\Delta$$

$$E_{2+} = (8/5)^{1/2}\Delta$$

9 pairs of Nambu partners

$$E_{2-}^2 + E_{2+}^2 = 4\Delta^2$$

$$E_{0-} + E_{0+} = 4\Delta^2$$

$$E_{1-} + E_{1+} = 4\Delta^2$$

FIG. 3. A schematic plot in ω vs T space for pair breaking, the squashing mode, and the real squashing modes. The Zeeman splitting (not to scale) of the collective modes in an applied magnetic field is shown in the plot. The dashed lines labeled ω_1 and ω_2 correspond to two sound frequencies.

Nambu sum rule: from 3He-B to Standard Model

relation between energies E_B & E_F of bosonic & fermionic excitations
in BCS type theories

$$E_{B1}^2 + E_{B2}^2 = 4 E_F^2$$

3He-B $E_{B1} = (8/5)^{1/2}\Delta$ $E_{B2} = (12/5)^{1/2}\Delta$ $E_F = \Delta$
real squashing mode squashing mode gap in quasiparticle spectrum

Application of Nambu rule to masses of Higgs fields and top quarks

$$m_{H1}^2 + m_{H2}^2 = 4 m_t^2$$

125 GeV 325 GeV 174 GeV

Sum rule for ${}^3\text{He-B}$: symmetry consideration

$$m_Q^2 = 2m_f^2 (1 +/\!-\! \eta^{(Q)})$$

$$m_{Q+}^2 + m_{Q-}^2 = 4 m_f^2$$

parameter dictated by symmetry

$$\eta^{(Q)} = \text{Tr} (\text{VO}^{(Q)} \text{VO}^{(Q)})$$

${}^3\text{He-B}$: quantum number $Q = J$ (total angular momentum), $m_f = \Delta$ gap

$$E_J^2 = 2 \Delta^2 (1 +/\!-\! \eta^{(J)})$$

$$J=2, \eta^{(J)} = 1/5$$

$$E_{2-} = (8/5)^{1/2} \Delta$$

real squashing mode

$$E_{2+} = (12/5)^{1/2} \Delta$$

squashing mode

$$J=0, \eta^{(J)} = 1$$

$$E_{1-} = 0$$

Goldstone
(sound mode)

$$E_{1+} = 2 \Delta$$

pair breaking mode

$$J=1, \eta^{(J)} = 1$$

$$E_{0-} = 0$$

3 Goldstone modes
(spin waves)

$$E_{0+} = 2 \Delta$$

pair breaking modes

which fermion is responsible?

Standard Model

Three generations
of matter (fermions)

	I	II	III	
mass →	2.4 MeV/c ²	1.27 GeV/c ²	171.2 GeV/c ²	
charge →	$\frac{2}{3}$	$\frac{2}{3}$	$\frac{2}{3}$	
spin →	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	
name →	u up	c charm	t top	γ photon
Quarks				
	4.8 MeV/c ² $-\frac{1}{3}$ $\frac{1}{2}$ d down	104 MeV/c ² $-\frac{1}{3}$ $\frac{1}{2}$ s strange	4.2 GeV/c ² $-\frac{1}{3}$ $\frac{1}{2}$ b bottom	0 0 1 g gluon
Leptons				
	<2.2 eV/c ² 0 $\frac{1}{2}$ ν _e electron neutrino	<0.17 MeV/c ² 0 $\frac{1}{2}$ ν _μ muon neutrino	<15.5 MeV/c ² 0 $\frac{1}{2}$ ν _τ tau neutrino	91.2 GeV/c ² 0 1 Z ⁰ Z boson
	0.511 MeV/c ² -1 $\frac{1}{2}$ e electron	105.7 MeV/c ² -1 $\frac{1}{2}$ μ muon	1.777 GeV/c ² -1 $\frac{1}{2}$ τ tau	80.4 GeV/c ² ± 1 1 W [±] W boson

if pairing occurs in one channel
fermion in this channel
will have largest gap (mass)

Gauge bosons

Nambu sum rule: from 3He-B to Standard Model

relation between energies E_B & E_F of bosonic & fermionic excitations
in BCS type theories

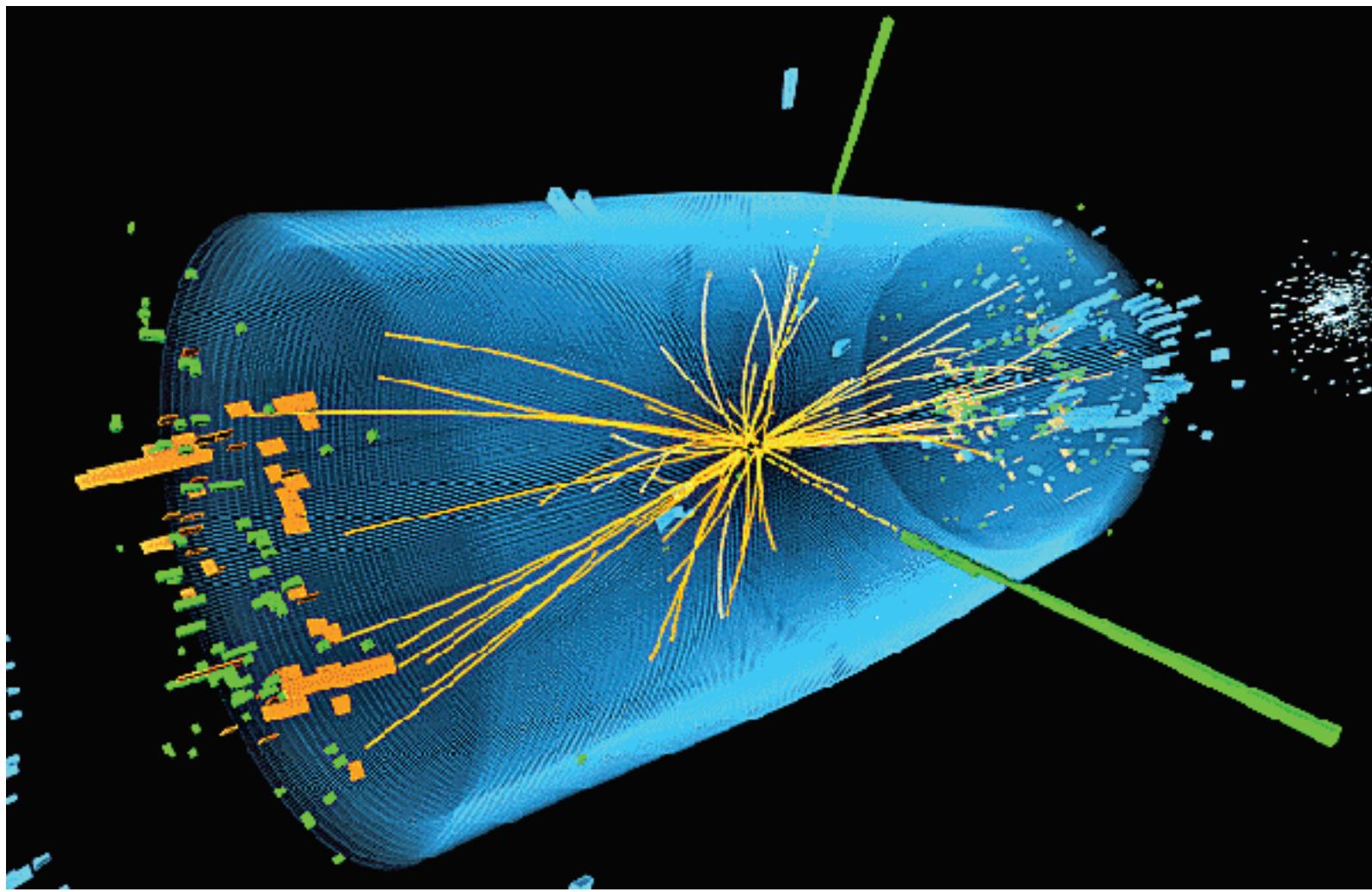
$$E_{B1}^2 + E_{B2}^2 = 4 E_F^2$$

3He-B $E_{B1} = (8/5)^{1/2}\Delta$ $E_{B2} = (12/5)^{1/2}\Delta$ $E_F = \Delta$
real squashing mode squashing mode gap in quasiparticle spectrum

Application of Nambu rule to masses of Higgs fields and top quarks

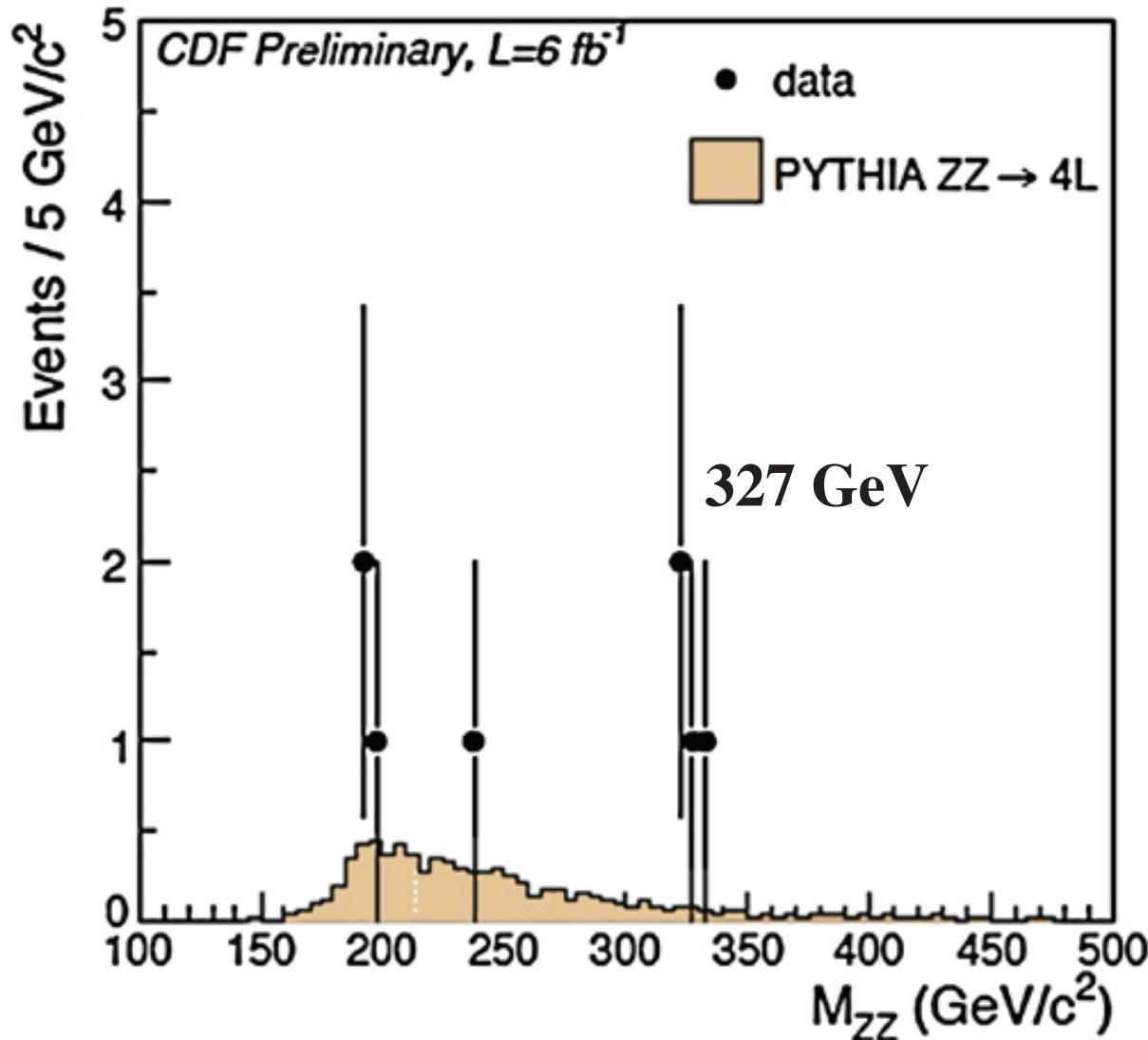
$$m_{H1}^2 + m_{H2}^2 = 4 m_t^2$$

125 GeV 325 GeV 174 GeV



A narrow scalar resonance at 325 GeV?

Krzysztof A. Meissner & Hermann Nicolai
Phys. Lett. B **718** (2013) 943



“We propose to identify the excess of events with four charged leptons at $E \sim 325$ GeV seen by the CDF Collaboration (2012) [1] & CMS Collaboration (2012) [2] with a new ‘sterile’ scalar particle characterized by a very narrow resonance of the same height and branching ratios as the Standard Model Higgs boson, as predicted in the framework of the so-called Conformal Standard Model (K.A. Meissner & H. Nicolai 2007) [3].”

Fig. 1. The four lepton events reported by the CDF Collaboration [1].

hints of 325 GeV Higgs boson in earlier experiments

* CDF/PUB/EXOTICS/PUBLIC/10603 July 17, 2011
Search for High-Mass Resonances Decaying
into ZZ in $p\bar{p}$ Collisions at $s^{1/2} = 1.96\text{TeV}$

"The invariant masses of four events
are clustered around 325 GeV"

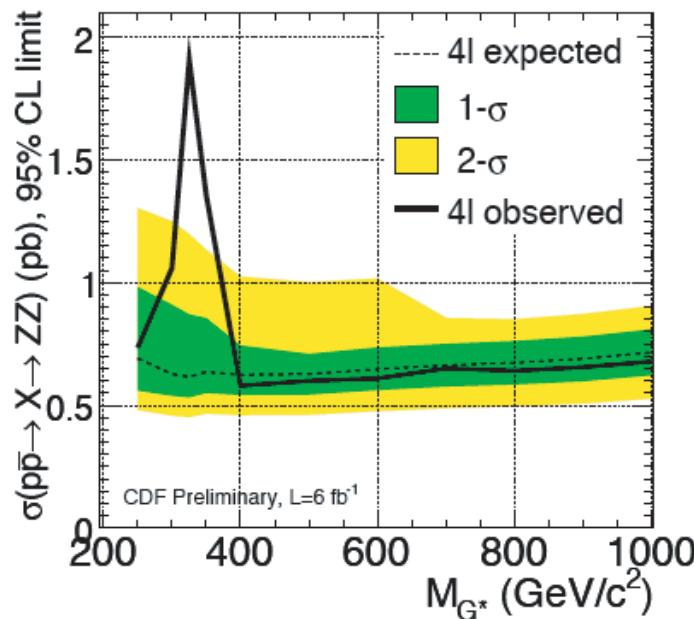
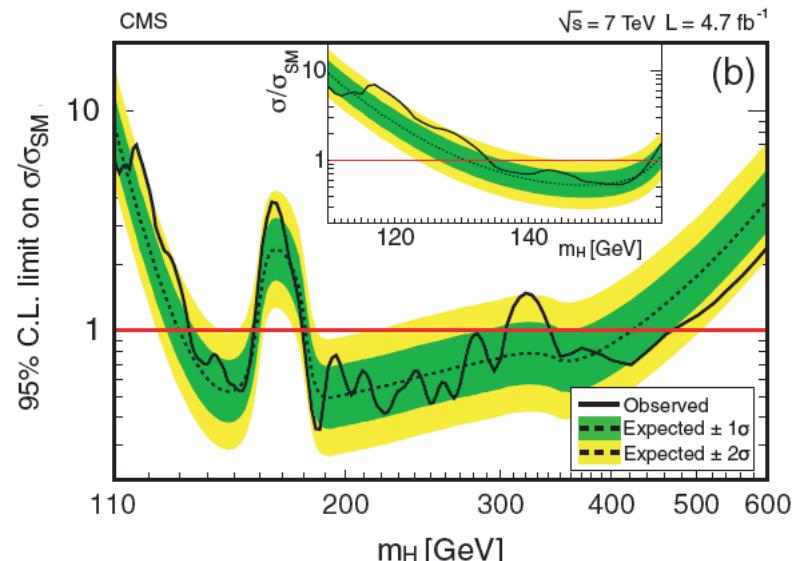


FIG. 13: Expected and observed 95% CL limits on $\sigma(p\bar{p} \rightarrow X \rightarrow ZZ)$ from the $ZZ \rightarrow \ell^+\ell^-\ell^+\ell^-$ channel; the four events with $M_{ZZ}=327 \text{ GeV}/c^2$ result in a deviation from the expected limit.

* Search for the Standard Model Higgs Boson
in CMS Collaboration @ LHS
(Compact Muon Solenoid))

PRL 108, 111804 (2012)

"Small excesses of events are observed around
masses of 119, 126 & 320 GeV"



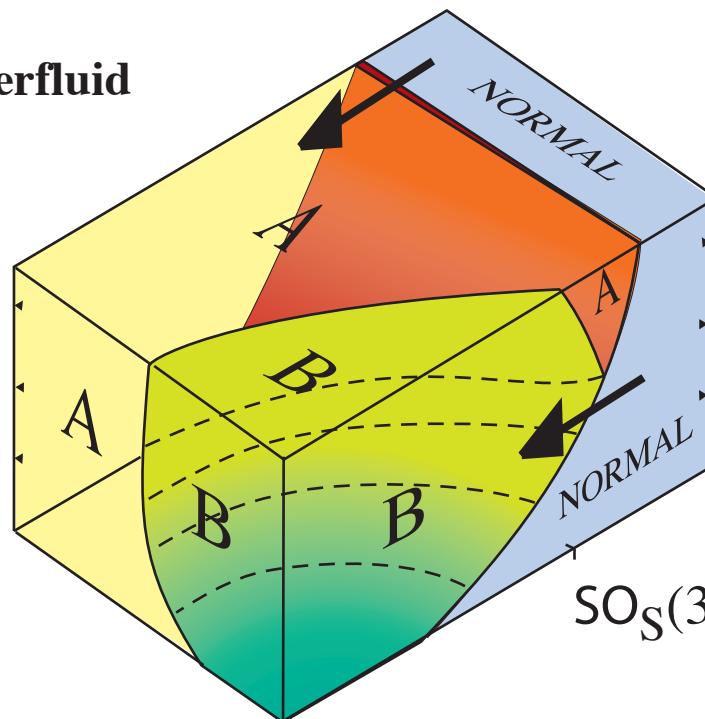
GUT in Standard Model
symmetry breaking phase transitions

$$SO(10) \rightarrow SU(3) \times SU_L(2) \times U(1) \rightarrow SU(3) \times U_Q(1)$$

symmetry breaking phase transitions
in superfluid ^3He

$$SO_S(3) \times SO_L(3) \times U(1) \rightarrow SO_S(2) \times U_Q(1) \quad Q \text{ is analog of electric charge}$$

$^3\text{He-A}$
topological chiral Weyl superfluid
two Dirac cones



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time-reversal symmetric
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$$SO_S(3) \times SO_L(3) \times U(1) \rightarrow SO_J(3)$$

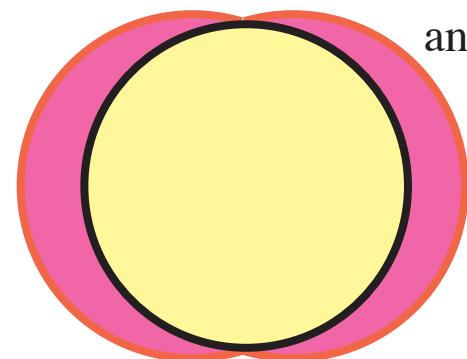
9 Goldstone + 9 Higgs modes in bulk $^3\text{He-A}$

$$\text{SO}_S(3) \times \text{SO}_L(3) \times \text{U}(1) \longrightarrow \text{SO}_S(2) \times \text{U}_Q(1)$$

$\text{U}_Q(1)$ - analog of electromagnetic group

quantum number Q - analog of electric charge

Modes	Variables	In the absence of dipole interaction and magnetic field			
		Q	p_z	s_z	p_t
Sound $E = 0$	$u_{23} - v_{13}$	0	-1	0	+1
Spin waves $E = 0$	$u_{11} + v_{21},$ $u_{12} + v_{22}$	0	+1	± 1	-
Orbital modes	u_{33}, v_{33}	± 1	-	0	-1
Spin-orbit modes $E = 0$	u_{31}, v_{31} u_{32}, v_{32}	± 1	-	± 1	-
$E^2 = 8/3 \Delta_0^2$	$u_{13} + v_{23}$	0	+1	0	+1
	$u_{22} - v_{12},$ $v_{11} - u_{21}$	0	-1	± 1	-
	$u_{23} + v_{13},$ $u_{13} - v_{23}$	± 2	-	0	+1
"Clapping" modes	$u_{11} - v_{21},$ $u_{21} + v_{11},$ $u_{12} - v_{22},$ $u_{22} + v_{12}$	± 2	-	± 1	-



anisotropic gap with nodes

$$\Delta(\theta) = \Delta_0 \sin \theta$$

$$\langle \Delta^2 \rangle = 2/3 \Delta_0^2$$

3 Goldstones + 9 Higgs
form 6 Nambu pairs

$$E_{0-}^2 + E_{0+}^2 = 4 \langle \Delta^2 \rangle$$

$$E_{0-}^2 + E_{0+}^2 = 4 \langle \Delta^2 \rangle$$

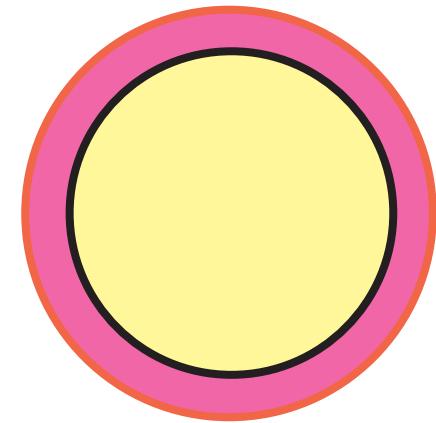
$$E_{-2}^2 + E_{+2}^2 = 4 \langle \Delta^2 \rangle$$

6 Goldstones
(2 "photons" + 4 from hidden symmetry)
have no Nambu partners

3 Goldstone + 9 Higgs modes in $^3\text{He-A}$ film

$$\text{SO}_S(3) \times \text{SO}_L(2) \times \text{U}(1) \longrightarrow \text{SO}_S(2) \times \text{U}_Q(1)$$

$$m_{Q+}^2 + m_{Q-}^2 = 4 m_f^2$$



isotropic gap Δ in 2D

$^3\text{He-A}$ film: quantum number Q is analog of electric charge, $m_f = \Delta$ gap

$Q=0$

$$E_{0-} = 0$$

Goldstone
(sound mode)

$$E_{0+} = 2 \Delta$$

pair breaking mode

$$E_{0+}^2 + E_{0-}^2 = 4 \Delta^2$$

$|Q|=1$

$$E_{-1} = 2^{1/2} \Delta$$

clapping modes
analogs of charged Higgs bosons

$$E_{+1} = 2^{1/2} \Delta$$

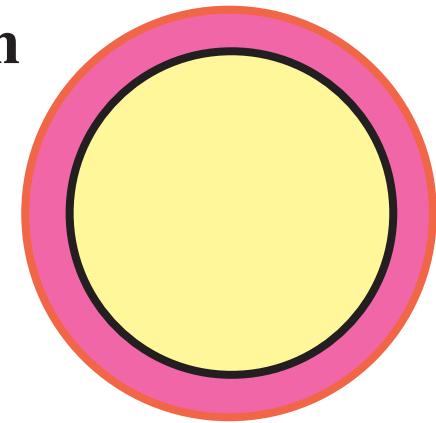
$$E_{-2}^2 + E_{+2}^2 = 4 \Delta^2$$

Nambu sum rule for ^3He -A film: symmetry consideration

$$m_Q^2 = 2m_f^2 (1 +/\!-\! \eta^{(Q)})$$

$$\eta^{(Q)} = \text{Tr} (\text{VO}^{(Q)} \text{V} \text{O}^{(Q)})$$

$$m_{Q+}^2 + m_{Q-}^2 = 4 m_f^2$$



isotropic gap Δ in 2D

^3He -A film: quantum number Q is analog of electric charge, $m_f = \Delta$ gap

$$E_Q^2 = 2 \Delta^2 (1 +/\!-\! \eta^{(Q)})$$

$$Q=0, \eta^{(Q)} = 1$$

$$E_{0-} = 0$$

Goldstone
(sound mode)

$$E_{0+} = 2 \Delta$$

pair breaking mode

$$E_{0+}^2 + E_{0-}^2 = 4 \Delta^2$$

$$Q=1, \eta^{(Q)} = 0$$

$$E_{-1} = 2^{1/2} \Delta$$

clapping modes
analogs of charged Higgs bosons

$$E_{+1} = 2^{1/2} \Delta$$

$$E_{-2}^2 + E_{+2}^2 = 4 \Delta^2$$

Nambu sum rule: from ${}^3\text{He}$ -A to Standard Model

relation between energies E_B & E_F of bosonic & fermionic excitations in BCS type theories

$$E_{B1}^2 + E_{B2}^2 = 4 E_F^2$$

Application of Nambu rule to masses of charged Higgs fields and top quarks

$$\begin{array}{ccc}
 \boxed{m_{H+}^2} & + & \boxed{m_{H-}^2} = \boxed{4 m_t^2} \\
 245 \text{ GeV} & & 245 \text{ GeV} & & 174 \text{ GeV} \\
 \text{charged Higgs} & & \text{charged Higgs} & & \text{top quark}
 \end{array}$$

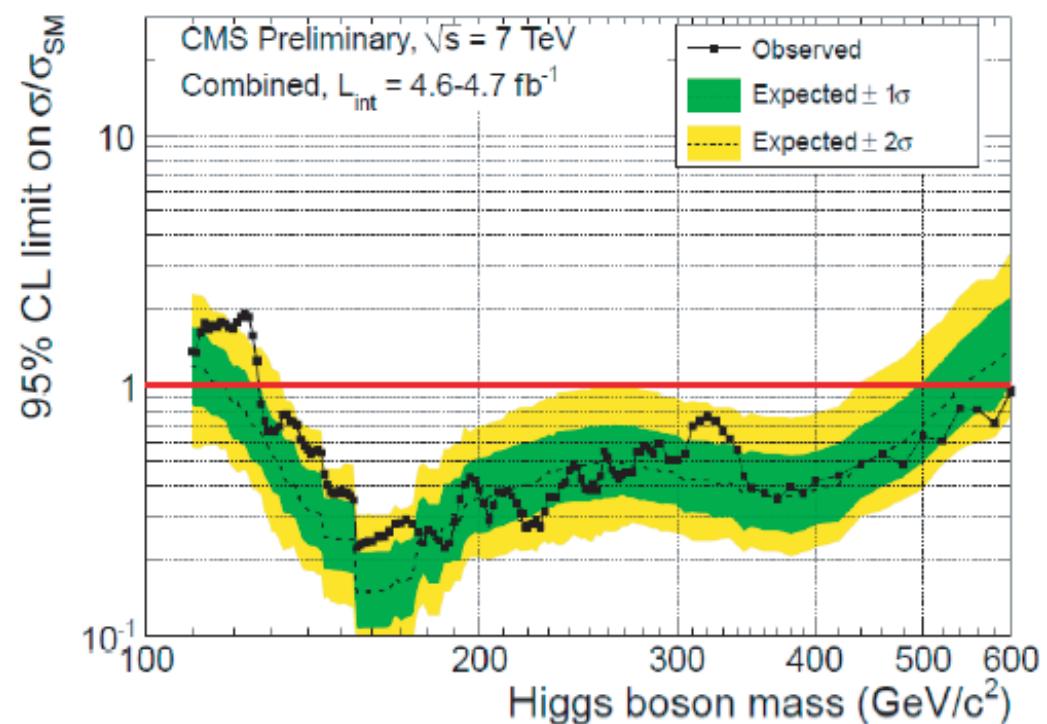
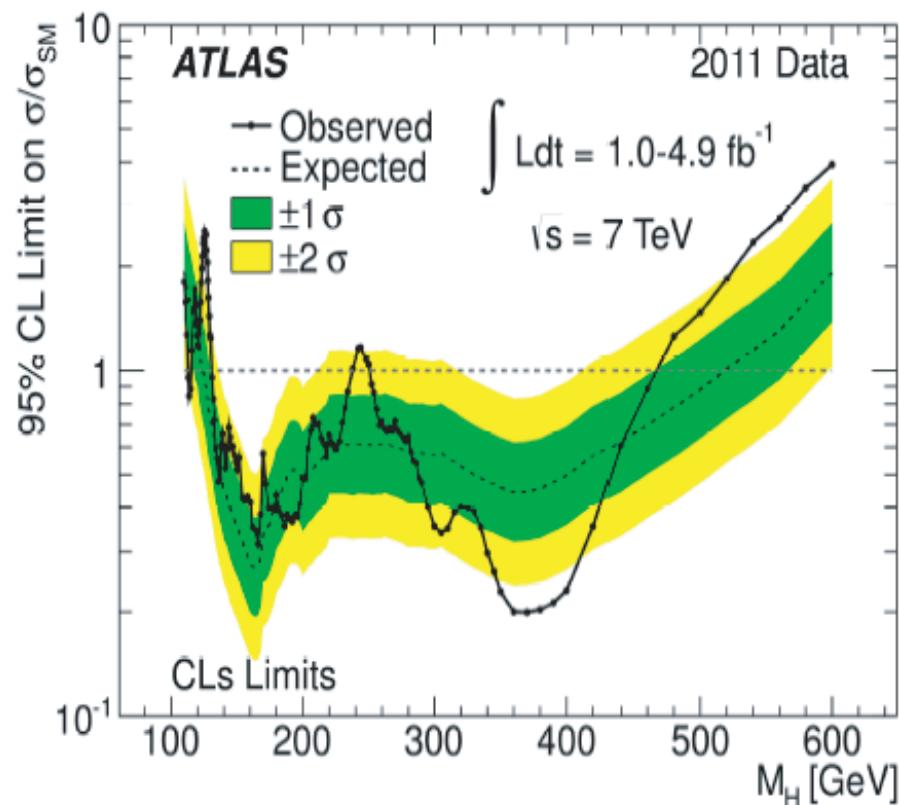
experimental evidence ?

hints of 245 GeV Higgs bosons in earlier experiments

* ATLAS-CONF-2011-135 August 21, 2011

Update of the Combination of Higgs Boson Searches

"The second was peaking at a Higgs boson mass hypothesis of 245 GeV"



Excess of events in $H \rightarrow \gamma\gamma, H \rightarrow ZZ$

ATLAS near $M_H \sim 126 \text{ GeV}$
and near $M_H \sim 245 \text{ GeV}$

CMS near $M_H \sim 124 \text{ GeV}$
and near $M_H \sim 119.5 \text{ GeV}$

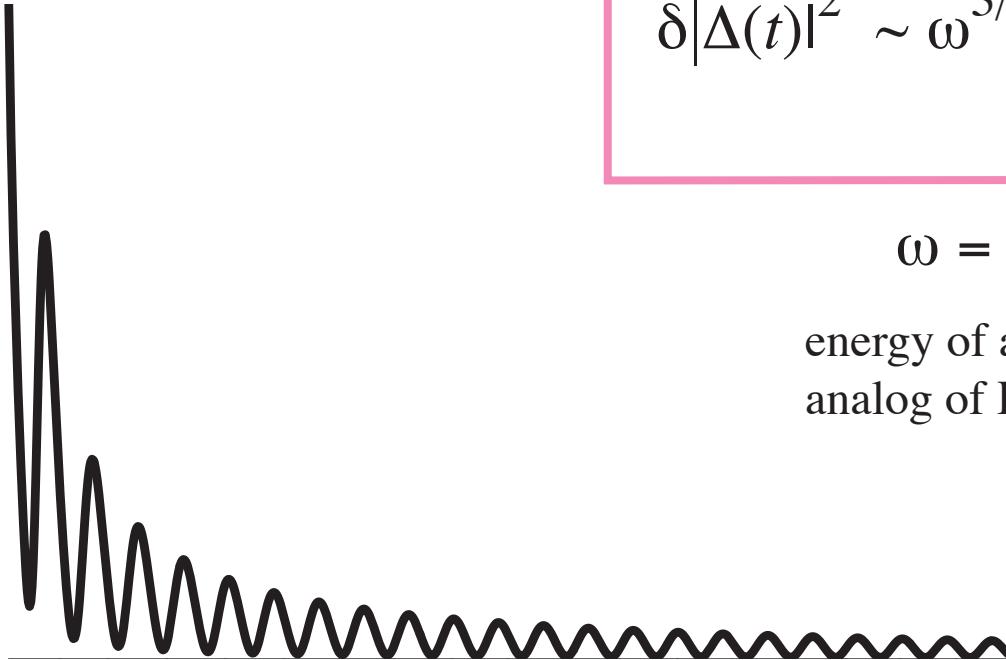
excitation of Higgs boson by quench

dynamics of vacuum energy Λ
in cosmology

$$\Lambda(t) \sim \omega^2 \frac{\sin^2 \omega t}{t^2}$$

$$\omega \sim M_{\text{inflaton}}$$

mass of Higgs inflation



Starobinsky Higgs inflation

dynamics of gap Δ
in superconductor

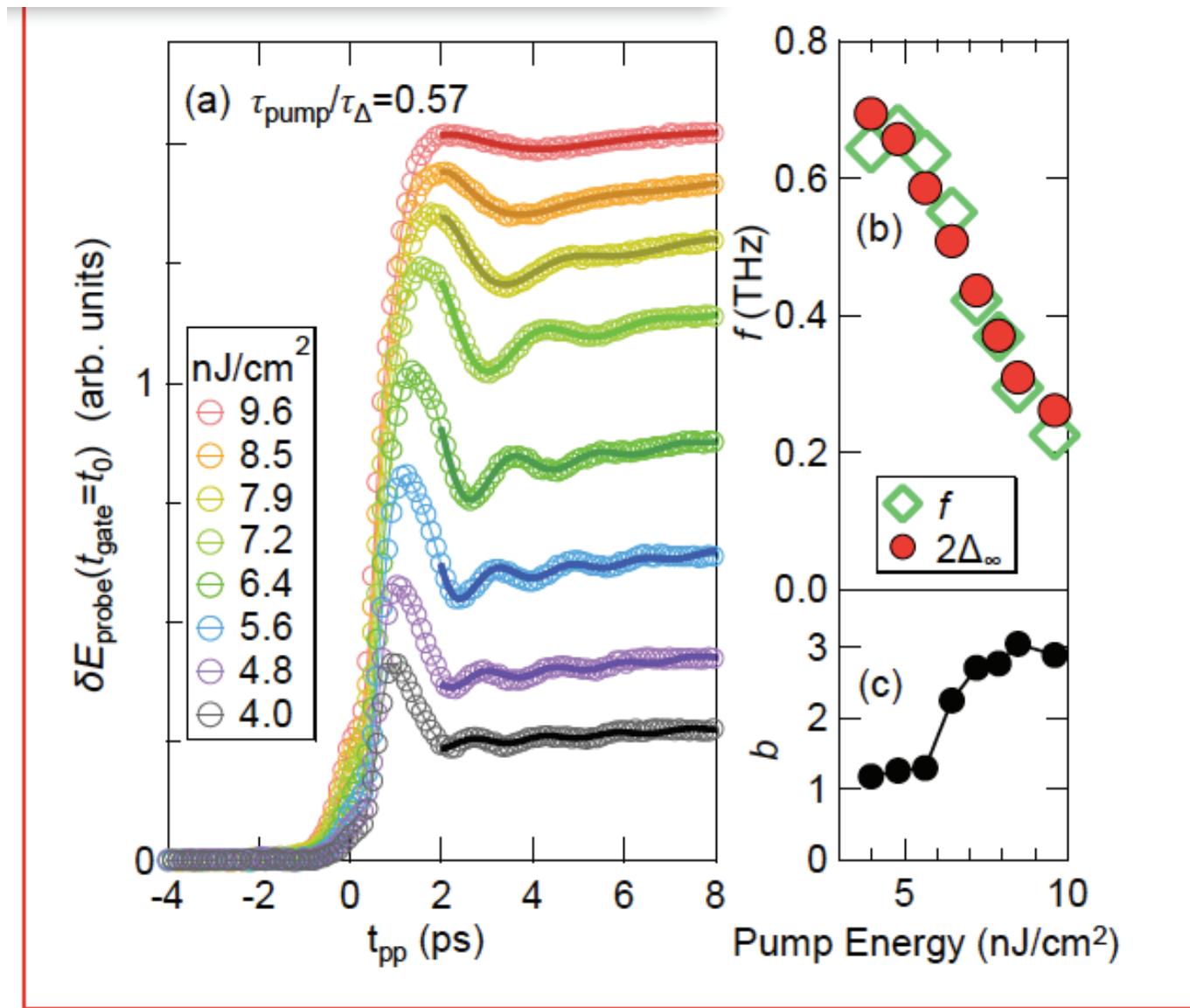
$$\delta|\Delta(t)|^2 \sim \omega^{3/2} \frac{\sin \omega t}{t^{1/2}}$$

$$\omega = 2\Delta$$

energy of amplitude mode:
analog of Higgs mass

V. Gurarie, 0905.4498, 1307.1485, 1307.2256
A.F. Volkov & S.M. Kogan, JETP **38**, 1018 (1974)
Barankov & Levitov, ...

Higgs amplitude mode after quench



Higgs Amplitude Mode in BCS Superconductors
 $\text{Nb}_{1-x}\text{Ti}_x$ induced by Terahertz Pulse Excitation
R. Matsunaga, et al. arXiv:1305.0381

Majorana fermions in particle physics & cond-mat

Majorana

(propagating fermions)



fermions in supersymmetric theories

Weyl “relativistic” fermions in bulk $^3\text{He-A}$

Majorana edge states on surface of $^3\text{He-B}$

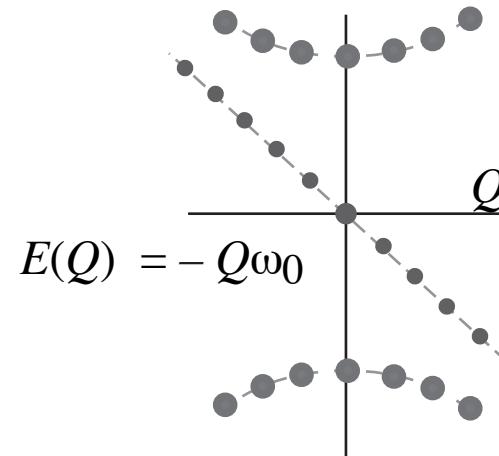
Majorana Valley: Fermi arc on surface of $^3\text{He-A}$

Majorino

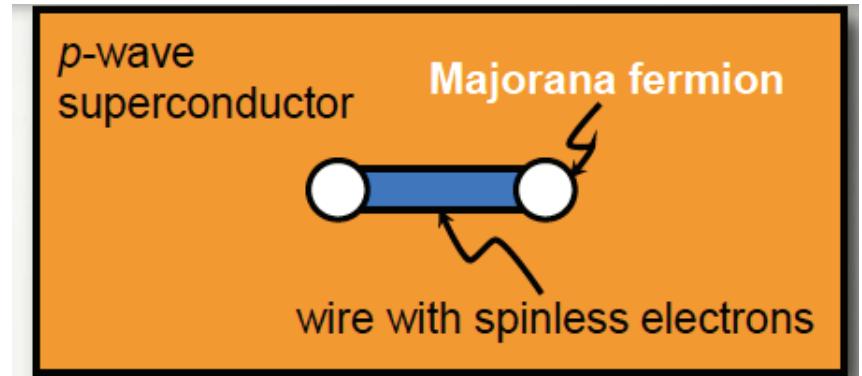
(zero-dimensional Majorana zero mode, Wilczek)

$$a^\dagger = a$$

zero mode in core of 2+1 p-wave superconductor

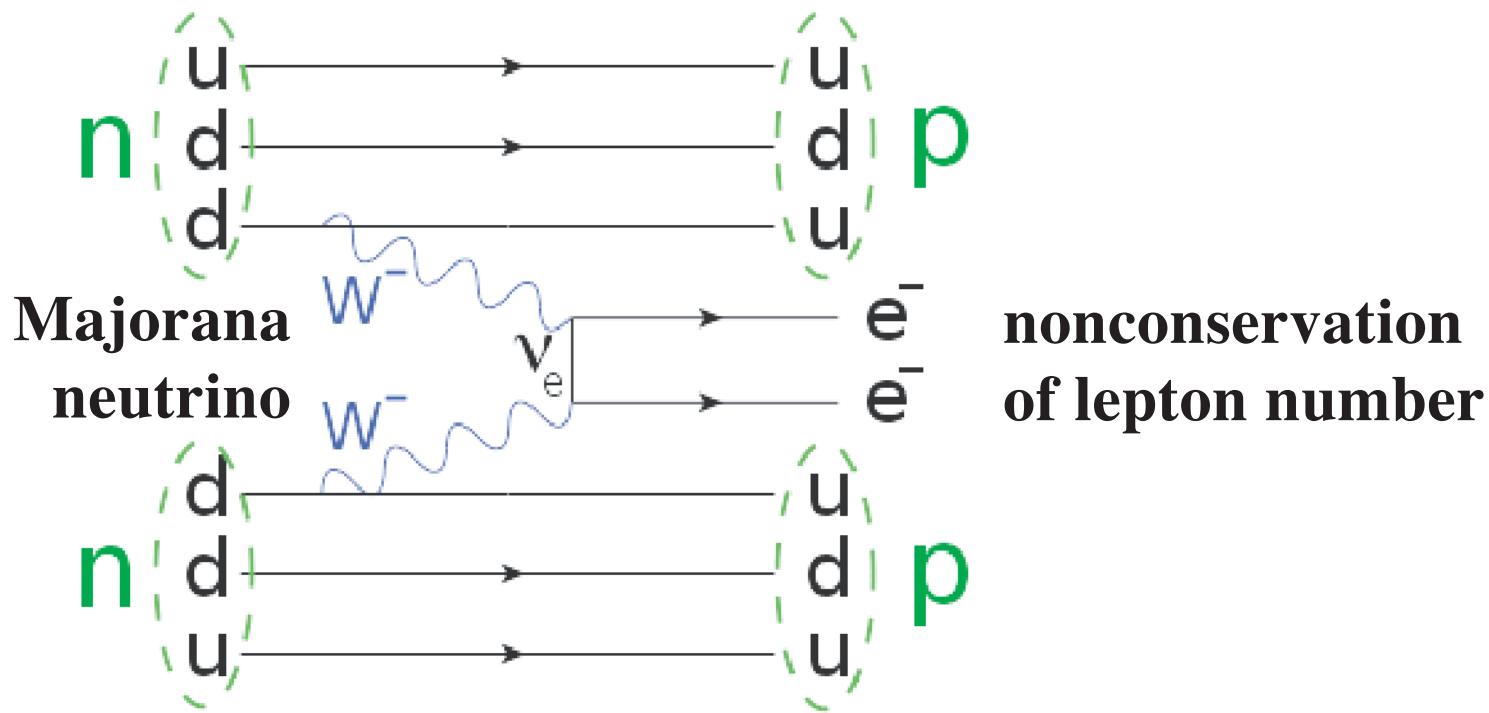


edge states in nanowires on superconductor

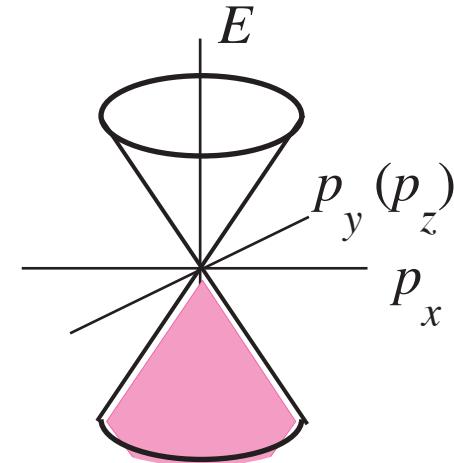
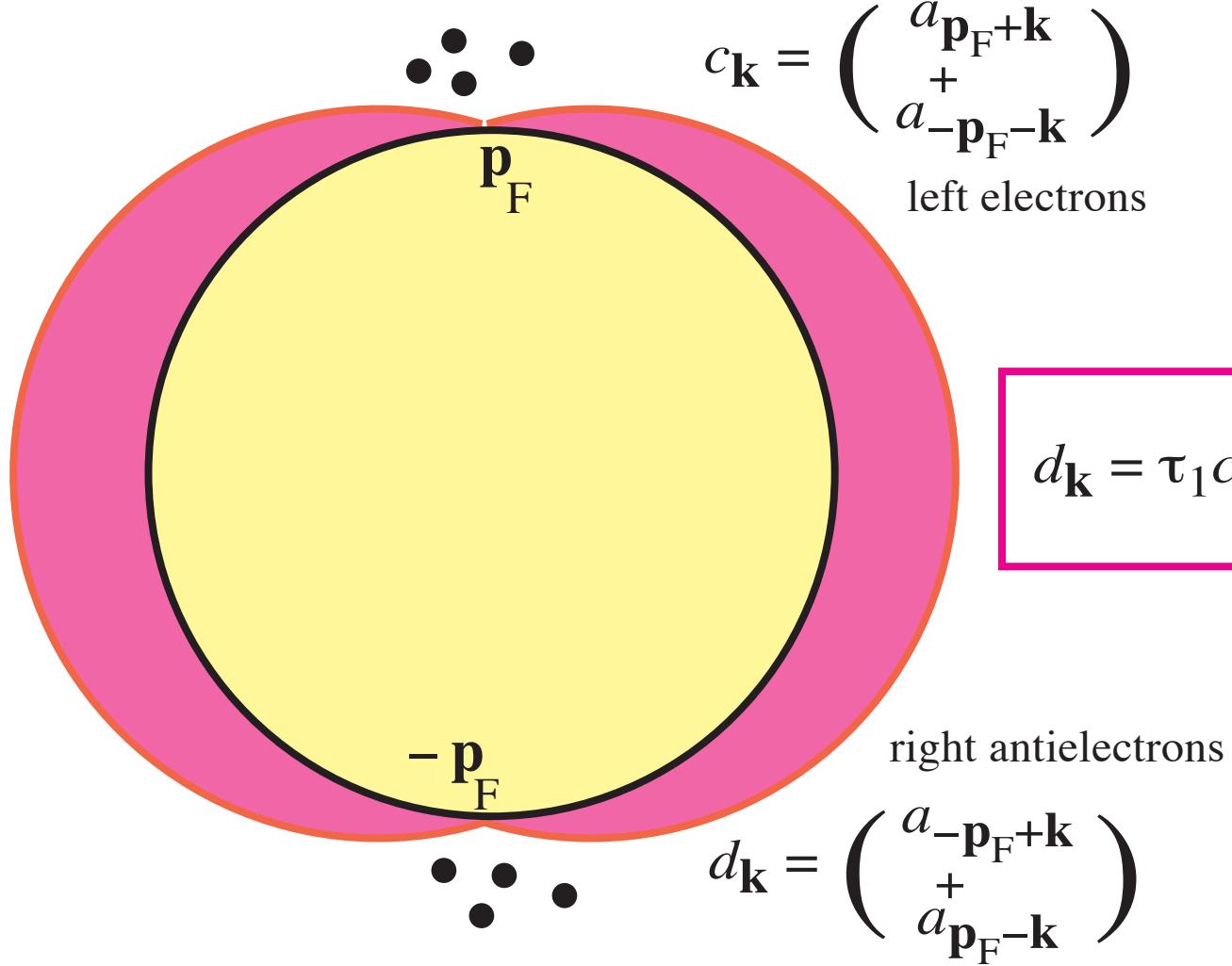


neutrinoless double beta decay

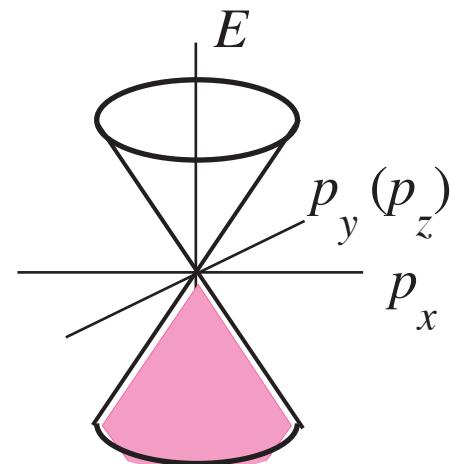
observation of neutrinoless double beta decay
necessarily implies Majorana nature of neutrinos



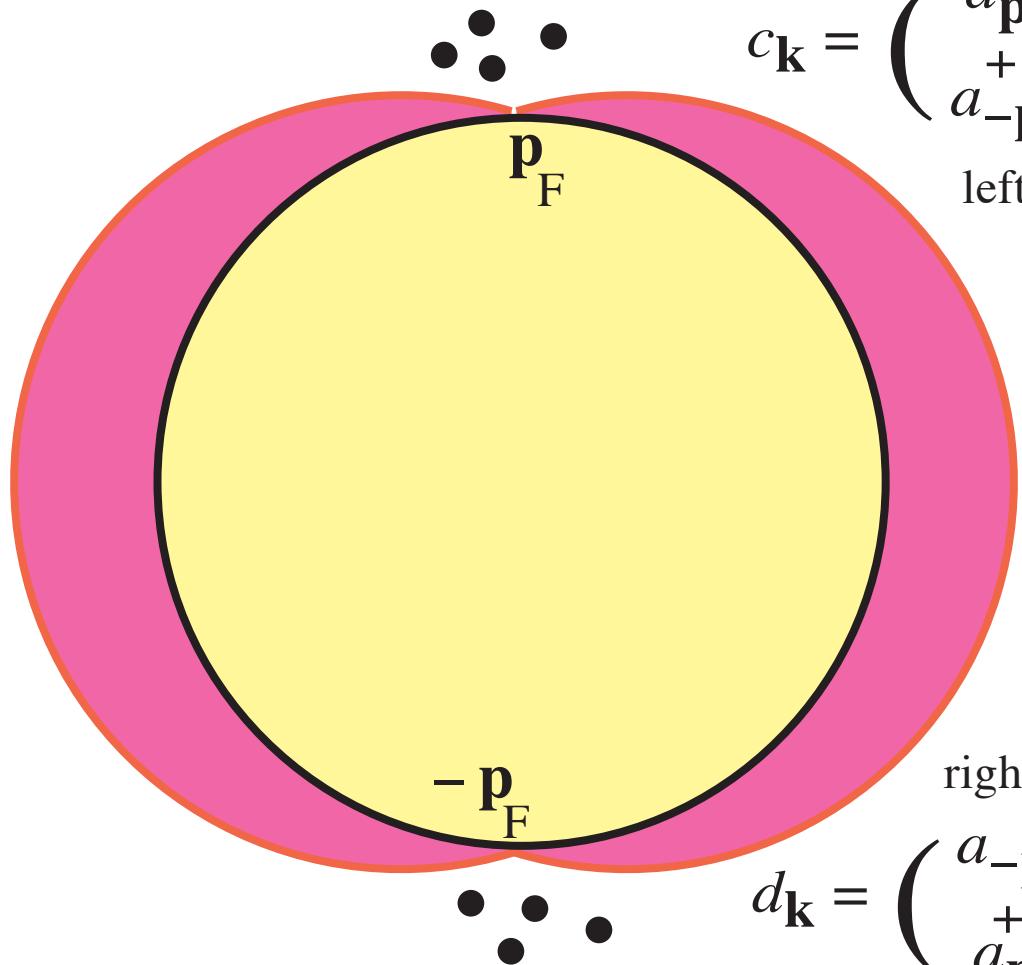
Weyl fermions in $^3\text{He-A}$



**particles in two valleys
are antiparticles**

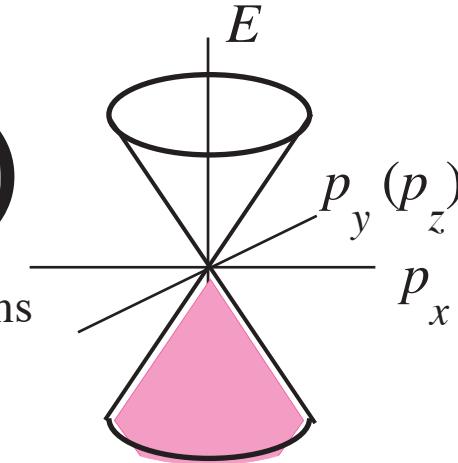


maybe all particles in our world are Majorana ?



$$c_{\mathbf{k}} = \begin{pmatrix} a_{\mathbf{p}_F + \mathbf{k}} \\ a_{-\mathbf{p}_F - \mathbf{k}}^+ \end{pmatrix}$$

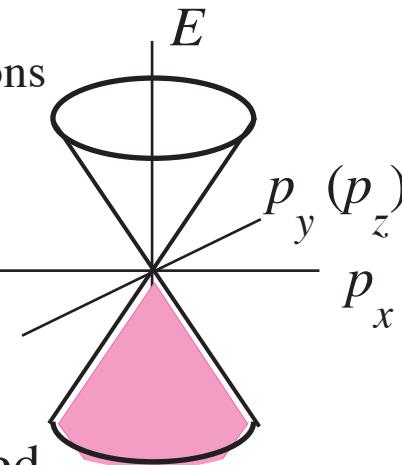
left electrons



$$d_{\mathbf{k}} = \tau_1 c_{-\mathbf{k}}^+$$

right antielectrons

$$d_{\mathbf{k}} = \begin{pmatrix} a_{-\mathbf{p}_F + \mathbf{k}} \\ a_{\mathbf{p}_F - \mathbf{k}}^+ \end{pmatrix}$$



$$H = d_{\mathbf{k}} d_{-\mathbf{k}} V_2 \mathbf{p}_F$$

in each of two worlds lepton number is not conserved
 neutrinoless double β -decay is possible
 but only due to very short wavelength perturbation

conclusion

- * composite Higgs field is extension of Standard Model:
from Ginzburg-Landau to BCS in Weyl semimetal
- * sum rule derived by Nambu for ${}^3\text{He}-\text{B}$ can be extended to Standard Model
- * Nambu sum rule (if valid) may express masses of neutral & charged Higgs bosons
via quark masses
- * more accurate analysis near 325 GeV and 245 GeV is needed
- * rigorous derivation of Nambu sum rule & its range of validity in fermionic systems:
relativistic QFT, superconductors, magnetics, cold Fermi gases, graphene,...